



Research article

Stability and bifurcation analysis of $SIQR$ for the COVID-19 epidemic model with time delay

Shishi Wang, Yuting Ding*, Hongfan Lu and Silin Gong

Department of Mathematics, Northeast Forestry University, Harbin, 150040, China

* **Correspondence:** E-mail: yuting840810@163.com.

Abstract: Based on the $SIQR$ model, we consider the influence of time delay from infection to isolation and present a delayed differential equation (DDE) according to the characteristics of the COVID-19 epidemic phenomenon. First, we consider the existence and stability of equilibria in the above delayed $SIQR$ model. Second, we analyze the existence of Hopf bifurcations associated with two equilibria, and we verify that Hopf bifurcations occur as delays crossing some critical values. Then, we derive the normal form for Hopf bifurcation by using the multiple time scales method for determining the stability and direction of bifurcation periodic solutions. Finally, numerical simulations are carried out to verify the analytic results.

Keywords: COVID-19 epidemic model; time delay; Hopf bifurcation; normal form; multiple time scales

1. Introduction

COVID-19 is a respiratory infectious disease and was first reported in Wuhan, China. It is caused by a novel coronavirus named SARS-CoV-2. SARS-CoV-2 appeared after severe acute respiratory syndrome coronavirus (SARS) and Middle East respiratory syndrome coronavirus (MERS) [1]. Coronaviruses are enveloped nonsegmented positive-sense RNA viruses and are broadly distributed in mammals, including humans. They can destroy the human respiratory systems and cause severe acute respiratory syndrome [2]. Infected individuals usually cough, have a fever, and have difficulty breathing. Some of them even die of COVID-19. Moreover, there are many ways to spread COVID-19, such as direct transmission, aerosol transmission and contact transmission. This caused COVID-19 to spread rapidly worldwide. As a global infectious disease, COVID-19 has caused more than one million deaths. We attempt to establish a model to analyze the spread of COVID-19 and provide some advice to make COVID-19 controllable.

In recent investigations, many researchers have studied different epidemic models applied to the

spread of COVID-19. In ref. [3] and ref. [4], Varotsos et al. established an *SPRD* model that included people that were susceptible, infected, recovered and dead. The model was based on COVID-19 data and was established to predict COVID-19 spread. We think there is no need to include people who died as a main part of the model because the proportion of people who died is very small. In ref. [5], León et al. developed an *SEIARD* mathematical model to investigate the COVID-19 outbreak in Mexico. The *SEIARD* model included infection with symptoms, infection without symptoms (asymptomatic), recovery from symptomatic infection and recovery from asymptomatic infection. If we consider only the infectivity of COVID-19, both infection with symptoms and infection without symptoms (asymptomatic) can infect the suspected cases. There is no need to divide them into infected with symptoms and infected without symptoms (asymptomatic). Both the recovery from symptomatic infection and the recovery from asymptomatic infection have very small possibilities for reinfection. There is no need to divide them into recovered from symptomatic infection and recovered from asymptomatic infection. In ref. [6], an *SEIR* epidemic model was developed for COVID-19. In ref. [7] and ref. [8], two *SIR* epidemic models were developed for COVID-19. For COVID-19, many countries have taken quarantine measures, which play an important part in the spread of COVID-19. These authors did not consider the influence of these quarantine measures. The quarantine measures taken by many countries make the proportion of quarantined people large. These quarantine measures can prevent infected individuals from infecting suspected individuals. In addition, exposure to air pollution can damage the heart and lungs and increase vulnerability to more serious coronavirus effects [9]. These quarantine measures can reduce the influence of air pollution. We think it is necessary to consider quarantined people because of the importance of these quarantine measures and the number of quarantined people. We attempt to construct an *SIQR* epidemic model including quarantined people to describe the spread of COVID-19.

In the spreading process of COVID-19, there is a time delay from infection to isolation. There are some studies about other epidemics with time delays. Liu et al. [10] developed an *SEIRU* epidemic model with a time delay before an infected person can transmit the infection to another person. They evaluated the effect of the latency period on the dynamics of COVID-19. Xu [11] and Yang et al. [12] also developed epidemic models with time delays before an infected person can transmit the infection to another person. Although this kind of time delay exists, human beings have difficulty changing this kind of time delay. In addition, Wang et al. [13] and Liu et al. [14] made time delays in their models the sojourn times in an infective state. Lu et al. [15] presented an *SIQR* model with a time delay from infection to recovery. The influence of the time delay from infection to recovery was discussed in detail. According to their conclusions, the smaller the time delay from infection to isolation was, the better COVID-19 was controlled. We decide to focus on the time delay from infection to isolation and discuss the effect of the time delay from infection to isolation on the spread of COVID-19, which is the main difference between our work and the works of others.

Stability and bifurcation have great significance to epidemic models, and some studies have analyzed stability and bifurcation for some epidemic models. In ref. [16], Greenhalgh investigated the stability of some *SEIRS* epidemiological models with vaccination and temporary immunity. Greenhalgh proposed that there was a threshold parameter R_0 , and the disease could persist if and only if R_0 exceeded one. In addition, disease-free equilibrium always existed and was locally stable if $R_0 < 1$ and unstable if $R_0 > 1$. However, for $R_0 > 1$, the endemic equilibrium was unique and locally asymptotically stable. Based on this conclusion, we think it is necessary for us to calculate the threshold

parameter R_0 in our model. In ref. [17], Xie et al. investigated the global stability of endemic equilibrium of an *SIS* epidemic model in complex networks. They proved that the endemic equilibrium was globally asymptotically stable by using a combination of the Lyapunov function method and a monotone iterative technique. However, the Hopf bifurcation analysis was not mentioned in this paper. In refs. [18, 19], two *SIS* models were proposed. In ref. [20], an *SIQRS* model was proposed. The stability of the models was analyzed, but bifurcation analysis was not mentioned. In refs. [21, 22, 23, 24], the stability and bifurcation of some *SIR* models were analyzed by different authors. In refs. [25, 26, 27, 28], different *SIR* models were developed. The stability and bifurcation of them were analyzed. In ref. [29], an *SIR* epidemic model with time delay was proposed. In ref. [30], an *SIR* model with information-dependent vaccination was proposed. We compare the methods they used to analyze the stability and bifurcation of these *SIR* models. Then, we use the multiple time scales method, which was also used in ref. [21]. The multiple time scales method is systematic. It is a standard method for calculating the normal form of Hopf bifurcation. It can be directly applied to the original nonlinear dynamical system, which is described by ordinary differential equations and delayed differential equations, without application of the center manifold theory [31].

It is obvious that the isolation measures taken by many countries play an important role in controlling the spread of COVID-19. The time delay also has a considerable influence on the spread. Thus, we attempt to develop one mathematical model with a time delay to investigate the spread of COVID-19 and analyze the stability of this model. In some studies such as refs. [3, 4, 5], the authors forecasted how the epidemic situation changed with detailed data. However, we decide to focus on the stability of equilibria in our model, calculate the critical time delay from infection to isolation and analyze the dynamical property of our model to discuss the tendency of the spread of COVID-19.

The main objective of this paper includes constructing the DDE model for COVID-19 and analyzing the stability of the model, the existence of Hopf bifurcation and the stability of the bifurcating periodic solution. The rest of the paper is organized as follows. In Section 2, we construct the DDE model according to the characteristics of the spread of COVID-19. In Section 3, we analyze the existence and stability of the equilibria and the existence of Hopf bifurcation associated with the COVID-19 model with a time delay. We calculate the critical time delay from infection to isolation and analyze the dynamic properties of our model. In Section 4, we derive the normal form of the Hopf bifurcation for the above model and analyze the stability of the bifurcating periodic solution. In Section 5, we present the results of simulations to verify the correctness of our analysis. Finally, the conclusion is drawn in Section 6.

2. Mathematical modeling

Suspected individuals can be infected with COVID-19 by infected individuals. After individuals are infected with COVID-19, some of them will be quarantined, and some of them will die of COVID-19. The quarantined individuals will not infect others. For COVID-19, we ignore the effect of reinfection because the probability of reinfection is very small. In addition, we ignore the probability of transforming the suspected individuals into recovered individuals directly since the probability is also small. We divide people into four kinds (suspected, infected, quarantined and recovered) based on their infectivity. The specific conversion between the four kinds of people is given in Figure 1.

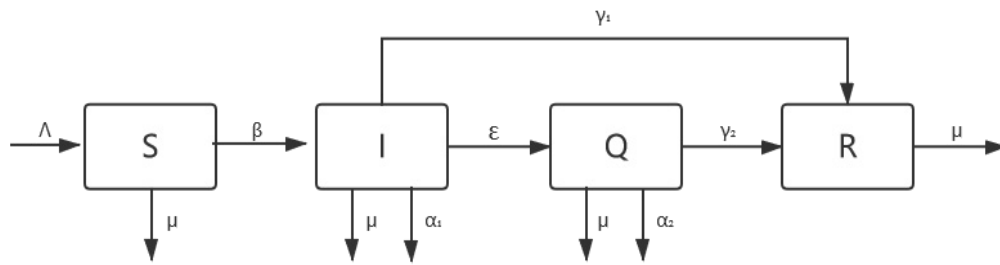


Figure 1. Flow chart for the model $SIQR$.

The variables and parameters in Figure 1 are given in Table 1.

Table 1. The definitions of parameters and variables.

Symbol	Definition
S	The number of suspected people
I	The number of infected people
Q	The number of quarantined people
R	The number of recovered people
Λ	The population growth
μ	The natural mortality
α_1	The rate of COVID-19 death of infected people
α_2	The rate of COVID-19 death of quarantined people
β	The transmission rate from S to I
ε	The transmission rate from I to Q
γ_1	The transmission rate from I to R
γ_2	The transmission rate from Q to R

For COVID-19, there is a latent period after the suspected individuals are infected. In the latent period, the infected individuals have no symptoms. However, these asymptomatic infected individuals can also infect suspected individuals. If we consider only their infectivity, there is little difference between asymptomatic infected individuals and symptomatic infected individuals. Thus, both of them can be divided into infected. Additionally, some suspected individuals may be quarantined, and they will not be infected. This means that the number of suspected persons who can be infected will decrease. Considering the influence of the suspected individuals who are quarantined, we can reduce the transmission rate from S to I , and there is no need to include the suspected individuals who are quarantined in the model.

Based on Figure 1, we construct the DDE model (2.1):

$$\begin{cases} S'(t) = \Lambda - \mu S(t) - \beta S(t)I(t), \\ I'(t) = \beta S(t)I(t) - \varepsilon I(t-\tau) - \mu I(t) - \alpha_1 I(t) - \gamma_1 I(t), \\ Q'(t) = \varepsilon I(t-\tau) - \gamma_2 Q(t) - \alpha_2 Q(t) - \mu Q(t), \\ R'(t) = \gamma_1 I(t) + \gamma_2 Q(t) - \mu R(t), \end{cases} \quad (2.1)$$

where the definitions of parameters and variables are presented in Figure 1 and Table 1, and $\tau > 0$ denotes the time delay from infection to isolation.

The initial condition of system (2.1) is $\phi_S(\theta) \geq 0, \phi_I(\theta) \geq 0, \phi_Q(\theta) \geq 0, \phi_R(\theta) \geq 0$, where $\Phi = (\phi_S(\theta), \phi_I(\theta), \phi_Q(\theta), \phi_R(\theta)) \in C([-\tau, 0], \mathbb{R}_{+0}^4)$ for $\theta \in [-\tau, 0]$, and $C([-\tau, 0], \mathbb{R}_{+0}^4)$ is the Ba-

nach space of continuous functions mapping interval $[-\tau, 0]$ into \mathbb{R}_{+0}^4 . For the system (2.1), $\mathbb{R}_{+0}^4 = \{(S, I, Q, R) | S \geq 0, I \geq 0, Q \geq 0, R \geq 0\}$.

According to the initial condition of system (2.1), we present a theorem addressing the nonnegativity and the boundedness of the solution to system(2.1).

Theorem 2.1. *If $S(0) \geq 0, I(0) \geq 0, Q(0) \geq 0, R(0) \geq 0$, the solution $S^*(t), I^*(t), Q^*(t), R^*(t)$ to system (2.1) with $\tau = 0$ is nonnegative and bounded when $t > 0$.*

Proof. First, we prove $S^*(t) \geq 0$ when $t \geq 0$ under the initial condition of system (2.1).

We assume that $S^*(t)$ is not always nonnegative for $t \geq 0$ and make t_1 the first time that $S^*(t_1) = 0$, $S'(t_1) < 0$. According to the first equation of system (2.1), we can obtain $S'(t_1) = \Lambda > 0$. The two conclusions we obtain are contradictory.

Therefore, $S^*(t) \geq 0$ when $t > 0$. In the same way, $I^*(t) \geq 0, Q^*(t) \geq 0, R^*(t) \geq 0$ when $t > 0$. The solution to system (2.1) is positive when $t > 0$.

Then, let $N(t) = S(t) + I(t) + Q(t) + R(t)$, and $N(t)$ represent the total number of people at time t .

Add four equations of Eq. (2.1), and we obtain $N'(t) = \Lambda - \mu(S(t) + I(t) + Q(t) + R(t))$. Then, we can obtain $\limsup_{t \rightarrow \infty} N(t) = \frac{\Lambda}{\mu}$. Thus, the solution $S^*(t), I^*(t), Q^*(t), R^*(t)$ to system (2.1) is bounded when $t > 0$. \square

Remark 1: We prove if $S(0) \geq 0, I(0) \geq 0, Q(0) \geq 0, R(0) \geq 0$, the solution $S^*(t), I^*(t), Q^*(t), R^*(t)$ to system (2.1) with $\tau = 0$ is positive. It is not easy for us to prove that the solution to system (2.1) is positive when $\tau > 0$. However, according to our numerical simulation, we can find that when system (2.1) is stable, the solution to system (2.1) is always positive, which is not contradictory to the positivity of the solution to system (2.1).

Next, we consider the dynamic phenomena of system (2.1).

3. Analysis of stability

3.1. Existence and stability of equilibria

In this section, the system (2.1) is considered, and we first determine the equilibria of the above system. Obviously, the system (2.1) has two equilibria:

$$E_1 = (S_1^*, I_1^*, Q_1^*, R_1^*), E_2 = (S_2^*, I_2^*, Q_2^*, R_2^*), \quad (3.1)$$

with

$$\begin{aligned} S_1^* &= \frac{\Lambda}{\mu}, I_1^* = 0, Q_1^* = 0, R_1^* = 0, S_2^* = \frac{\varepsilon + \mu + \alpha_1 + \gamma_1}{\beta}, I_2^* = \frac{\Lambda}{\varepsilon + \mu + \alpha_1 + \gamma_1} - \frac{\mu}{\beta}, \\ Q_2^* &= \frac{\varepsilon}{\alpha_2 + \gamma_2 + \mu} \left(\frac{\Lambda}{\varepsilon + \mu + \alpha_1 + \gamma_1} - \frac{\mu}{\beta} \right), R_2^* = \frac{1}{\mu} \left(\gamma_1 + \frac{\varepsilon \gamma_2}{\alpha_2 + \gamma_2 + \mu} \right) \left(\frac{\Lambda}{\varepsilon + \mu + \alpha_1 + \gamma_1} - \frac{\mu}{\beta} \right). \end{aligned} \quad (3.2)$$

Transferring the equilibria $E_k (k = 1, 2)$ to the origin, we make $w = S + S_k^*, x = I + I_k^*, y = Q + Q_k^*, z = R + R_k^*$. We can obtain $S = w - S_k^*, I = x - I_k^*, Q = y - Q_k^*$ and $R = z - R_k^*$. Substituting these equations into the model (2.1), we use S, I, Q and R to represent w, x, y and z . We obtain the following

model:

$$\begin{cases} S'(t) = -\mu S(t) - \beta S(t)I(t) - \beta S_k^* I(t) - \beta S(t)I_k^*, \\ I'(t) = \beta S(t)I(t) + \beta S_k^* I(t) + \beta S(t)I_k^* - \varepsilon I(t-\tau) - \mu I(t) - \alpha_1 I(t) - \gamma_1 I(t), \\ Q'(t) = \varepsilon I(t-\tau) - \gamma_2 Q(t) - \alpha_2 Q(t) - \mu Q(t), \\ R'(t) = \gamma_1 I(t) + \gamma_2 Q(t) - \mu R(t), k = 1, 2. \end{cases} \quad (3.3)$$

In this subsection, we first demonstrate the stability of the equilibrium $E_1 = (S_1^*, I_1^*, Q_1^*, R_1^*) = (\frac{\Lambda}{\mu}, 0, 0, 0)$. We obtain the characteristics equation of the linearized system of Eq. (3.3) at the equilibrium E_1 as follows:

$$(\lambda + \mu)^2 (\lambda + \mu + \alpha_2 + \gamma_2) (\lambda + \mu + \alpha_1 + \gamma_1 - \beta S_1^* + \varepsilon e^{-\lambda\tau}) = 0. \quad (3.4)$$

When $\tau = 0$, Eq. (3.4) becomes:

$$(\lambda + \mu)^2 (\lambda + \mu + \alpha_2 + \gamma_2) (\lambda + \mu + \alpha_1 + \gamma_1 - \beta S_1^* + \varepsilon) = 0. \quad (3.5)$$

We show the following assumption:

$$(H1) \quad \beta S_1^* - (\varepsilon + \mu + \alpha_1 + \gamma_1) < 0.$$

Under (H1), the four roots of Eq. (3.5) have negative real parts due to $\mu > 0, \alpha_1 > 0, \gamma_1 > 0$, and the equilibrium E_1 is locally asymptotically stable when $\tau = 0$.

Similarly, for the stability of the other equilibrium $E_2 = (S_2^*, I_2^*, Q_2^*, R_2^*)$, we obtain the characteristic equation of the linearized system of Eq. (3.3) at the equilibrium E_2 as follows:

$$(\lambda + \mu) (\lambda + \mu + \alpha_2 + \gamma_2) [(\lambda - \varepsilon + \varepsilon e^{-\lambda\tau}) (\lambda + \mu + \beta I_2^*) + \beta^2 S_2^* I_2^*] = 0. \quad (3.6)$$

When $\tau = 0$, Eq. (3.6) becomes:

$$(\lambda + \mu) (\lambda + \mu + \alpha_2 + \gamma_2) [\lambda^2 + \lambda(\mu + \beta I_2^*) + \beta^2 S_2^* I_2^*] = 0. \quad (3.7)$$

If (H1) is not satisfied, the four roots of Eq. (3.7) have negative real parts due to $\mu > 0, \alpha_2 > 0, \gamma_2 > 0$, and the equilibrium E_2 of system (2.1) is locally asymptotically stable when $\tau = 0$.

We can find that E_1 is the disease-free equilibrium, and E_2 is the endemic equilibrium. We calculate the basic reproduction number R_0 , which is the number of suspected individuals who are infected by the same infectious individual and can estimate the infectiousness of an infectious disease. According to the system (2.1), we can obtain the new infections matrix \mathcal{F} and the transition matrix \mathcal{V} .

$$\mathcal{F} = \begin{bmatrix} 0 \\ \beta S(t)I(t) \\ 0 \\ 0 \end{bmatrix}, \mathcal{V} = \begin{bmatrix} -\Lambda + \mu S(t) + \beta S(t)I(t) \\ \varepsilon I(t) + \mu I(t) + \alpha_1 I(t) + \gamma_1 I(t) \\ -\varepsilon I(t) + \gamma_2 Q(t) + \alpha_2 Q(t) + \mu Q(t) \\ -\gamma_1 I(t) - \gamma_2 Q(t) + \mu R(t) \end{bmatrix}.$$

Then, we make F_0 represent the derivative of \mathcal{F} at E_1 and V_0 represent the derivative of \mathcal{V} at E_1 :

$$F_0 = \begin{bmatrix} 0 & 0 & 0 & 0 \\ 0 & \beta S_1^* & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix}, V_0 = \begin{bmatrix} \mu & \beta S_1^* & 0 & 0 \\ 0 & \varepsilon + \mu + \alpha_1 + \gamma_1 & 0 & 0 \\ 0 & -\varepsilon & \gamma_2 + \alpha_2 + \mu & 0 \\ 0 & -\gamma_1 & -\gamma_2 & \mu \end{bmatrix}.$$

The inverse of V_0 is:

$$V_0^{-1} = \begin{bmatrix} \frac{1}{\mu} & \frac{-\beta S_1^*}{\mu(\varepsilon + \mu + \alpha_1 + \gamma_1)} & 0 & 0 \\ 0 & \frac{1}{\varepsilon + \mu + \alpha_1 + \gamma_1} & 0 & 0 \\ 0 & \frac{\varepsilon}{(\varepsilon + \mu + \alpha_1 + \gamma_1)(\gamma_2 + \alpha_2 + \mu)} & \frac{1}{\mu(\gamma_2 + \alpha_2 + \mu)} & 0 \\ 0 & \frac{1}{\mu} \left[\frac{\gamma_1}{\varepsilon + \mu + \alpha_1 + \gamma_1} + \frac{\varepsilon \gamma_2}{(\varepsilon + \mu + \alpha_1 + \gamma_1)(\gamma_2 + \alpha_2 + \mu)} \right] & \frac{\gamma_2 + \alpha_2 + \mu}{\mu(\gamma_2 + \alpha_2 + \mu)} & \frac{1}{\mu} \end{bmatrix}.$$

And we can obtain:

$$F_0 V_0^{-1} = \begin{bmatrix} 0 & 0 & 0 & 0 \\ 0 & \frac{\beta S_1^*}{\varepsilon + \mu + \alpha_1 + \gamma_1} & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix}.$$

The maximum eigenvalue of $F_0 V_0^{-1}$ is R_0 :

$$R_0 = \rho(F_0 V_0^{-1}) = \frac{\beta S_1^*}{\varepsilon + \mu + \alpha_1 + \gamma_1}.$$

We compare R_0 with (H1), and we find that if $R_0 < 1$, the inequality of (H1) exists constantly, and if $R_0 > 1$, (H1) is not satisfied. This represents if $R_0 < 1$, the disease-free equilibrium E_1 is locally asymptotically stable when $\tau = 0$ and if $R_0 > 1$, the endemic equilibrium E_2 is locally asymptotically stable when $\tau = 0$.

3.2. Existence of Hopf bifurcation

For equilibrium E_1 , when $\tau > 0$, let $\lambda = i\omega_0$ ($\omega_0 > 0$) be a root of the following equation:

$$\lambda + \mu + \alpha_1 + \gamma_1 - \beta S_1^* + \varepsilon e^{-\lambda\tau} = 0. \quad (3.8)$$

Substituting $\lambda = i\omega_0$ into Eq. (3.8) and separating the real and imaginary parts, we have

$$\begin{cases} \omega_0 = \varepsilon \sin(\omega_0\tau), \\ \beta S_1^* - (\mu + \alpha_1 + \gamma_1) = \varepsilon \cos(\omega_0\tau). \end{cases} \quad (3.9)$$

Thus, we obtain:

$$\begin{cases} Q_1 \triangleq \sin(\omega_0\tau) = \frac{\omega_0}{\varepsilon}, \\ P_1 \triangleq \cos(\omega_0\tau) = \frac{\beta S_1^* - (\mu + \alpha_1 + \gamma_1)}{\varepsilon}. \end{cases} \quad (3.10)$$

Adding the square of two equations (3.9), we can obtain

$$\omega_0^2 = \varepsilon^2 - [\beta S_1^* - (\mu + \alpha_1 + \gamma_1)]^2. \quad (3.11)$$

Thus, $\omega_0 = \sqrt{\varepsilon^2 - [\beta S_1^* - (\mu + \alpha_1 + \gamma_1)]^2}$ when $\varepsilon^2 - [\beta S_1^* - (\mu + \alpha_1 + \gamma_1)]^2 \geq 0$, and substituting it into Eq. (3.10). Since $\omega_0 > 0$ and $\varepsilon \geq 0$, we can obtain $Q_1 > 0$ and obtain:

$$\tau_1^{(j)} = \frac{1}{\omega_0} [\arccos(P_1) + 2j\pi], \quad Q_1 \geq 0, \quad (3.12)$$

where Q_1, P_1 are given in Eq. (3.10), $j = 0, 1, 2, \dots$.

Let $\lambda(\tau) = \alpha(\tau) + i\omega(\tau)$ be the root of Eq. (3.8) satisfying $\alpha(\tau_1^{(j)}) = 0, \omega(\tau_1^{(j)}) = \omega_0, j = 0, 1, 2, \dots$. Then, we have the transversality conditions

$$\operatorname{Re}\left(\frac{d\lambda}{d\tau}\right)^{-1}\bigg|_{\lambda=i\omega_0, \tau=\tau_1^{(j)}} = \frac{1}{\varepsilon^2} > 0, j = 0, 1, 2, \dots$$

Thus, the system (2.1) undergoes a Hopf bifurcation at equilibrium E_1 of system (2.1) when $\tau = \tau_1^{(j)}, j = 0, 1, 2, \dots$.

For the other equilibrium E_2 , when $\tau > 0$, let $\lambda = i\omega$ ($\omega > 0$) be the root of the following equation:

$$(\lambda - \varepsilon + \varepsilon e^{-\lambda\tau})(\lambda + \mu + \beta I_2^*) + \beta^2 S_2^* I_2^* = 0. \quad (3.13)$$

Substituting it into the equation and separating the real and imaginary parts, we have

$$\begin{cases} \omega^2 + \mu\varepsilon + \varepsilon\beta I_2^* - \beta^2 S_2^* I_2^* = (\omega\varepsilon) \sin(\omega\tau) + (\mu\varepsilon + \beta\varepsilon I_2^*) \cos(\omega\tau), \\ \mu\omega - \varepsilon\omega + \beta\omega I_2^* = (-\omega\varepsilon) \cos(\omega\tau) + (\mu\varepsilon + \beta\varepsilon I_2^*) \sin(\omega\tau). \end{cases} \quad (3.14)$$

Equation (3.14) can be presented as:

$$\begin{cases} Q_2 \triangleq \sin(\omega\tau) = \frac{(\omega^2 + \varepsilon(\beta I_2^* + \mu) - \beta^2 S_2^* I_2^*)\omega - (\varepsilon - (\beta I_2^* + \mu))(\beta I_2^* + \mu)\omega}{\varepsilon[\omega^2 + (\beta I_2^* + \mu)^2]}, \\ P_2 \triangleq \cos(\omega\tau) = \frac{(\omega^2 + \varepsilon(\beta I_2^* + \mu) - \beta^2 S_2^* I_2^*)(\beta I_2^* + \mu) + (\varepsilon - (\beta I_2^* + \mu))\omega^2}{\varepsilon[\omega^2 + (\beta I_2^* + \mu)^2]}. \end{cases} \quad (3.15)$$

Adding the square of two equations Eq. (3.14), let $z = \omega^2$, and

$$h(z) = z^2 + c_1 z + c_0 = 0, \quad (3.16)$$

where $c_1 = (\beta I_2^* + \mu)^2 - 2\beta^2 S_2^* I_2^*, c_0 = (\beta^2 S_2^* I_2^*)^2 - 2\beta^2 S_2^* I_2^* (\beta I_2^* + \mu)\varepsilon$.

Therefore, we have the following assumption:

(H2) $c_1^2 - 4c_0 > 0, c_1 < 0, c_0 > 0$.

(H3) $c_0 < 0$.

If (H2) holds, then Eq. (3.16) has two positive roots, z_1 and z_2 ($z_1 < z_2$). If (H3) holds, then Eq. (3.16) has only one positive root z_3 . Without loss of generality, substituting $\omega_k = \sqrt{z_k}$ ($k = 1, 2, 3$) into Eq. (3.15), we obtain:

$$\tau_{2,k}^{(j)} = \begin{cases} \frac{1}{\omega_k} [\arccos(P_{2,k}) + 2j\pi], & Q_{2,k} \geq 0, \\ \frac{1}{\omega_k} [2\pi - \arccos(P_{2,k}) + 2j\pi], & Q_{2,k} < 0, j = 0, 1, 2, \dots, \end{cases} \quad (3.17)$$

where Q_2, P_2 are given in Eq. (3.15) and $Q_{2,k} = Q_2|_{\omega=\omega_k, \tau=\tau_{2,k}^{(j)}}, P_{2,k} = P_2|_{\omega=\omega_k, \tau=\tau_{2,k}^{(j)}}, k = 1, 2, 3$.

Note that since $Q_1 = \frac{\omega_0}{\varepsilon}$ is always nonnegative, we can obtain $\tau_1^{(j)}$ directly. However, whether Q_2 is positive depends on the parameters of Eq. (3.15), and we need to obtain $\tau_{2,k}^{(j)}$ ($j = 1, 2, 3$) according to the value of Q_2 .

Furthermore, let $\lambda(\tau) = \alpha(\tau) + i\omega_k(\tau)$ be the root of Eq. (3.17) satisfying $\alpha(\tau_{2,k}^{(j)}) = 0, \omega(\tau_{2,k}^{(j)}) = \omega_k$ ($k = 1, 2, 3; j = 0, 1, 2, \dots$).

If (H2) or (H3) holds and $z_k = \omega_k^2 (k = 1, 2, 3)$, then we can deduce $\text{Re}\left(\frac{d\tau}{d\lambda}\right)\Big|_{\tau=\tau_{2,k}^{(j)}}$:

$$\text{Re}\left(\frac{d\tau}{d\lambda}\right)\Big|_{\tau=\tau_{2,k}^{(j)}} = \frac{\omega_k^2 [(\beta I_2^* + \mu)^2 - 2\beta^2 S_2^* I_2^*]}{\varepsilon^2 [\omega_k^2 + (\beta I_2^* + \mu)^2]} = \frac{\omega_k^2 h'(z_k)}{\varepsilon^2 [\omega_k^2 + (\beta I_2^* + \mu)^2]} \neq 0, k = 1, 2, 3.$$

There is a Hopf bifurcation at equilibrium E_2 of system (2.1) when $\tau = \tau_{2,k}^{(j)}$, $j = 0, 1, 2, \dots$.

Then, we present a stability theorem of system (2.1)'s equilibria E_1 , E_2 and the existence of Hopf bifurcation.

Theorem 3.1. *We consider the model (2.1):*

(1) *If (H1) holds, the equilibrium E_1 of the model (2.2) undergoes a Hopf bifurcation when $\tau = \tau_1^{(j)}$, $j = 0, 1, 2, \dots$, and we obtain the following. When $\tau \in [0, \tau_1^{(0)})$, the equilibrium E_1 is locally asymptotically stable. When $\tau \in [\tau_1^{(0)}, +\infty)$, the equilibrium E_1 is locally asymptotically unstable.*

(2) *Under the condition that (H1) is not satisfied, and if (H2) or (H3) holds, the equilibrium E_2 of the model (2.1) undergoes a Hopf bifurcation when $\tau = \tau_{2,k}^{(j)}$ ($k = 1, 2, 3$), and we obtain the following.*

(a) *If (H2) holds, $h(z)$ has two positive roots, z_1 and z_2 ; then we assume $z_1 < z_2$, and we obtain $h'(z_1) < 0$, $h'(z_2) > 0$. Then, $\exists m \in \mathbb{N}$, which can make $0 < \tau_{2,2}^{(0)} < \tau_{2,1}^{(0)} < \tau_{2,2}^{(1)} < \tau_{2,1}^{(1)} < \dots < \tau_{2,1}^{(m-1)} < \tau_{2,2}^{(m)} < \tau_{2,2}^{(m+1)}$. When $\tau \in [0, \tau_{2,2}^{(0)}) \cup \bigcup_{l=1}^m (\tau_{2,1}^{(l-1)}, \tau_{2,2}^{(l)})$, the equilibrium of the model is locally asymptotically*

stable. When $\tau \in \bigcup_{l=0}^{m-1} (\tau_{2,2}^{(l)}, \tau_{2,1}^{(l)}) \cup (\tau_{2,2}^{(m)}, +\infty)$, the equilibrium is locally asymptotically unstable.

(b) *If (H3) holds, $h(z)$ has only one positive root z_3 , when $\tau \in [0, \tau_{2,3}^{(0)})$, the equilibrium E_2 is locally asymptotically stable. When $\tau \in (\tau_{2,3}^{(0)}, +\infty)$, the equilibrium E_2 is unstable.*

4. Normal form of Hopf bifurcation

Without loss of generality, we denote the critical value $\tau = \tau^*$ when the characteristic Eq. (3.4) and Eq. (3.6) have eigenvalue $\lambda = i\omega_k^*$, $k = 1, 2$, at which system (2.1) undergoes a Hopf bifurcation at equilibrium $E_k = (S_k^*, I_k^*, Q_k^*, R_k^*)$, $k = 1, 2$.

First, we transform the equilibrium to the origin; then, we use the multiple time scales method, and system (2.1) can be written as

$$\dot{X}(t) = AX(t) + BX(t - \tau) + F(X(t), X(t - \tau)), \quad (4.1)$$

where $X(t) = (\hat{S}_k(t), \hat{I}_k(t), \hat{Q}_k(t), \hat{R}_k(t))^T$, $X(t - \tau) = (\hat{S}_k(t - \tau), \hat{I}_k(t - \tau), \hat{Q}_k(t - \tau), \hat{R}_k(t - \tau))^T$,

$$A = \begin{bmatrix} -\mu - \beta I_k^* & -\beta S_k^* & 0 & 0 \\ \beta I_k^* & \beta S_k^* - (\mu + \alpha_1 + \gamma_1) & 0 & 0 \\ 0 & 0 & -(\gamma_2 + \alpha_2 + \mu) & 0 \\ 0 & \gamma_1 & \gamma_2 & -\mu \end{bmatrix}, B = \begin{bmatrix} 0 & 0 & 0 & 0 \\ 0 & -\varepsilon & 0 & 0 \\ 0 & \varepsilon & 0 & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix},$$

$$F(X(t), X(t - \tau)) = \begin{bmatrix} F_S \\ F_I \\ F_Q \\ F_R \end{bmatrix} = \begin{bmatrix} \Lambda - \mu S_k^* - \beta S_k^* I_k - \beta S_k^* I_k^* \\ \beta S_k^* I_k + \beta S_k^* I_k^* - (\varepsilon + \mu + \alpha_1 + \gamma_1) I_k^* \\ \varepsilon I_k^* - (\mu + \alpha_2 + \gamma_2) Q_k^* \\ \gamma_1 I_k^* + \gamma_2 Q_k^* - \mu R_k^* \end{bmatrix}.$$

We make $\widetilde{X}(t) = AX(t) + BX(t - \tau)$ the linear system of system (4.1). We assume that $h_k (k = 1, 2)$ is the eigenvector of the eigenvalue $\lambda = i\omega_k^* (k = 1, 2)$ of the linear system and that h_k^* is the eigenvector of the eigenvalue $\lambda = -i\omega_k^*$ of the linear system. We make \widetilde{I} the unit matrix, and h_k satisfies $(\widetilde{I} - A - Be^{-\lambda\tau})h_k = 0$, where $\lambda = i\omega_k^*$, and h_k^* satisfies $(\widetilde{I} - A - Be^{-\lambda\tau})^T h_k^* = 0$, where $\lambda = -i\omega_k^*$. Additionally, $\langle h_k^*, h_k \rangle = \overline{h_k^*}^T h_k = 1$, and we can obtain

$$h_k = (h_{k1}, h_{k2}, h_{k3}, h_{k4})^T = \left(\frac{-\beta S_k^*}{\lambda + \mu + \beta I_k^*}, 1, \frac{\epsilon e^{-\lambda\tau}}{\lambda + \mu + \alpha_2 + \gamma_2}, \frac{\gamma_1}{\lambda + \mu} + \frac{\gamma_2 \epsilon e^{-\lambda\tau}}{(\lambda + \mu)(\lambda + \mu + \alpha_2 + \gamma_2)} \right)^T, \quad (4.2)$$

$$h_k^* = (h_{k1}^*, h_{k2}^*, h_{k3}^*, h_{k4}^*)^T = d_k \left(\frac{\beta I_k^*}{\lambda + \mu + \beta I_k^*}, 1, 0, 0 \right)^T,$$

where $d_k = \left(\frac{-\beta^2 S_k^* I_k^*}{(\mu + \beta I_k^*)^2 + \omega_k^{*2}} + 1 \right)^{-1}$, $k = 1, 2$.

We use the multiple time scale method to deduce the normal form of the Hopf bifurcation and assume the solution to Eq. (4.1) is

$$X(t) = X(T_0, T_1, T_2, \dots) = \sum_{k=1}^{\infty} \epsilon^k X_k(T_0, T_1, T_2, \dots), \quad (4.3)$$

where $X(T_0, T_1, T_2, \dots) = [S(T_0, T_1, T_2, \dots), I(T_0, T_1, T_2, \dots), Q(T_0, T_1, T_2, \dots), R(T_0, T_1, T_2, \dots)]^T$, $X_k(T_0, T_1, T_2, \dots) = [S_k(T_0, T_1, T_2, \dots), I_k(T_0, T_1, T_2, \dots), Q_k(T_0, T_1, T_2, \dots), R_k(T_0, T_1, T_2, \dots)]^T$, $k = 1, 2$, $T_i = \epsilon^i t$, $i = 0, 1, 2, \dots$ and T_i is the scaling transform in the time direction.

The derivative with respect to t is

$$\frac{d}{dt} = \frac{\partial}{\partial T_0} + \epsilon \frac{\partial}{\partial T_1} + \epsilon^2 \frac{\partial}{\partial T_2} + \dots = D_0 + \epsilon D_1 + \epsilon^2 D_2 + \dots,$$

where $D_i = \frac{\partial}{\partial T_i}$, $i = 0, 1, 2, \dots$.

Note that $X_j = (S_j, I_j, Q_j, R_j)^T = X_j(T_0, T_1, T_2, \dots)$, $X_{j,\tau_c} = (S_{j,\tau_c}, I_{j,\tau_c}, Q_{j,\tau_c}, R_{j,\tau_c})^T = X_j(T_0 - \tau_c, T_1, T_2, \dots)$, $j = 1, 2, \dots$.

Then, we can obtain

$$\dot{X}(t) = \epsilon D_0 X_1 + \epsilon^2 D_1 X_1 + \epsilon^3 D_2 X_1 + \epsilon^2 D_0 X_2 + \epsilon^3 D_1 X_2 + \epsilon^3 D_0 X_3 + \dots \quad (4.4)$$

We consider that τ is the bifurcation parameter, and we set $\tau = \tau_c + \epsilon\tau_\epsilon$, where $\tau_c = \tau_k^{(j)}$ ($k = 1, 2$) is the critical value of the Hopf bifurcation, τ_ϵ is the perturbation parameter, and ϵ is the dimensionless parameter.

Using Taylor series expansion of $X(t - \tau)$, we obtain that

$$X(t - \tau) = \epsilon X_{1,\tau_c} + \epsilon^2 X_{2,\tau_c} + \epsilon^3 X_{3,\tau_c} - \epsilon^2 \tau_\epsilon D_0 X_{1,\tau_c} - \epsilon^3 \tau_\epsilon D_0 X_{2,\tau_c} - \epsilon^2 \tau_c D_1 X_{1,\tau_c} - \epsilon^3 \tau_\epsilon D_1 X_{1,\tau_c} - \epsilon^3 \tau_c D_2 X_{1,\tau_c} - \epsilon^3 \tau_c D_1 X_{2,\tau_c} + \dots, \quad (4.5)$$

where $X_{j,\tau_c} = X_j(T_0 - \tau_c, T_1, T_2, \dots)$, $j = 1, 2, \dots$.

Then, substituting the solution with the multiple scales Eqs. (4.3)-(4.5) into Eq. (4.1) and balancing the coefficients of ϵ^n ($n = 1, 2, 3$), we obtain a set-ordered linear differential equation:

$$\begin{cases} D_0 S_{k1} + (\mu + \beta I_k^*) S_{k1} + \beta S_k^* I_{k1} = 0, \\ D_0 I_{k1} - \beta I_k^* S_{k1} + (\mu + \alpha_1 + \gamma_1) I_{k1} - \beta S_k^* I_{k1} + \epsilon I_{k1,\tau_c} = 0, \\ D_0 Q_{k1} + (\mu + \alpha_2 + \gamma_2) Q_{k1} - \epsilon Q_{k1,\tau_c} = 0, \\ D_0 R_{k1} - \gamma_1 I_{k1} - \gamma_2 Q_{k1} + \mu R_{k1} = 0, k = 1, 2. \end{cases} \quad (4.6)$$

Since $\pm i\omega_k^*$ ($k = 1, 2$) are the eigenvalues of the characteristic Eq. (4.1), the solution to Eq. (4.6) can be expressed in the form of Eq. (4.7):

$$X_1(T_1, T_2, T_3, \dots) = G_1(T_1, T_2, T_3, \dots) e^{i\omega_k^* T_0} h_k + \bar{G}_1(T_1, T_2, T_3, \dots) e^{-i\omega_k^* T_0} \bar{h}_k, k = 1, 2. \quad (4.7)$$

where G_1 and \bar{G}_1 represent the coordinates on the center manifold, and the linear part of the center manifold can be presented as equation (4.7).

Next, for the ϵ^2 -order terms, we can obtain the following equations:

$$\begin{cases} D_0 S_{k2} + (\mu + \beta I_k^*) S_{k2} + \beta S_k^* I_{k2} = -D_1 S_{k1} - \beta S_{k1} I_{k1}, \\ D_0 I_{k2} - \beta I_k^* S_{k2} + (\mu + \alpha_1 + \gamma_1) I_{k2} - \beta S_k^* I_{k2} + \epsilon I_{k2,\tau_c} = \epsilon \tau_c D_0 I_{k1,\tau_c} + \epsilon \tau_c D_1 I_{k1,\tau_c} \\ \qquad \qquad \qquad \qquad \qquad \qquad \qquad \qquad \qquad \qquad \qquad \qquad \qquad \qquad \qquad \qquad - D_1 I_{k1} + \beta S_{k1} I_{k1}, \\ D_0 Q_{k2} + (\mu + \alpha_2 + \gamma_2) Q_{k2} - \epsilon Q_{k2,\tau_c} = -\epsilon \tau_c D_0 Q_{k1,\tau_c} - \epsilon \tau_c D_1 Q_{k1,\tau_c} - D_1 Q_{k1}, \\ D_0 R_{k2} - \gamma_1 I_{k2} - \gamma_2 Q_{k2} + \mu R_{k2} = -D_1 R_{k1}, k = 1, 2. \end{cases} \quad (4.8)$$

We substitute Eq. (4.7) into the right expression of Eq. (4.8), and the coefficients before $e^{i\omega_k^* T_0}$ are denoted by the vector m_1 . According to the solvability condition $\langle h_k^*, m_1 \rangle = 0$, we can solve $\frac{\partial G}{\partial T_1}$ as follows:

$$\frac{\partial G_1}{\partial T_1} = M_k \tau_c G_1, \quad (4.9)$$

where $M_k = \epsilon e^{-i\omega_k^* \tau_c} \left[\frac{\beta I_k^* (\omega_k^* + \mu + \beta I_k^*)}{(\mu + \beta I_k^*)^2 + \omega_k^{*2}} \frac{\beta S_k^*}{i\omega_k^* + \mu + \beta I_k^*} + 1 - \epsilon \tau_c e^{-i\omega_k^* \tau_c} \right]^{-1}$, $k = 1, 2$.

We assume:

$$\begin{aligned} S_{k2} &= f_{k1} e^{i\omega_k^* T_0} G_1 + \bar{f}_{k1} e^{-i\omega_k^* T_0} \bar{G}_1 + g_{k1} e^{2i\omega_k^* T_0} G_1^2 + \bar{g}_{k1} e^{-2i\omega_k^* T_0} \bar{G}_1^2 + l_{k1} G_1 \bar{G}_1, \\ I_{k2} &= f_{k2} e^{i\omega_k^* T_0} G_1 + \bar{f}_{k2} e^{-i\omega_k^* T_0} \bar{G}_1 + g_{k2} e^{2i\omega_k^* T_0} G_1^2 + \bar{g}_{k2} e^{-2i\omega_k^* T_0} \bar{G}_1^2 + l_{k2} G_1 \bar{G}_1, \\ Q_{k2} &= f_{k3} e^{i\omega_k^* T_0} G_1 + \bar{f}_{k3} e^{-i\omega_k^* T_0} \bar{G}_1 + g_{k3} e^{2i\omega_k^* T_0} G_1^2 + \bar{g}_{k3} e^{-2i\omega_k^* T_0} \bar{G}_1^2 + l_{k3} G_1 \bar{G}_1, \\ R_{k2} &= f_{k4} e^{i\omega_k^* T_0} G_1 + \bar{f}_{k4} e^{-i\omega_k^* T_0} \bar{G}_1 + g_{k4} e^{2i\omega_k^* T_0} G_1^2 + \bar{g}_{k4} e^{-2i\omega_k^* T_0} \bar{G}_1^2 + l_{k4} G_1 \bar{G}_1. \end{aligned} \quad (4.10)$$

Substituting Eq. (4.10) into Eq. (4.8), we obtain:

$$\begin{aligned}
f_{k1} &= U^{-1} \tau_\varepsilon (i\omega_k^* + \mu + \alpha_2 + \gamma_2 - \varepsilon e^{-i\omega_k^* \tau_c}) [(-M_k h_{k1})(i\omega_k^* + \mu + \alpha_1 + \gamma_1 - \beta S_k^* + \varepsilon e^{-i\omega_k^* \tau_c}) \\
&\quad - \beta S_k^* (i\omega_k^* \varepsilon e^{-i\omega_k^* \tau_c} h_{k2} + \varepsilon \tau_c M_k h_{k2} - M_k h_{k2})] (i\omega_k^* + \mu), \\
f_{k2} &= U^{-1} \tau_\varepsilon (i\omega_k^* + \mu + \alpha_2 + \gamma_2 - \varepsilon e^{-i\omega_k^* \tau_c}) [(i\omega_k^* + \mu + \beta I_k^*) (i\omega_k^* \varepsilon e^{-i\omega_k^* \tau_c} h_{k2} + \varepsilon \tau_c M_k h_{k2} \\
&\quad - M_k h_{k2}) + \beta I_k^* (-M_k h_{k1})] (i\omega_k^* + \mu), \\
f_{k3} &= \tau_\varepsilon (i\omega_k^* + \mu) [(i\omega_k^* + \mu + \beta I_k^*) (i\omega_k^* + \mu + \alpha_1 + \gamma_1 - \beta S_k^* + \varepsilon e^{-i\omega_k^* \tau_c}) (i\omega_k^* + \mu + \alpha_2 - \varepsilon e^{-i\omega_k^* \tau_c} \\
&\quad + \gamma_2) + \beta^2 S_k^* I_k^* (i\omega_k^* + \mu + \alpha_2 + \gamma_2 - \varepsilon e^{-i\omega_k^* \tau_c})] U^{-1}, \\
f_{k4} &= \tau_\varepsilon \{ (i\omega_k^* + \mu + \beta I_k^*) [(-M_k h_{k4})(i\omega_k^* + \mu + \alpha_1 + \gamma_1 - \beta S_k^* + \varepsilon e^{-i\omega_k^* \tau_c}) (i\omega_k^* + \mu - \varepsilon e^{-i\omega_k^* \tau_c} \\
&\quad + \alpha_2 + \gamma_2) + \gamma_1 (i\omega_k^* \varepsilon e^{-i\omega_k^* \tau_c} h_{k2} + \varepsilon \tau_c M_k h_{k2} - M_k h_{k2}) (i\omega_k^* + \mu + \alpha_2 + \gamma_2 - \varepsilon e^{-i\omega_k^* \tau_c}) \\
&\quad + \gamma_2 (i\omega_k^* + \mu + \alpha_1 + \gamma_1 - \beta S_k^* + \varepsilon e^{-i\omega_k^* \tau_c}) (-M_k h_{k3} - \varepsilon \tau_c M_k e^{-i\omega_k^* \tau_c} h_{k3} - i\omega_k^* \varepsilon e^{-i\omega_k^* \tau_c} h_{k3})] \\
&\quad + \beta I_k^* [\beta S_k^* (i\omega_k^* + \mu + \alpha_2 + \gamma_2 - \varepsilon e^{-i\omega_k^* \tau_c}) (-M_k h_{k4}) + \gamma_1 (i\omega_k^* \varepsilon e^{-i\omega_k^* \tau_c} h_{k2} + \varepsilon \tau_c M_k h_{k2} \\
&\quad - M_k h_{k2}) (i\omega_k^* + \mu + \alpha_2 + \gamma_2 - \varepsilon e^{-i\omega_k^* \tau_c}) + \gamma_2 \beta S_k^* (-M_k h_{k3} - i\omega_k^* \varepsilon M_k e^{-i\omega_k^* \tau_c} h_{k3} \\
&\quad - \varepsilon e^{-i\omega_k^* \tau_c} h_{k3})] \} U^{-1}, \\
g_{k1} &= (2i\omega_k^* + \mu) (2i\omega_k^* + \mu + \alpha_2 + \gamma_2 - \varepsilon e^{-2i\omega_k^* \tau_c}) [(-\beta h_{k1} h_{k2}) (2i\omega_k^* + \mu + \alpha_1 + \gamma_1 + \varepsilon e^{-2i\omega_k^* \tau_c} \\
&\quad - \beta S_k^*) + \beta^2 S_k^* h_{k1} h_{k2}] V^{-1}, \\
g_{k2} &= (2i\omega_k^* + \mu) (2i\omega_k^* + \mu + \alpha_2 + \gamma_2 - \varepsilon e^{-2i\omega_k^* \tau_c}) [\beta h_{k1} h_{k2} (2i\omega_k^* + \mu + \beta I_k^*) - \beta^2 I_k^* h_{k1} h_{k2}] V^{-1}, \\
g_{k3} &= 0, \\
g_{k4} &= (2i\omega_k^* + \mu + \alpha_2 + \gamma_2 - \varepsilon e^{-2i\omega_k^* \tau_c}) [(2i\omega_k^* + \mu + \beta I_k^*) (\beta h_{k1} \bar{h}_{k2} + \beta h_{k2} \bar{h}_{k1}) + \beta I_k^* (-\beta h_{k1} \bar{h}_{k2} \\
&\quad - \beta h_{k2} \bar{h}_{k1})] \gamma_1 V^{-1}, \\
l_{k1} &= \mu (\mu + \alpha_2 + \gamma_2 - \varepsilon) [(-\beta h_{k1} \bar{h}_{k2} - \beta h_{k2} \bar{h}_{k1}) (\mu + \alpha_1 + \gamma_1 - \beta S_k^* + \varepsilon) + \beta S_k^* (\beta h_{k1} \bar{h}_{k2} \\
&\quad + \beta h_{k2} \bar{h}_{k1})] W^{-1}, \\
l_{k2} &= W^{-1} \mu (\mu + \alpha_2 + \gamma_2 - \varepsilon) [(\beta h_{k1} \bar{h}_{k2} + \beta h_{k2} \bar{h}_{k1}) (\mu + \beta I_k^*) + \beta I_k^* (-\beta h_{k1} \bar{h}_{k2} - \beta h_{k2} \bar{h}_{k1})], \\
l_{k3} &= 0, \\
l_{k4} &= (\mu + \alpha_2 + \gamma_2 - \varepsilon) \gamma_1 [(\beta h_{k1} \bar{h}_{k2} + \beta h_{k2} \bar{h}_{k1}) (\mu + \beta I_k^*) + \beta I_k^* (-\beta h_{k1} \bar{h}_{k2} - \beta h_{k2} \bar{h}_{k1})] W^{-1}, \\
U &= (i\omega_k^* + \mu) (i\omega_k^* + \mu + \alpha_2 + \gamma_2 - \varepsilon e^{-i\omega_k^* \tau_c}) [(i\omega_k^* + \mu + \beta I_k^*) (i\omega_k^* + \mu + \alpha_1 + \gamma_1 - \beta S_k^* \\
&\quad + \varepsilon e^{-i\omega_k^* \tau_c}) + \beta^2 S_k^* I_k^*], \\
V &= (2i\omega_k^* + \mu) (2i\omega_k^* + \mu + \alpha_2 + \gamma_2 - \varepsilon e^{-2i\omega_k^* \tau_c}) [(2i\omega_k^* + \mu + \beta I_k^*) (2i\omega_k^* + \mu + \alpha_1 + \gamma_1 - \beta S_k^* \\
&\quad + \varepsilon e^{-2i\omega_k^* \tau_c}) + \beta^2 S_k^* I_k^*], \\
W &= \mu (\mu + \alpha_2 + \gamma_2 - \varepsilon) [(\mu + \beta I_k^*) (\mu + \alpha_1 + \gamma_1 - \beta S_k^* + \varepsilon) + \beta^2 S_k^* I_k^*], \quad k = 1, 2.
\end{aligned} \tag{4.11}$$

Next, balancing the ε^3 -order terms, we can obtain equations as follows:

$$\begin{cases} D_0 S_3 + (\mu + \beta I_k^*) S_3 + \beta S_k^* I_3 = -D_2 S_1 - D_1 S_2 - \beta S_2 I_1 - \beta S_1 I_2, \\ D_0 I_3 - \beta I_k^* S_3 + (\mu + \alpha_1 + \gamma_1) I_3 - \beta S_k^* I_3 + \varepsilon I_{3,\tau_c} = -D_2 I_1 - D_1 I_2 + \beta S_2 I_1 + \beta S_1 I_2 \\ \quad + \varepsilon[\tau_\varepsilon(D_0 I_{2,\tau_c} - D_1 I_{1,\tau_c}) - \tau_c(D_2 I_{1,\tau_c} + D_1 I_{2,\tau_c})], \\ D_0 Q_3 + (\mu + \alpha_2 + \gamma_2) Q_3 - \varepsilon Q_{3,\tau_c} = -\varepsilon(\tau_\varepsilon D_0 Q_{2,\tau_c} - \tau_\varepsilon D_1 Q_{1,\tau_c} - \tau_c D_2 Q_{1,\tau_c} - \tau_c D_1 Q_{2,\tau_c}) \\ \quad - D_2 Q_1 - D_1 Q_2, \\ D_0 R_3 - \gamma_1 I_3 - \gamma_2 Q_3 + \mu R_3 = -D_2 R_1 - D_1 R_2, k = 1, 2. \end{cases} \quad (4.12)$$

Substituting Eq. (4.7) and Eq. (4.10) into the right expression of Eq. (4.12), and the coefficients before $e^{i\omega_k T_0}$ are denoted by the vector m_2 . According to the solvability condition $\langle h_1^*, m_2 \rangle = 0$ and note that τ_ε^2 is small enough for small unfolding parameter τ_ε , we ignore the term $\tau_\varepsilon^2 G_1$. Then, we have:

$$\frac{\partial G_1}{\partial T_2} = H_k \chi_k G_1^2 \bar{G}_1, \quad (4.13)$$

where

$$H_k = \frac{1}{\frac{\beta I_k^* (\mu + \beta I_k^* - i\omega_k^*)}{(\mu + \beta I_k^*)^2 + \omega_k^{*2}} h_{k1} + h_{k2} + \varepsilon \tau_c e^{-i\omega_k^* \tau_c} h_{k2}},$$

$$\chi_k = -\frac{\beta I_k^* (\mu + \beta I_k^* - i\omega_k^*)}{(\mu + \beta I_k^*)^2 + \omega_k^{*2}} \beta g_{k1} \bar{h}_{k2} - \frac{\beta I_k^* (\mu + \beta I_k^* - i\omega_k^*)}{(\mu + \beta I_k^*)^2 + \omega_k^{*2}} \beta g_{k2} \bar{h}_{k2} + \beta g_{k1} \bar{h}_{k2} + \beta g_{k2} \bar{h}_{k2}, k = 1, 2,$$

where g_k are given in Eq. (4.11) and h_k are given in Eq. (4.2).

Let $G_1 \mapsto (G_1/\varepsilon)$, and we can obtain the normal form of Hopf bifurcation of system (2.1) as:

$$\dot{G}_1 = M_k \tau_\varepsilon G_1 + H_k \chi_k G_1^2 \bar{G}_1, \quad (4.14)$$

where M_k is given in Eq. (4.9), and H_k, χ_k are given in Eq. (4.13).

Let $G = r e^{i\theta_k^*}$ ($k = 1, 2$) and substitute it into Eq. (4.14), and we can obtain the normal form of Hopf bifurcation in polar coordinates:

$$\begin{cases} \dot{r} = \text{Re}(M_k) \tau_\varepsilon r + \text{Re}(H_k \chi_k) r^3, \\ \dot{\theta}_k^* = \text{Im}(M_k) \tau_\varepsilon + \text{Im}(H_k \chi_k) r^2, \end{cases} \quad (4.15)$$

where M_k is given in Eq. (4.9), and H_k, χ_k are given in Eq. (4.13) ($k = 1, 2$).

According to the normal form of bifurcation in polar coordinates, there is a theorem as follows:

Theorem 4.1. For system (4.15), if $\frac{\text{Re}(M_k) \tau_\varepsilon}{\text{Re}(H_k \chi_k)} < 0$ ($k = 1, 2$) holds, then system (2.1) exists periodic solutions near equilibrium E_k :

- (1) If $\text{Re}(M_k) \tau_\varepsilon < 0$, the bifurcating periodic solutions are unstable.
- (2) If $\text{Re}(M_k) \tau_\varepsilon > 0$, the bifurcating periodic solutions are locally asymptotically stable.

5. Numerical simulations

In this section, according to the natural mortality in China and the data presented in ref. [32], we choose $\Lambda = 100, \mu = 0.00713, \varepsilon = 0.8, \gamma_1 = 0.1, \gamma_2 = 0.1, \alpha_1 = 0.025, \alpha_2 = 0.025, \beta = 0.00001$. According to Eq. (3.2), we obtain $S_1^* = \frac{\Lambda}{\mu} = 14025, I_1^* = 0, Q_1^* = 0, R_1^* = 0, S_2^* = 93213, I_2^* = -605.719, Q_2^* = -3667.4, R_2^* = -59932$. Obviously, assumption (H1) holds and equilibrium E_1 is locally asymptotically stable when $\tau = 0$. However, the equilibrium E_2 is unstable for $I_2^* < 0$.

Remark 2: If $I_2^* < 0$, the number of infected individuals is less than zero, which is inconsistent with the facts. Additionally, according to our theoretical analysis, when $I_2^* < 0$, the equilibrium E_2 is unstable, which is consistent with the facts.

Substituting these parameter values into Eqs. (3.10)-(3.12), by using MATLAB, we can obtain $\omega_0 = 0.8000, Q_1 = 0.0102, P_1 = 0.9999, \tau_1^{(0)} = 1.9509$.

Thus, the equilibrium E_1 is locally asymptotically stable when $\tau \in [0, \tau_1^{(0)})$, and Hopf bifurcation occurs near the equilibrium E_1 when $\tau = \tau_1^{(0)}$. According to Eq. (4.9) and Eq. (4.13), we obtain $\text{Re}(M_1) < 0, \text{Re}(H_1\chi_1) > 0$. The bifurcating periodic solution is unstable due to Theorem 4.1. We find that if the time delay from infection to isolation is over the critical value, the epidemic will not be controlled under a stable situation, and the epidemic situation will be more severe.

When $\tau = 0$, we choose the initial value (1500, 1000, 100, 200), and the solution corresponding to a locally asymptotically stable equilibrium is shown in Figure 2. Thus, when $\tau = 0$, once the individuals are infected, they will be quarantined and they will not infect others, and the epidemic situation can be controlled into a stable situation.

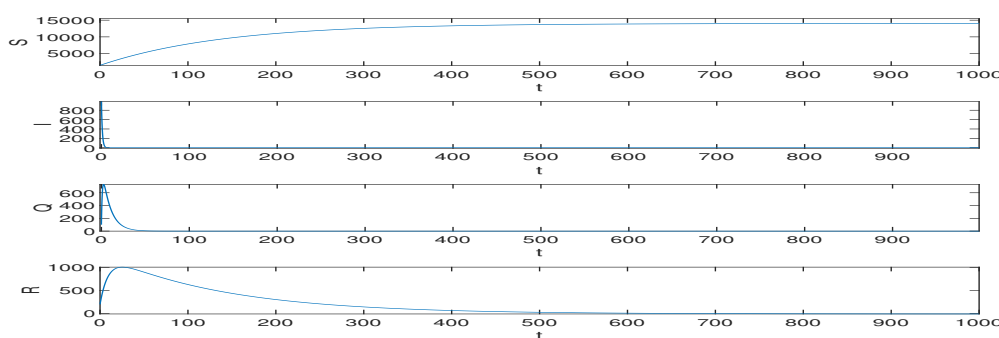


Figure 2. When $\tau = 0$, the equilibrium E_1 of system (2.1) is locally asymptotically stable.

When $\tau = 1.5 \in (0, \tau_1^{(0)})$, we choose initial values (1, 500, 1, 000, 100, 200), and the equilibrium E_1 of system (2.1) is locally asymptotically stable. See Figure 3.

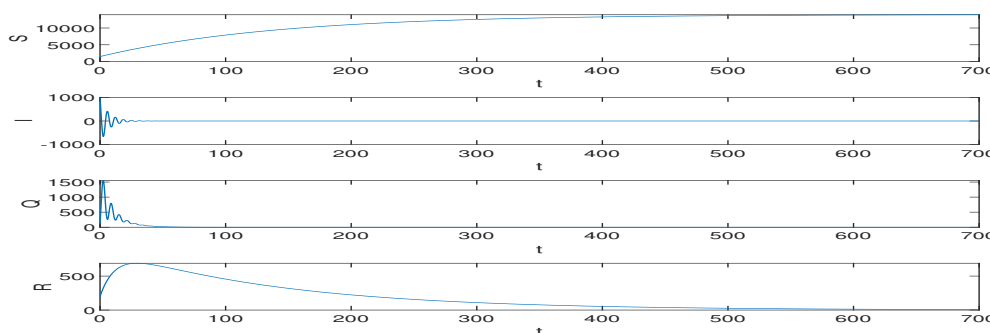


Figure 3. When $\tau = 1.5$, the equilibrium E_1 of system (2.1) is locally asymptotically stable.

When $\tau = 2 \in (\tau_1^{(0)}, +\infty)$, we choose the initial values $(1, 500, 1, 000, 100, 200)$, and the equilibrium E_1 is unstable, as shown in Figure 4. This means that if τ is larger than the critical value $\tau_1^{(0)}$, the epidemic situation cannot be controlled into a stable state.

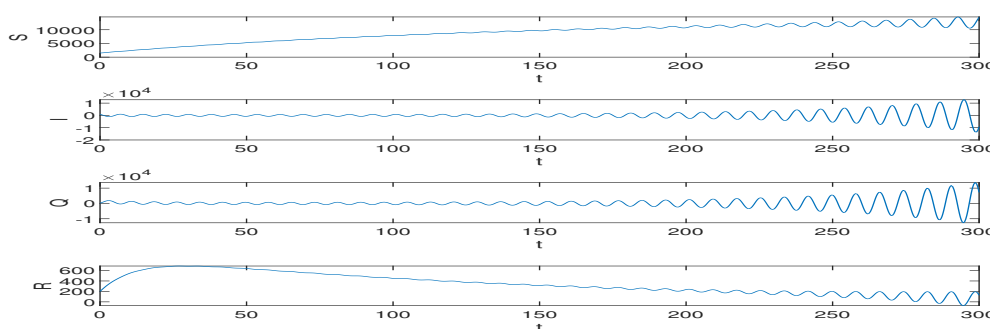


Figure 4. When $\tau = 2$, the equilibrium E_1 of system (2.1) is unstable.

According to Figures 2-4, we find that the time delay from infection to isolation has a great influence on the epidemic situation in E_1 , and with the increasing time delay, the epidemic situation will become increasingly severe.

According to the natural mortality in China and the data presented in ref. [32], we choose another group of parameters, namely, $\Lambda = 200, \mu = 0.00713, \varepsilon = 0.8, \gamma_1 = 0.05, \gamma_2 = 0.05, \alpha_1 = 0.025, \alpha_2 = 0.025, \beta = 0.0001$. According to Eq. (3.2), we obtain $S_1^* = 28,050, I_1^* = 0, Q_1^* = 0, R_1^* = 0, S_2^* = 8,821.300, I_2^* = 155.424, Q_2^* = 1,513.9, R_2^* = 11,703$. We can find that assumption (H1) does not hold and the equilibrium E_1 is unstable. However, the equilibrium E_2 is locally asymptotically stable when $\tau = 0$.

Substituting these parameter values into Eq. (3.16), we can obtain $c_1 = -0.0269, c_2 = -0.00030938$, and Eq. (3.16) has only one positive root, $z_1 = 0.0356$. According to Eq. (3.15) and Eq. (3.17), we can obtain $\omega_2 = 0.1887, Q_2 = 0.9892, P_1 = 0.1463, \tau_{2,3}^{(0)} = 0.7782$ by using MATLAB.

Thus, the equilibrium E_2 is locally asymptotically stable for $\tau \in [0, \tau_{2,3}^{(0)})$, and Hopf bifurcation occurs near the equilibrium E_2 when $\tau = \tau_{2,3}^{(0)}$. According to Eq. (4.9) and Eq. (4.13), we obtain $\text{Re}(M_1) > 0, \text{Re}(H_1 \chi_1) > 0$. Thus, $\text{Re}(M_1) \tau_\varepsilon < 0$, the bifurcation solution is unstable. In E_2 , the epidemic situation cannot be controlled under a stable situation.

When $\tau = 0$, we choose the initial values $(1,000, 50, 2,000, 1,000)$, and the solution corresponding to a locally asymptotically stable equilibrium is shown in Figure 5.

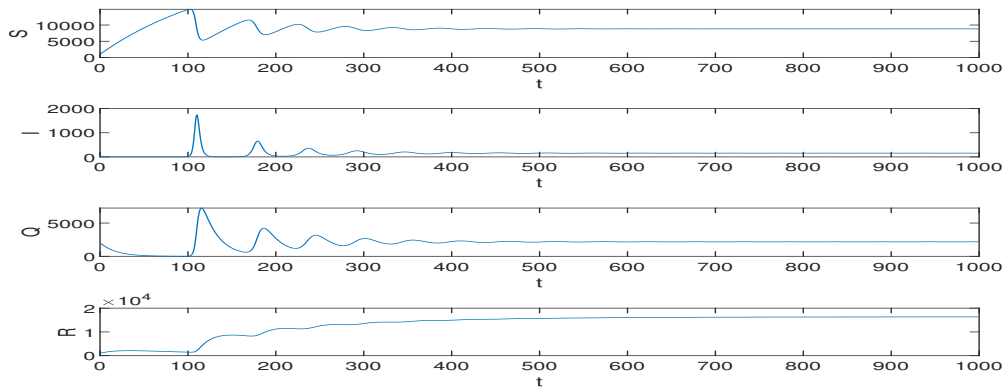


Figure 5. When $\tau = 0$, the equilibrium E_2 of system (2.1) is locally asymptotically stable.

When $\tau = 0.5 \in (0, \tau_{2,3}^{(0)})$, we choose the initial value $(1,000, 50, 2,000, 1,000)$, and the equilibrium E_2 of system (2.1) is locally asymptotically stable. See Figure 6.

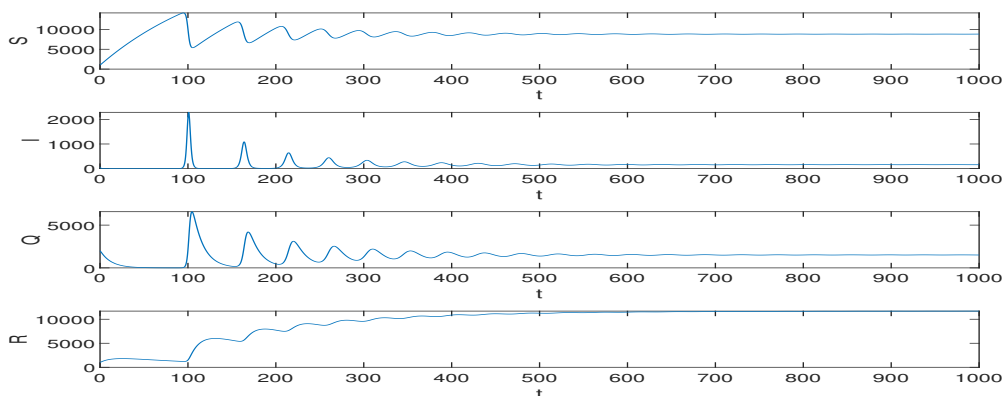


Figure 6. When $\tau = 0.5$, the equilibrium E_2 of system (2.1) is locally asymptotically stable.

When $\tau \in (\tau_{2,3}^{(0)}, +\infty)$, $\tau_{2,3}^{(0)} = 0.7782$, the equilibrium E_2 is unstable. We choose $\tau = 0.8 > \tau_{2,3}^{(0)} = 0.7782$ and the initial value $(1,000, 50, 2,000, 1,000)$. See Figure 7.

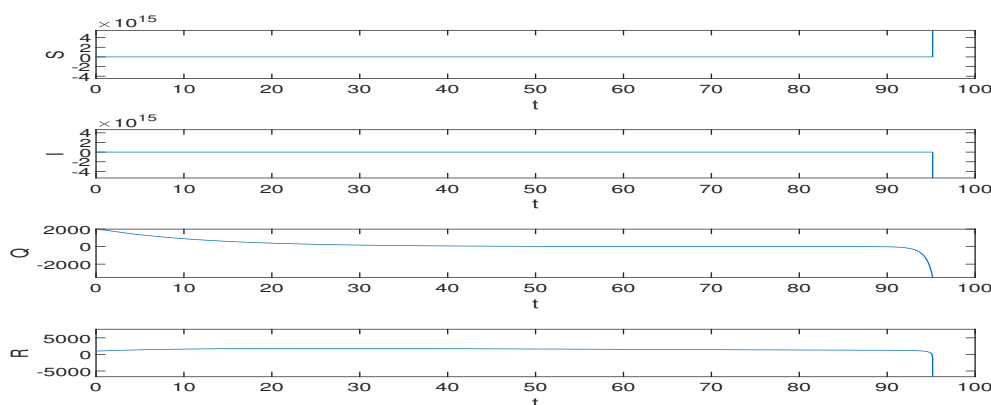


Figure 7. When $\tau = 0.8$, the equilibrium E_2 of system (2.1) is unstable.

According to numerical simulations, we can find that the smaller τ is, the better the control effect is. The results of the numerical simulation are consistent with this fact. Then, we provide some explanations for the above results.

Remark 3: For the two groups of parameter values that are $\Lambda = 100, \mu = 0.00713, \varepsilon = 0.8, \gamma_1 = 0.1, \gamma_2 = 0.1, \alpha_1 = 0.025, \alpha_2 = 0.025, \beta = 0.00001$ and $\Lambda = 200, \mu = 0.00713, \varepsilon = 0.8, \gamma_1 = 0.05, \gamma_2 = 0.05, \alpha_1 = 0.025, \alpha_2 = 0.025, \beta = 0.0001$, we have the following conclusions.

For the first group of parameter values, we can control $\tau < 1.9509$, and the epidemic situation can gradually become stable. However, if $\tau > 1.9509$, the epidemic cannot be controlled into a stable situation. This means that if we can make the time delay from infection to isolation within 1.9509, we can make the number of infected individuals gradually tend to zero, and the epidemic situation will be controllable. However, if we cannot make the time delay from infection to isolation within 1.9509, the epidemic situation will be uncontrollable.

For the second group of parameter values, we can control $\tau < 0.7782$, and the epidemic situation can be gradually stable. If $\tau > 0.7782$, the epidemic cannot be controlled into a stable situation. This means that if we can make the time delay from infection to isolation within 0.7782, we can make the number of infected individuals gradually tend to stabilize, and the epidemic situation will be controllable. However, if we cannot make the time delay from infection to isolation within 0.7782, the epidemic situation will be uncontrollable.

Remark 4: According to our theoretical analysis, with the increase in τ , the epidemic situation will be increasingly severe, and COVID-19 will further threaten the health of humans. To control the epidemic situation, we need to take some measures to shorten the time from infection to isolation. In addition, for epidemic situations in different areas, we can deduce the corresponding parameter values and use the model we constructed to obtain the stability determination and the strategies to control epidemic situations.

Our work is reproducible. For different regions, we can calculate the different critical time delays by changing variables such as the transmission rate from S to I , the transmission rate from I to Q and the transmission rate from I to R . We can obtain the critical time delay based on Eqs. (3.10)-(3.12) and Eqs. (3.15)-(3.17). Then, according to Theorem 3.1, we know that the equilibria are locally asymptotically stable if the time delay is within the critical time delay. Finally, we can obtain the

normal form of Hopf bifurcation by using the multiple time scales method and determine whether the epidemic situation can be controlled into a stable situation based on Theorem 4.1. The time delay from infection to isolation is the most important variable in our model. We calculate the critical time delay based on the parameter values. We choose parameter values according to the natural mortality in China and the data presented in ref. [32]. Then, we change the value of the time delay from infection to isolation, and we know that the equilibria are locally asymptotically stable if the time delay is within the critical time delay through numerical simulations.

6. Conclusion

In this paper, we constructed a DDE according to the characteristics of the COVID-19 epidemic phenomenon based on the *SIQR* model. We considered the existence and stability of equilibria in the above delayed *SIQR* model. We also analyzed the existence and dynamic properties of Hopf bifurcation associated with both equilibria. We chose two groups of parameter values according to the natural mortality in China and the data presented in ref. [32]. Numerical simulations were carried out to verify the analytical results.

According to the results of numerical simulations, the epidemic solution would be more severe with increasing the time delay from infection to isolation. In addition, we predicted the stability of epidemic solutions in different areas and provided effective strategies to control the epidemic.

Acknowledgments

This study was funded by Fundamental Research Funds for the Central Universities of China. (Grant No. 2572019BC14), the Heilongjiang Provincial Natural Science Foundation of China (Grant No. LH2019A001) and College Students Innovations Special Project funded by Northeast Forestry University of China (No. 202010225035).

Conflict of interest

All authors declare no conflicts of interest in this paper.

References

1. J. Liu, W. Xie, Y. Wang, Y. Xiong, S. Chen, J. Han, et al., A comparative overview of COVID-19, MERS and SARS: Review article, *Int. J. Surg.*, **81** (2020), 1–8.
2. C. Huang, Y. Wang, X. Li, L. Ren, J. Zhao, Y. Hu, et al., Clinical features of patients infected with 2019 novel coronavirus in Wuhan, China, *lancet*, **395** (2020), 497–506.
3. C. A. Varotsos, V. F. Krapivin, Y. Xue, Diagnostic model for the society safety under COVID-19 pandemic conditions, *Saf. Sci.*, **136** (2021), 105164.
4. C. A. Varotsos, V. F. Krapivin, A new model for the spread of COVID-19 and the improvement of safety, *Saf. Sci.*, **132** (2020), 104962.

5. U. Avila-Ponce de León, Á. G. C. Pérez, E. Avila-Vales, An SEIARD epidemic model for COVID-19 in Mexico: Mathematical analysis and state-level forecast, *Chaos Solitons Fractals*, **140** (2020), 110165.
6. S. Annas, Muh. Isbar Pratama, Muh. Rifandi, W. Sanusi, S. Side, Stability analysis and numerical simulation of SEIR model for pandemic COVID-19 spread in Indonesia, *Chaos Solitons Fractals*, **139** (2020), 110072.
7. I. Cooper, A. Mondal, C. G. Antonopoulos, A SIR model assumption for the spread of COVID-19 in different communities, *Chaos Solitons Fractals*, **139** (2020), 110057.
8. A. Comunian, R. Gaburro, M. Giudici, Inversion of a SIR-based model: A critical analysis about the application to COVID-19 epidemic, *Physica D*, **413** (2020), 132674.
9. C. Varotsos, J. Christodoulakis, A. Kouremadas, E. F. Fotaki, The Signature of the Coronavirus Lockdown in Air Pollution in Greece, *Water Air Soil Pollut.*, **232** (2021), 119.
10. Z. Liu, P. Magal, O. Seydi, G. Webb, A COVID-19 epidemic model with latency period, *Infect. Dis. Model.*, **5** (2020), 323–337.
11. R. Xu, Global dynamics of an SEIS epidemic model with saturation incidence and latent period, *Appl. Math. Comput.*, **218** (2012), 7927–7938.
12. P. Yang, Y. Wang, Dynamics for an SEIRS epidemic model with time delay on a scale-free network, *Physica A*, **527** (2019), 121290.
13. Y. Wang, J. Cao, G. Q. Sun, J. Li, Effect of time delay on pattern dynamics in a spatial epidemic model, *Physica A*, **412** (2014), 137–148.
14. Q.-M. Liu, C.-S. Deng, M.-C. Sun, The analysis of an epidemic model with time delay on scale-free networks, *Physica A*, **410** (2014), 79–87.
15. H. Lu, Y. Ding, S. Gong, S. Wang, Mathematical modeling and dynamic analysis of SIQR model with delay for pandemic COVID-19, *Math. Biosci. Eng.*, **18** (2021), 3197–3214.
16. D. Greenhalgh, Hopf Bifurcation in Epidemic Models with a Latent Period and Nonpermanent Immunity, *Math. Comput. Model.*, **25** (1997), 85–107.
17. Y. Xie, Z. Wang, J. Lu, Y. Li, Stability analysis and control strategies for a new SIS epidemic model in heterogeneous networks, *Appl. Math. Comput.*, **383** (2020), 125381.
18. F. Wei, R. Xue, Stability and extinction of SEIR epidemic models with generalized nonlinear incidence, *Math. Comput. Simul.*, **170** (2020), 1–15.
19. Y. Yu, L. Ding, L. Lin, P. Hu, X. An, Stability of the SNIS epidemic spreading model with contagious incubation period over heterogeneous networks, *Physica A*, **526** (2019), 120878.
20. X. Wei, G. Xu, W. Zhou, Global stability of endemic equilibrium for a SIQRS epidemic model on complex networks, *Physica A*, **512** (2018), 203–214.
21. Z. Jiang, J. Wei, Stability and bifurcation analysis in a delayed SIR model, *Chaos Solitons Fractals*, **35** (2008), 609–619.
22. Y. A. Kuznetsov, C. Piccardi, Bifurcation analysis of periodic SEIR and SIR epidemic models, *J. Math. Biol.*, **32** (1994), 109–121.

23. H. L. Smith, Subharmonic Bifurcation in an S-I-R Epidemic Model, *J. Math. Biol.*, **17** (1983), 163–177.
24. P. E. Kloeden, C. Pötzsche, Nonautonomous bifurcation scenarios in SIR models, *Math. Meth. Appl. Sci.*, **38** (2015), 3495–3518.
25. J. Li, Z. Teng, G. Wang, L. Zhang, C. Hu, Stability and bifurcation analysis of an SIR epidemic model with logistic growth and saturated treatment, *Chaos Solitons Fractals*, **99** (2017), 63–71.
26. X.-Z. Li, W.-S. Li, M. Ghosh, Stability and bifurcation of an SIR epidemic model with nonlinear incidence and treatment, *Appl. Math. Comput.*, **210** (2009), 141–150.
27. Z. Hu, Z. Teng, L. Zhang, Stability and bifurcation analysis in a discrete SIR epidemic model, *Math. Comput. Simul.*, **97** (2014), 80–93.
28. J. Wang, S. Liu, B. Zheng, Y. Takeuchi, Qualitative and bifurcation analysis using an SIR model with a saturated treatment function, *Math. Comput. Model.*, **55** (2012), 710–722.
29. T. K. Kar, P. K. Mondal, Global dynamics and bifurcation in delayed SIR epidemic model, *Non-linear Anal. Real World Appl.*, **12** (2011), 2058–2068.
30. A. d’Onofrio, P. Manfredi, E. Salinelli, Bifurcation Thresholds in an SIR Model with Information-Dependent Vaccination, *Math. Model. Nat. Phenom.*, **2** (2007), 3495–3518.
31. Y. Ding, W. Jiang, P. Yu, Hopf-zero bifurcation in a generalized Gopalsamy neural network model, *Chaos Solitons Fractals*, **70** (2012), 1037–1050.
32. Y. Wei, Z. Lu, Z. Du, Z. Zhang, Y. Zhao, S. Shen, Fitting and forecasting the trend of COVID-19 by SEIR^{+CAQ} dynamic model, *Chin. J. Epidemiol.*, **41** (2020), 470–475, (Chinese).



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