Research article

Symmetry solutions and conservation laws of a new generalized 2D Bogoyavlensky-Konopelchenko equation of plasma physics

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Abstract: In physics as well as mathematics, nonlinear partial differential equations are known as veritable tools in describing many diverse physical systems, ranging from gravitation, mechanics, fluid dynamics to plasma physics. In consequence, we analytically examine a two-dimensional generalized Bogoyavlensky-Konopelchenko equation in plasma physics in this paper. Firstly, the technique of Lie symmetry analysis of differential equations is used to find its symmetries and perform symmetry reductions to obtain ordinary differential equations which are solved to secure possible analytic solutions of the underlying equation. Then we use Kudryashov’s and \((G'/G)\)-expansion methods to acquire analytic solutions of the equation. As a result, solutions found in the process include exponential, elliptic, algebraic, hyperbolic and trigonometric functions which are highly important due to their various applications in mathematic and theoretical physics. Moreover, the obtained solutions are represented in diagrams. Conclusively, we construct conservation laws of the underlying equation through the use of multiplier approach. We state here that the results secured for the equation under study are new and highly useful.

Keywords: two-dimensional generalized Bogoyavlensky-Konopelchenko equation; Lie point symmetries; analytic solutions; conservation laws

Mathematics Subject Classification: 35L65, 35B06

1. Introduction

Plasma physics simply refers to the study of a state of matter consisting of charged particles. Plasmas are usually created by heating a gas until the electrons become detached from their parent atom or molecule. In addition, plasma can be generated artificially when a neutral gas is heated or subjected to a strong electromagnetic field. The presence of free charged particles makes plasma electrically
conductive with the dynamics of individual particles and macroscopic plasma motion governed by collective electromagnetic fields [1].

Nonlinear partial differential equations (NPDE) in the fields of mathematics and physics play numerous important roles in theoretical sciences. They are the most fundamental models essential in studying nonlinear phenomena. Such phenomena occur in plasma physics, oceanography, aerospace industry, meteorology, nonlinear mechanics, biology, population ecology, fluid mechanics to mention a few. We have seen in [2] that the authors studied a generalized advection-diffusion equation which is a NPDE in fluid mechanics, characterizing the motion of buoyancy propelled plume in a bent-on absorptive medium. Moreover, in [3], a generalized Korteweg-de Vries-Zakharov-Kuznetsov equation was studied. This equation delineates mixtures of warm adiabatic fluid, hot isothermal as well as cold immobile background species applicable in fluid dynamics. Furthermore, the authors in [4] considered a NPDE where they explored important inclined magneto-hydrodynamic flow of an upper-convected Maxwell liquid through a leaky stretched plate. In addition, heat transfer phenomenon was studied with heat generation and absorption effect. The reader can access more examples of NPDEs in [5–16].

In order to really understand these physical phenomena it is of immense importance to solve NPDEs which govern these aforementioned phenomena. However, there is no general systematic theory that can be applied to NPDEs so that their analytic solutions can be obtained. Nevertheless, in recent times scientists have developed effective techniques to obtain viable analytical solutions to NPDEs, such as inverse scattering transform [16], simple equation method [17], Bäcklund transformation [18], F-expansion technique [19], extended simplest equation method [20], Hirota technique [21], Lie symmetry analysis [22–27], bifurcation technique [28, 29], the (G'/G)-expansion method [30], Darboux transformation [31], sine-Gordon equation expansion technique [32], Kudryashov’s method [33], and so on.

The (2+1)-dimensional Bogoyavlensky-Konopelchenko (BK) equation given as

\[
\begin{align*}
 u_{tx} + 6\alpha u_{ux} + 3\beta u_{uxy} + 3\alpha u_{uxx} + \alpha u_{xxxx} + \beta u_{xxx} &= 0, \\
 \end{align*}
\]  

(1.1)

where parameters \( \alpha \) and \( \beta \) are constants, is a special case of the KdV equation in [34] which was introduced as a (2+1)-dimensional version of the KdV and it is described as an interaction of a long wave propagation along \( x \)-axis and a Riemann wave propagation along the \( y \)-axis [35]. In addition to that, few particular properties of the equation have been explored. The authors in [36] provided a Darboux transformation for the BK equation and the obtained transformation was used to construct a family of solutions of this equation. In [37], with \( 3\beta \) replaced by \( 4\beta \) and \( u_y = v_x \) in (1.1), the authors integrated the result once to get

\[
\begin{align*}
 u_t + \alpha u_{xxx} + \beta v_{xxx} + 3\alpha u_x^2 + 4\beta u_x v_x &= 0, \\
 u_y - v_x &= 0.
\end{align*}
\]  

(1.2)

Further, they utilized Lie group theoretic approach to obtain solutions of the system of Eq (1.2). They also engaged the concept of nonlinear self-adjointness of differential equations in conjunction with formal Lagrangian of (1.2) for constructing nonlocal conservation laws of the system. In addition, various applications of BK equation (1.1) were highlighted in [37]. Further investigations on certain particular cases of (1.1) were also carried out in [38, 39].

In [40], the 2D generalized BK equation that reads

\[
\begin{align*}
 u_{tx} + k_1 u_{xxxx} + k_2 u_{xxyy} + \frac{2k_1k_3}{k_2} u_x u_{xx} + k_3 (u_x u_y)_x + \gamma_1 u_{xx} + \gamma_2 u_{xy} + \gamma_3 u_{yy} &= 0
\end{align*}
\]  

(1.3)
was studied and lump-type and lump solutions were constructed by invoking the Hirota bilinear method. Liu et al. [41] applied the Lie group analysis together with \((G'/G)\)-expansion and power series methods and obtained some analytic solutions of (1.3).

Yang et al. [42] recently examined a generalized combined fourth-order soliton equation expressed as

\[
\alpha (6u_xu_{xx} + u_{xxxx}) + \beta [3(u_xu_{x})_x + u_{xxxx}] + \gamma [3(u_xu_y)_x + u_{xxy}] + \delta_1 u_{yt} + \delta_2 u_{xx} + \delta_3 u_{xt} + \delta_4 u_{xy} + \delta_5 u_{yy} + \delta_6 u_{tt} = 0,
\]

with constant parameters \(\alpha, \beta\) and \(\gamma\) which are not all zero, whereas all constant coefficients \(\delta_i, 1 \leq i \leq 6\), are arbitrary. It was observed that Eq (1.4) comprises three fourth-order terms and second-order terms that consequently generalizes the standard Kadomtsev-Petviashvili equation. Soliton equations are known to have applications in plasma physics and other nonlinear sciences such as fluid mechanics, atomic physics, biophysics, nonlinear optics, classical and quantum fields theories.

Assuming \(\alpha = 0, \beta = 1, \gamma = 0\) and \(\delta_1 = \delta_2 = 1, \delta_3 = \delta_4 = \delta_5 = \delta_6 = 0\), the authors gain an integrable \((1+2)\)-dimensional extension of the Hirota-Satsuma equation commonly referred to as the Hirota-Satsuma-Ito equation in two dimensions [43] given as

\[
u_{ty} + u_{xx} + 3(u_xu_x)_x + u_{xxt} = 0
\]

that satisfies the Hirota three-soliton condition and also admits a Hirota bilinear structure under logarithmic transformation presented in the form

\[
u = 2 (\ln f)_x, \text{ where } \left(D_x^3D_t + D_xD_t + D^2_t\right)f \cdot f = 0,
\]

whose lump solutions have been calculated in [44]. On taking parameters \(\alpha = 1, \beta = 0, \gamma = 0\) with \(\delta_1 = \delta_4 = \delta_6 = 0\) whereas \(\delta_2 = \delta_3 = \delta_5 = 1\), they eventually came up with a two dimensional equation [42]:

\[
u_{tx} + 6u_xu_{xx} + u_{xxxx} + 3\left(u_xu_y\right)_x + u_{xx} + u_{yy} = 0,
\]

which is called a two-dimensional generalized Bogoyavlensky-Konopelchenko (2D-gBK) equation. We notice that if one takes \(\alpha = \beta = 1\) in Eq (1.1) with the introduction of two new terms \(u_{xx}\) and \(u_{yy}\), the new generalized version (1.7) is achieved.

In consequence, we investigate explicit solutions of the new two-dimensional generalized Bogoyavlensky-Konopelchenko equation (1.7) of plasma physics in this study. In order to achieve that, we present the paper in the subsequent format. In Section 2, we employ Lie symmetry analysis to carry out the symmetry reductions of the equation. In addition, direct integration method will be employed in order to gain some analytic solutions of the equation by solving the resulting ordinary differential equations (ODEs) from the reduction process. We achieve more analytic solutions of (1.7) via the conventional \((G'/G)\)-expansion method as well as Kudryashov’s technique. In addition, by choosing suitable parametric values, we depict the dynamics of the solutions via 3-D, 2-D as well as contour plots. Section 3 presents the conservation laws for 2D-gBK equation (1.7) through the multiplier method and in Section 4, we give the concluding remarks.

2. Symmetry analysis and analytic solutions of (1.7)

In this section we in the first place compute the Lie point symmetries of Eq (1.7) and thereafter engage them to generate analytic solutions.
2.1. Lie point symmetries of (1.7)

A one-parameter Lie group of symmetry transformations associated with the infinitesimal generators related to (gbk) can be presented as

\[ \begin{align*}
\bar{t} &= t + \epsilon \xi^1(t, x, y, u) + O(\epsilon^2), \\
\bar{x} &= x + \epsilon \xi^2(t, x, y, u) + O(\epsilon^2), \\
\bar{y} &= y + \epsilon \xi^3(t, x, y, u) + O(\epsilon^2), \\
\bar{u} &= u + \epsilon \phi(t, x, y, u) + O(\epsilon^2).
\end{align*} \]

(2.1)

We calculate symmetry group of 2D-gBK equation (1.7) using the vector field

\[ R = \xi^1(t, x, y, u) \frac{\partial}{\partial t} + \xi^2(t, x, y, u) \frac{\partial}{\partial x} + \xi^3(t, x, y, u) \frac{\partial}{\partial y} + \phi(t, x, y, u) \frac{\partial}{\partial u}, \]

(2.2)

where \( \xi^i, i = 1, 2, 3 \) and \( \phi \) are functions depending on \( t, x, y \) and \( u \). We recall that (2.2) is a Lie point symmetry of Eq. (1.7) if

\[ R^{(4)}(u_{tx} + 6u_xu_{xx} + u_{xxxx} + 3(u_xu_y)_x + u_{xx} + u_{yy})|_{\epsilon=0} = 0, \]

(2.3)

where \( Q = u_{tx} + 6u_xu_{xx} + u_{xxxx} + 3(u_xu_y)_x + u_{xx} + u_{yy} \). Here, \( R^{(4)} \) denotes the fourth prolongation of \( R \) defined by

\[ R^{(4)} = R + \eta^t \partial_{u_t} + \eta^x \partial_{u_x} + \eta^y \partial_{u_y} + \eta^x \partial_{u_{tx}} + \eta^y \partial_{u_{ty}} + \eta^{xy} \partial_{u_{txy}} + \eta^{ttx} \partial_{u_{txx}}, \]

(2.4)

where coefficient functions \( \eta^t, \eta^x, \eta^y, \eta^t^x, \eta^t^y, \eta^{xy}, \eta^{ttx} \) and \( \eta^{ttxy} \) can be calculated from [22–24].

Writing out the expanded form of the determining equation (2.3), splitting over various derivatives of \( u \) and with the help of Mathematica, we achieve the system of linear partial differential equations (PDEs):

\[
\begin{align*}
\xi_1^2 &= 0, \quad \xi_2^2 = 0, \quad \xi_3^2 = 0, \quad \xi_4^2 = 0, \\
\xi_1^3 + 5\xi_2^3 = 0, \quad \xi_2^3 = 0, \quad 5\xi_3^3 = 0, \\
5\xi_4^3 - 2\xi_5^3 = 0, \quad 5\xi_5^3 - 3\xi_6^3 = 0, \quad 3\phi_{xx} - 3\xi_6^3 = 0, \quad \xi_7^3 - 3\phi_x + 2\xi_8^3 = 0, \\
4\xi_9^3 - 3\phi_{xx} + 15\phi_y = 0, \quad \phi_{tx} + \phi_{xx} + \phi_{xxxx} + \phi_{xyy} + \phi_{yy} = 0.
\end{align*}
\]

The solution of the above system of PDEs is

\[
\begin{align*}
\xi_1^1 &= A_1 + A_2t, \quad \xi_2^2 = F(t) + \frac{1}{5}A_2(x + 2y), \\
\xi_3^3 &= A_4 - \frac{4}{5}A_2t + 3A_3t + \frac{3}{5}A_2y, \\
\eta &= G(t) - \frac{1}{5}A_2u + A_3x - \frac{4}{15}A_2y - 2A_3y + \frac{1}{3}yF'(t),
\end{align*}
\]

where \( A_1 \)–\( A_4 \) are arbitrary constants and \( F(t), G(t) \) are arbitrary functions of \( t \). Consequently, we secure the Lie point symmetries of (1.7) given as

\[
\begin{align*}
R_1 &= \frac{\partial}{\partial t}, \quad R_2 = \frac{\partial}{\partial y}, \quad R_3 = 3F(t) \frac{\partial}{\partial x} + yF'(t) \frac{\partial}{\partial u}, \quad R_4 = 3t \frac{\partial}{\partial y} + (x - 2y) \frac{\partial}{\partial u}, \\
R_5 &= G(t) \frac{\partial}{\partial u}, \quad R_6 = 15t \frac{\partial}{\partial t} + (3x + 6y) \frac{\partial}{\partial x} + (9y - 12t) \frac{\partial}{\partial y} - (4y + 3u) \frac{\partial}{\partial u}.
\end{align*}
\]

(2.5)
2.2. Lie group transformations associated to (2.5)

We contemplate the exponentiation of the vector fields (2.5) by computing the flow or one parameter group generated by (2.5) via the Lie equations [22, 23]:

\[ \frac{d\tilde{t}}{de} = \xi^1(\tilde{t}, \tilde{x}, \tilde{y}, \tilde{u}), \quad \tilde{t}|_{e=0} = t, \]
\[ \frac{d\tilde{x}}{de} = \xi^2(\tilde{t}, \tilde{x}, \tilde{y}, \tilde{u}), \quad \tilde{x}|_{e=0} = x, \]
\[ \frac{d\tilde{y}}{de} = \xi^3(\tilde{t}, \tilde{x}, \tilde{y}, \tilde{u}), \quad \tilde{y}|_{e=0} = y, \]
\[ \frac{d\tilde{u}}{de} = \phi(\tilde{t}, \tilde{x}, \tilde{y}, \tilde{u}), \quad \tilde{u}|_{e=0} = u. \]

Therefore, by taking \( F(t) = G(t) = t \) in (2.5), one computes a one parameter transformation group of 2D-gBK (1.7). Thus, we present the result in the subsequent theorem.

**Theorem 2.1.** Let \( T^i_\epsilon(t, x, y, u), i = 1, 2, 3, \ldots, 6 \) be transformations group of one parameter generated by vectors \( R_1, R_2, R_3, \ldots, R_6 \) in (2.5), then, for each of the vectors, we have accordingly

\[ T^1_\epsilon : (\tilde{t}, \tilde{x}, \tilde{y}, \tilde{u}) \longrightarrow (t + \epsilon_1, x, y, u), \]
\[ T^2_\epsilon : (\tilde{t}, \tilde{x}, \tilde{y}, \tilde{u}) \longrightarrow (t, x + \epsilon_2, y, u), \]
\[ T^3_\epsilon : (\tilde{t}, \tilde{x}, \tilde{y}, \tilde{u}) \longrightarrow (t, 3\epsilon_3t + x, y + \epsilon_3y + u), \]
\[ T^4_\epsilon : (\tilde{t}, \tilde{x}, \tilde{y}, \tilde{u}) \longrightarrow (t, x, 3\epsilon_4t + y + u) + (x - 2y)\epsilon_4 - 3\epsilon_4^2t), \]
\[ T^5_\epsilon : (\tilde{t}, \tilde{x}, \tilde{y}, \tilde{u}) \longrightarrow (t, x, \epsilon_5t + u), \]
\[ T^6_\epsilon : (\tilde{t}, \tilde{x}, \tilde{y}, \tilde{u}) \longrightarrow \bigg( t^e_{156} + (2e_9^{6} - e_3^{6} - e_5^{156})t + xe_3^{6} + (e_9^{6} - e_3^{6})y, \bigg( 2e_9^{6} - 2e_5^{156})t + ye_9^{6}, \frac{1}{9} \left[ (4e_9^{186} - 6e_9^{126} + 2t) + (3 - 3e_9^{126})y + 9u \right] e_3^{6} \bigg), \]

where \( \epsilon \in \mathbb{R} \) is regarded as the group parameter.

**Theorem 2.2.** Hence, suppose \( u(t, x, y) = \Theta(t, x, y) \) satisfies the 2D-gBK (1.7), in the same vein, the functions given in the structure

\[ u^1(t, x, y) = \Theta(t - \epsilon_1, x, y, z), \]
\[ u^2(t, x, y) = \Theta(t, x - \epsilon_2, y, u), \]
\[ u^3(t, x, y) = \Theta(t, x - 3\epsilon_3t, y) - \epsilon_3y, \]
\[ u^4(t, x, y) = \Theta(t, x, 3\epsilon_4t + y) - (x - 2y)\epsilon_4 + 3\epsilon_4^2t, \]
\[ u^5(t, x, y) = \Theta(t, x, \epsilon_5t) - \epsilon_5t, \]
\[ u^6(t, x, y) = \Theta \left[ t^e_{156} + (2e_9^{6} - e_3^{6} - e_5^{156})t + xe_3^{6} + (e_9^{6} - e_3^{6})y, \bigg( 2e_9^{6} - 2e_5^{156})t + ye_9^{6}, \frac{1}{9} \left[ (4e_9^{186} - 6e_9^{126} + 2t) + (3 - 3e_9^{126})y + 9u \right] e_3^{6} \bigg), \]

will do, where \( u^i(t, x, y) = T^i_\epsilon \cdot \Theta(t, x, y), i = 1, 2, 3, \ldots, 6 \) with \( \epsilon << 1 \) regarded as any positive real number.
2.3. Symmetry reduction of 2D-gBK equation (1.7)

In this subsection, we utilize symmetries (2.5) with a view to reduce Eq (1.7) to ordinary differential equations and thereafter obtain the analytic solutions of Eq (1.7) by solving the respective ODEs.

**Case 1. Invariant solutions via \(R_1-R_3\)**

Taking \(F(t) = 1/3\), we linearly combine translational symmetries \(R_1-R_3\) as \(R = bR_1 + cR_2 + aR_3\) with nonzero constant parameters \(a, b\) and \(c\). Subsequently utilizing the combination reduces 2D-gBK equation (1.7) to a PDE with two independent variables. Thus, solution to the characteristic equation associated with the symmetry \(R\) leaves us with invariants

\[
r = ct - ay, \quad s = cx - by, \quad \theta = u.
\]

(2.6)

Now treating \(\theta\) above as the new dependent variable as well as \(r, s\) as new independent variables, (1.7) then transforms into the PDE:

\[
c^2 \theta_{rs} + 6c^3 \theta_s \theta_{ss} - 3ac^2 \theta_s \theta_{sr} - 6bc^2 \theta_s \theta_{ss} - 3ac^2 \theta_s \theta_{ss} - ac^3 \theta_{sssr} - bc^3 \theta_{ssss} + c^2 \theta_{ss} + a^2 \theta_{rr} + 2ab \theta_{sr} + b^2 \theta_{ss} = 0.
\]

(2.7)

We now utilize the Lie point symmetries of (2.7) in a bid to transform it to an ODE. From (2.7), we achieve three translation symmetries:

\[
Q_1 = \frac{\partial}{\partial r}, \quad Q_2 = \frac{\partial}{\partial s}, \quad Q_3 = \frac{\partial}{\partial \theta}.
\]

The linear combination \(Q = Q_1 + \omega Q_2\) (\(\omega \neq 0\) being an arbitrary constant) leads to two invariants:

\[
z = s - \omega r, \quad \theta = \Theta,
\]

(2.8)

that secures group-invariant solution \(\Theta = \Theta(z)\). Thus, on using these invariants, (2.7) is transformed into the fourth-order nonlinear ODE:

\[
(c^2 - \omega c^2 + a^2 \omega^2 - 2b \omega + b^2) \Theta''(z) - 6(\beta bc^2 - \beta c^2 a \omega - c^3) \Theta'(z) \Theta''(z)
\]

\[
+ (c^3 a \omega + c^4 - bc^3) \Theta'''(z) = 0,
\]

which we rewrite in a simple structure as

\[
A \Theta''(z) - B \Theta'(z) \Theta''(z) + C \Theta'''(z) = 0,
\]

(2.9)

where \(A = c^2 - \omega c^2 + a^2 \omega^2 - 2b \omega + b^2\), \(B = 6(bc^2 - c^2 a \omega - c^3)\), \(C = c^3 a \omega + c^4 - bc^3\) and \(z = cx + (aw - b)y - cot\).

2.4. Some analytic solutions of 2D-gBK equation (1.7)

In this section, we seek travelling wave solutions of the 2D-gBK equation (1.7).

**A. Elliptic function solution of (1.7)**

On integrating equation (2.9) once, we accomplish a third-order ODE:

\[
A \Theta'(z) - \frac{1}{2} B \Theta''(z) + C \Theta'''(z) + C_1 = 0,
\]

(2.10)
where $C_1$ is a constant of integration. Multiplying Eq (2.10) by $\Theta''(z)$, integrating once and simplifying the resulting equation, we have the second-order nonlinear ODE:

$$
\frac{1}{2} A \Theta'(z)^2 - \frac{1}{6} B \Theta'(z)^3 + \frac{1}{2} C \Theta''(z)^2 + C_1 \Theta'(z) + C_2 = 0,
$$

where $C_2$ is a constant of integration. The above equation can be rewritten as

$$
\Theta''(z)^2 = \frac{B}{3C} \Theta'(z)^3 - \frac{A}{C} \Theta'(z)^2 - \frac{2C_1}{C} \Theta'(z) - \frac{2C_2}{C}.
$$

(2.11)

Letting $U(z) = \Theta'(z)$, Eq (2.11) becomes

$$
U'(z)^2 = \frac{B}{3C} U(z)^3 - \frac{A}{C} U(z)^2 - \frac{2C_1}{C} U(z) - \frac{2C_2}{C}.
$$

(2.12)

Suppose that the cubic equation

$$
U(z)^3 - \frac{3A}{B} U(z)^2 - \frac{6C_1}{B} U(z) - \frac{6C_2}{B} = 0
$$

(2.13)

has real roots $c_1 - c_3$ such that $c_1 > c_2 > c_3$, then Eq (2.12) can be written as

$$
U'(z)^2 = \frac{B}{3C} (U(z) - c_1)(U(z) - c_2)(U(z) - c_3),
$$

(2.14)

whose solution with regards to Jacobi elliptic function [45, 46] is

$$
U(z) = c_2 + (c_1 - c_2) \text{cn}^2 \left\{ \frac{B(c_1 - c_2)}{12C}, \Delta^2 \right\}, \Delta^2 = \frac{c_1 - c_2}{c_1 - c_3},
$$

(2.15)

with (cn) being the elliptic cosine function. Integration of (2.15) and reverting to the original variables secures a solution of 2D-gBK equation (1.7) as

$$
u(t, x, y) = \sqrt{\frac{12C(c_1 - c_2)^2}{B(c_1 - c_3)\Delta^8}} \left\{ \text{EllipticE} \left[ \text{sn} \left( \frac{B(c_1 - c_3)}{12C}, \Delta^2 \right), \Delta^2 \right] \right\}

+ \left\{ c_2 - (c_1 - c_2) \frac{1 - \Delta^4}{\Delta^4} \right\} z + C_3,
$$

(2.16)

with $z = cx + (a\omega - b)y - c\omega t$ and $C_3$ a constant of integration. We note that (2.16) is a general solution of (1.7), where EllipticE[$p; q$] is the incomplete elliptic integral [46, 47] expressed as

$$
\text{EllipticE}[p; q] = \int_0^p \sqrt{1 - q^2 r^2} dr.
$$

We present wave profile of periodic solution (2.16) in Figure 1 with 3D, contour and 2D plots with parametric values $a = -4$, $b = 0.2$, $c = -0.1$, $\omega = 0.1$, $c_1 = 100$, $c_2 = 50.05$, $c_3 = -60$, $B = 10$, $C = 70$, where $t = 1$ and $-10 \leq x, y \leq 10$. 
However, contemplating a special case of (2.9) with $B = 0$, we integrate the equation twice and so we have

$$C \Theta''(z) + A \Theta(z) + K_1 z + K_2 = 0,$$  \hspace{1cm} (2.17)

where $K_1$ and $K_2$ are integration constants. Solving the second-order linear ODE (2.17) and reverting to the basic variables, we achieve the trigonometric solution of 2D-gBK equation (1.7) as

$$u(t, x, y) = A_1 \sin \left( \frac{\sqrt{a^2 \omega^2 - \omega(2ab + c^2)} + b^2 + c^2 \ z}{\sqrt{c^3(a \omega - b + c)}} \right) + A_2 \cos \left( \frac{\sqrt{a^2 \omega^2 - \omega(2ab + c^2)} + b^2 + c^2 \ z}{\sqrt{c^3(a \omega - b + c)}} \right) - \frac{K_1 z + K_2}{a^2 \omega^2 - \omega(2ab + c^2) + b^2 + c^2},$$

(2.18)

with $A_1$ and $A_2$ as the integration constants as well as $z = cx + (a \omega - b)y - c \omega t$. We depict the wave dynamics of periodic solution (2.18) in Figure 2 via 3D, contour and 2D plots with dissimilar parametric values $a = 1, b = 0.2, c = -0.1, \omega = 0.1, A_1 = 20, A_2 = -2, K_1 = 1, K_2 = 10$, where $t = 2$ and $-10 \leq x, y \leq 10$.  

**Figure 2.** Wave profile of the trigonometric function solution (2.18) at $t = 2$. 

**B. Weierstrass elliptic solution of 2D-gBK equation (1.7)**

We further explore Weierstrass elliptic function solution of (1.7), which is a technique often involved in getting general exact solutions to NPDEs [47, 48]. In order to accomplish this, we use
the transformation

\[ U(z) = W(z) + \frac{A}{B} \]  

(2.19)

and transform the nonlinear ordinary differential equation (NODE) (2.12) to

\[ W_\xi^2 = 4W^3 - g_2 W - g_3, \quad \xi = \sqrt{\frac{B}{12C}} z, \]  

(2.20)

with the invariants \( g_2 \) and \( g_3 \) given by

\[ g_2 = -\frac{12A^2}{B^2} - \frac{24C_1}{B} \quad \text{and} \quad g_3 = -\frac{8A^3}{B^3} - \frac{24AC_1}{B^2} - \frac{24C_2}{B}. \]

Thus, we have the solution of NODE (2.12) as

\[ U(z) = \frac{A}{B} + w \left( \sqrt{\frac{1}{12C}} (z - z_0); g_2, g_3 \right), \]  

(2.21)

where \( w \) denotes the Weierstrass elliptic function [46]. In consequence, integration of (2.21) and reverting to the basic variables gives the solution of 2D-gBK equation (1.7) as

\[ u(t, x, y) = \frac{A}{B} (z - z_0) - \sqrt{\frac{12B}{C} \zeta} \left[ \sqrt{\frac{B}{12C}} (z - z_0); g_2, g_3 \right], \]  

(2.22)

with arbitrary constant \( z_0, z = cx + (a\omega - b)y - c\omega t \) and \( \zeta \) being the Weierstrass zeta function [46]. We give wave profile of Weierstrass function solution (2.22) in Figure 3 with 3D, contour and 2D plots using parameter values \( a = 1, b = 0.2, c = -0.1, \omega = 0.1, A = 10, B = -2, z_0 = 0, C = 1, C_1 = 1, C_2 = 10, \) where \( t = 2 \) and \(-10 \leq x, y \leq 10\).

2.4.1. Solution of (1.7) by Kudrayshov’s approach

This part of the study furnishes the solution of 2D-gBK equation (1.7) through the use of Kudrayshov’s approach [33]. This technique is one of the most prominent way to obtain closed-form solutions of NPDEs. Having reduced Eq (1.7) to the NODE (2.9), we assume the solution of (2.9) as

\[ \Theta(z) = \sum_{n=0}^{N} B_n Q^n(z), \]  

(2.23)
with $Q(z)$ satisfying the first-order NODE
\[ Q'(z) = Q^2(z) - Q(z). \] (2.24)

We observe that the solution of (2.24) is
\[ Q(z) = \frac{1}{1 + \exp(z)}. \] (2.25)

The balancing procedure for NODE (2.9) produces $N = 1$. Hence, from (2.23), we have
\[ \Theta(z) = B_0 + B_1 Q(z). \] (2.26)

Now substituting (2.26) into (2.9) and using (2.24), we gain a long determining equation and splitting on powers of $Q(z)$, we get algebraic equations for the coefficients $B_0$ and $B_1$ as
\[
\begin{align*}
Q(z)^5 &: 2aB_1 c^3 \omega + aB_1^2 c^2 \omega - 2bB_1 c^3 - bB_1^2 c^2 + 2B_1 c^4 + B_1^2 c^3 = 0, \\
Q(z)^4 &: 2bB_1 c^3 - 2abB_1 c^3 \omega - aB_1^2 c^2 \omega + bB_1^2 c^2 - 2B_1 c^4 - B_1^2 c^3 = 0, \\
Q(z)^3 &: a^2 B_1 \omega^2 - 2abB_1 \omega + 25aB_1 c^3 \omega + 12aB_1^2 c^2 \omega + b^2 B_1 - 25bB_1 c^3 \\
& \quad - 12bB_1^2 c^2 + 25B_1 c^4 + 12B_1^2 c^3 - B_1 c^2 \omega + B_1 c^2 = 0, \\
Q(z)^2 &: 2abB_1 \omega - a^2 B_1 \omega^2 - 5aB_1 c^3 \omega - 2aB_1^2 c^2 \omega - b^2 B_1 + 5bB_1 c^3 \\
& \quad + 2bB_1^2 c^2 - 5B_1 c^3 - 2B_1^2 c^3 + B_1 c^2 \omega - B_1 c^2 = 0, \\
Q(z) &: a^2 B_1 \omega^2 - 2abB_1 \omega + aB_1 c^3 \omega + b^2 B_1 - bB_1 c^3 + B_1 c^4 - B_1 c^2 \omega \\
& \quad + B_1 c^2 = 0.
\end{align*}
\] (2.27)

The solution of the above system gives
\[ B_0 = 0, \quad B_1 = -2c, \quad a = \frac{2b\omega - c^3 \omega \pm \sqrt{c^2 \omega^2 (c^4 - 4c^2 + 4\omega - 4)}}{2\omega^2}. \] (2.28)

Hence, the solution of 2D-gBK equation (1.7) associated with (2.28) is given as
\[ u(t, x, y) = -\frac{2c}{1 + \exp\{cx + (a\omega - b)y - cw\}}. \] (2.29)

The wave profile of solution (2.29) is shown in Figure 4 with 3D, contour and 2D plots using parameter values $a = 1, b = -0.2, c = 20, \omega = 0.05, B_0 = 0$ with $t = 7$ and $-6 \leq x, y \leq 6$.

\begin{figure}[h]
\centering
\includegraphics[width=\textwidth]{figure4.png}
\caption{The wave profile of solution (2.29) at $t = 7$.}
\end{figure}
2.4.2. Solution of (1.7) through \((G'/G)\)-expansion technique

We reckon the \((G'/G)\)-expansion technique [30] in the construction of analytic solutions of 2D-gBK equation (1.7) and so we contemplate a solution structured as

\[
\Theta(z) = \sum_{j=0}^{M} B_j \left( \frac{Q'(z)}{Q(z)} \right)^j,
\]

(2.30)

where \(Q(z)\) satisfies

\[
Q''(z) + \lambda Q'(z) + \mu Q(z) = 0
\]

(2.31)

with \(\lambda\) and \(\mu\) taken as constants. Here, \(B_0, \ldots, B_M\) are parameters to be determined. Utilization of balancing procedure for (2.9) produces \(M = 1\) and as a result, the solution of (1.7) assumes the form

\[
\Theta(z) = B_0 + B_1 \left( \frac{Q'(z)}{Q(z)} \right).
\]

(2.32)

Substituting the value of \(\Theta(z)\) from (2.32) into (2.9) and using (2.31) and following the steps earlier adopted, leads to an algebraic equation in \(B_0\) and \(B_1\), which splits over various powers of \(Q(z)\) to give the system of algebraic equations whose solution is secured as

\[
B_0 = 0,\; B_1 = 2c,\; \alpha = \frac{16b\omega - B_1^1 \lambda^2 \omega \pm \sqrt{\Omega_0 + 64B_1^2 \lambda^2 \omega^2 - 64B_1^2 \lambda^2 \omega^2 + 4B_1^3 \mu \omega}}{16\omega^2},
\]

where \(\Omega_0 = B_1^1 \lambda^4 \omega^2 - 8B_1^1 \lambda^2 \mu \omega^2 - 16B_1^1 \lambda^2 \omega^2 + 16B_1^2 \mu \omega^2 + 64B_1^4 \mu \omega^2\). Thus, we have three types of solutions of the 2D-gBK equation (1.7) given as follows:

When \(\lambda^2 - 4\mu > 0\), we gain the hyperbolic function solution

\[
u(t, x, y) = B_0 + B_1 \left( \Delta_1 \frac{A_1 \sinh \left( \Delta_1 z \right) + A_2 \cosh \left( \Delta_1 z \right)}{A_1 \cosh \left( \Delta_1 z \right) + A_2 \sinh \left( \Delta_1 z \right)} - \frac{\lambda}{2} \right),
\]

(2.33)

with \(z = cx + (a\omega - b)y - c\omega t, \Delta_1 = \frac{1}{2} \sqrt{\lambda^2 - 4\mu}\) together with \(A_1, A_2\) being arbitrary constants. The wave profile of solution (2.33) is shown in Figure 5 with 3D, contour and 2D plots using parameter values \(a = 3, b = 0.5, c = 10, \omega = -0.1, B_0 = 0, \lambda = -0.971, \mu = 10, A_1 = 5, A_2 = 1\), where \(t = 10\) and \(-10 \leq x, y \leq 10\).

Figure 5. The wave profile of solution (2.33) at \(t = 10\).
When $\lambda^2 - 4\mu < 0$, we achieve the trigonometric function solution

$$u(t, x, y) = B_0 + B_1 \left( \frac{A_2 \cos (\Delta z) - A_1 \sin (\Delta z)}{A_1 \cos (\Delta z) + A_2 \sin (\Delta z)} - \frac{\lambda}{2} \right),$$

(2.34)

with $z = cx + (a\omega - b)y - c\omega t$, $\Delta_2 = \frac{1}{2} \sqrt{4\mu - \lambda^2}$ together with $A_1$ and $A_2$ are arbitrary constants. The wave profile of solution (2.34) is shown in Figure 6 with 3D, contour and 2D plots using parameter values $a = 1$, $b = 0.5$, $c = 0.3$, $\omega = 0.3$, $B_0 = 0$, $\lambda = -0.971$, $\mu = 2$, $A_1 = 5$, $A_2 = 1$ with $t = 10$ and $-10 \leq x, y \leq 10$.

![Figure 6](image1)

**Figure 6.** The wave profile of solution (2.34) at $t = 10$.

When $\lambda^2 - 4\mu = 0$, we gain the rational function solution

$$u(t, x, y) = B_0 + B_1 \left( \frac{A_2}{A_1 + A_2 z} - \frac{\lambda}{2} \right),$$

(2.35)

with $z = cx + (a\omega - b)y - c\omega t$ and $A_1$, $A_2$ being arbitrary constants. We plot the graph of solution (2.35) in Figure 7 via 3D, contour and 2D plots using parametric values $a = 1$, $b = 1.01$, $c = 100$, $\omega = 0.1$, $B_0 = 10$, $\lambda = 10$, $A_1 = 3$, $A_2 = 10$, where $t = 2.4$ and $-5 \leq x, y \leq 5$.

![Figure 7](image2)

**Figure 7.** The wave profile of solution (2.35) at $t = 2.4$. 
Case 2. Group-invariant solutions via $R_4$

Lagrange system associated with the symmetry $R_4 = 3t \partial / \partial y + (x - 2y) \partial / \partial u$ is

\[
\frac{dt}{0} = \frac{dx}{0} = \frac{dy}{3t} = \frac{du}{(x - 2y)},
\]

(2.36)

which leads to the three invariants $T = t$, $X = x$, $Q = u + (y^2/3t) - (xy/3t)$. Using these three invariants, the 2D-gBK equation (1.7) is reduced to

\[
18T Q_X Q_{XX} + 3T Q_T X + 3X Q_{XX} + 3Q_X + 3T Q_{XXX} - 2 = 0.
\]

(2.37)

Case 3. Group-invariant solutions via $R_1$, $R_2$ and $R_5$

We take $G(t) = 1$ and by combining the generators $R_1$, $R_2$ as well as $R_5$, we solve the characteristic equations corresponding to the combination and get the invariants $X = x$, $Y = y - t$ with group-invariant $u = Q(X, Y) + t$. With these invariants, the 2D-gBK equation (1.7) transforms to the NPDE

\[
Q_{XX} + Q_{YY} - Q_{XY} + 3Q_X Q_{XY} + 3Q_Y Q_{XX} + 6Q_X Q_{XX} + Q_{XXX} + Q_{XXY} = 0,
\]

(2.38)

whose solution is given by

\[
Q(X, Y) = 2A_2 \tanh \left[ A_2 X + A_2 \left( \frac{1}{2} - \frac{1}{2} \sqrt{16A_2^4 - 24A_2^2 - 3 - 2A_2^2} \right) Y + A_1 \right] + A_3,
\]

(2.39)

with arbitrary constants $A_1$–$A_3$. Thus, we achieve the hyperbolic solution of (1.7) as

\[
u(t, x, y) = t + 2A_2 \tanh \left[ \frac{1}{2} A_2 (t - y) \sqrt{16A_2^4 - 24A_2^2 - 3} + \frac{1}{2} (4t - 4y)A_2^3 \right.

\]

\[
+ \frac{1}{2}(y + 2x - t)A_2 + A_1 \right] + A_3.
\]

(2.40)

The wave profile of solution (2.40) is shown in Figure 8 with 3D, contour and 2D plots using parameter values $A_1 = 70.1$, $A_2 = -30$, $A_3 = 0$, where $t = 0.5$ and $-10 \leq x, y \leq 10$.

Figure 8. The wave profile of solution (2.40) at $t = 0.5$. 

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Besides, symmetries of (2.38) are found as
\[ P_1 = \frac{\partial}{\partial X}, \quad P_2 = \frac{\partial}{\partial Y}, \quad P_3 = \frac{\partial}{\partial Q}, \quad P_4 = \left( \frac{1}{3} X + \frac{2}{3} Y \right) \frac{\partial}{\partial X} + Y \frac{\partial}{\partial Y} + \left( \frac{2}{3} X - 2 Y - \frac{1}{3} Q \right) \frac{\partial}{\partial Q}. \]

Now, the symmetry \( P_1 \) furnishes the solution \( Q(X, Y) = f(z), \, z = Y \). So, Eq (2.38) gives the ODE \( f''(z) = 0 \). Hence, we have a solution of (1.7) as
\[ u(t, x, y) = t + A_0(y - t) + A_1, \quad (2.41) \]
with \( A_0, A_1 \) as constants. Further, the symmetry \( P_2 \) yields \( Q(X, Y) = f(z), \, z = X \) and so Eq (2.38) reduces to
\[ f'''(z) + 6f'(z)f''(z) + f'''(z) = 0. \quad (2.42) \]
Integration of the above equation three times with respect to \( z \) gives
\[ f'(z)^2 + 2f(z)^3 + f(z)^2 + 2A_0f(z) + 2A_1 = 0, \quad (2.43) \]
and taking constants \( A_0 = A_1 = 0 \) and then integrating it results in the solution of (1.7) as
\[ u(t, x, y) = t - \frac{1}{2} \left\{ 1 + \tan \left( \frac{1}{2} A_1 - \frac{1}{2} x \right) \right\}. \quad (2.44) \]
The wave profile of solution (2.44) is shown in Figure 9 with 3D, contour and 2D plots using parameter values \( A_1 = 40, \, t = 3.5 \) and \(-10 \leq x \leq 10\).

![Figure 9. The wave profile of solution (2.44) at \( t = 3.5 \).](image-url)

On combining \( P_1 - P_3 \) as \( P = c_0 P_1 + c_1 P_2 + c_2 P_3 \), we accomplish
\[ Q(X, Y) = \frac{c_2}{c_0} X + f(z), \quad \text{where} \quad z = c_0 Y - c_1 X. \quad (2.45) \]
Using the newly acquired invariants (2.45), Eq (2.38) transforms to the NODE:
and Eq (2.38) reduces to the NODE

\begin{equation}
\begin{aligned}
c_0c_1^2f''(z) + 6c_1^2c_2f'''(z) - 3c_0c_1c_2f''(z) + c_0^3c_1^2f''(z) + c_0^3f'''(z) + 6c_1^2c_1f'(z)f''(z) \\
- 6c_0c_1^3f'(z)f''(z) + c_0c_1^4f''''(z) - c_0^3c_1f''''(z) = 0.
\end{aligned}
\end{equation}

Engaging the Lie point symmetry \( P_4 \), we obtain

\[ Q(X, Y) = X - 2Y + Y^{-1/3} f(z) \text{ with } z = Y^{-1/3}(X - Y), \quad (2.47) \]

and Eq (2.38) reduces to the NODE

\[ 6zf''(z) + z^2f'''(z) - 18f'(z)^2 - 9f(z)f''(z) + 4f(z) - 18z^2f''(z) - 12f'''(z) - 3zf''''(z) = 0. \quad (2.48) \]

**Case 4.** Group-invariant solutions via \( R_6 \)

Lie point symmetry \( R_6 \) dissociates to the Lagrange system

\[ \frac{dt}{15t} = \frac{dx}{3x + 6y} = \frac{dy}{9y - 12t} = \frac{du}{-(4y + 3u)}, \]

which gives

\[ u = t^{-1/5} Q(T, X) - \frac{2}{9} t - \frac{1}{3} y, \text{ with } T = (2t + y)t^{-3/5} \text{ and } X = (x - t - y)t^{-1/5}. \quad (2.49) \]

Substituting the expression of \( u \) in (1.7), we obtain the NPDE

\[ 5Q_{TT} - 3TQ_{TX} - XQ_{XX} - 2Q_X + 15Q_XQ_{TX} + 15Q_TQ_{XX} + 5Q_{TXXX} = 0, \quad (2.50) \]

which has two symmetries:

\[ P_1 = \frac{\partial}{\partial Q}, \quad P_2 = \frac{\partial}{\partial X} + \frac{1}{15} T \frac{\partial}{\partial Q}. \]

The symmetry \( P_2 \) gives \( Q(X, Y) = f(z) + (1/15)TX, z = T \) and hence (2.50) reduces to

\[ 75f''(z) - 4z = 0. \]

Solving the above ODE and reverting to the basic variables gives the solution of (1.7) as

\[ u(t, x, y) = \frac{1}{\sqrt{t}} \left\{ \frac{(2t + y)(x - t - y)}{15t^{1/5}} + \frac{2(2t + y)^3}{225t^{1/5}} + \frac{(2t + y)}{t^{3/5}} A_1 + A_2 \right\} - \frac{2t}{9} - \frac{y}{3}, \quad (2.51) \]

where \( A_1 \) and \( A_2 \) are integration constants. The wave profile of solution (2.51) is shown in Figure 10 with 3D, contour and 2D plots using parameter values \( A_1 = -0.3, A_2 = -50 \) with \( t = 1.1 \) and \(-10 \leq x, y \leq 10\).
Next, we invoke the symmetry $P_1 + P_2$. This yields $Q(X, Y) = f(z) + X + (1/15)TX$, $z = T$. Consequently, we have the transformed version of (2.50) as

$$75f'''(z) - 4z - 15 = 0.$$ 

Solving the above ODE and reverting to basic variables gives the solution of (1.7) as

$$u(t, x, y) = \frac{1}{\sqrt{t}} \left\{ \frac{(2t + y)(x - t - y)}{15t^{4/5}} + \frac{2(2t + y)^3}{225t^{9/5}} + \frac{(2t + y)^2}{10t^{6/5}} + \frac{(2t + y)}{t^{3/5}}A_1 + A_2 \right\} - \frac{2t}{9} - \frac{y}{3}. \tag{2.52}$$

The wave profile of solution (2.52) is shown in Figure 11 with 3D, contour and 2D plots using parameter values $A_1 = -3.6, A_2 = 50$ with $t = 1.1$ and $-10 \leq x, y \leq 10$.

3. Conservation laws of (1.7)

In this section, we construct the conservation laws for 2D-gBK equation (1.7) by making use of the multiplier approach [22, 49, 50], but first we give some basic background of the method that we are adopting.
3.1. Fundamental operators and their relationship

Consider the $n$ independent variables $x = (x^1, x^2, \ldots, x^n)$ and $m$ dependent variables $u = (u^1, u^2, \ldots, u^m)$. The derivatives of $u$ with respect to $x$ are defined as

\[ u^\alpha_i = D_i(u^\alpha), \quad u^\alpha_{ij} = D_j D_i(u^\alpha), \ldots, \quad i = 1, \ldots, n, \]

where

\[ D_i = \frac{\partial}{\partial x^i} + u^\alpha_i \frac{\partial}{\partial u^\alpha} + u^\alpha_{ij} \frac{\partial}{\partial u^\alpha_j} + \cdots, \quad i = 1, \ldots, n, \]

is the operator of total differentiation. The collection of all first derivatives $u^\alpha_i$ is denoted by $u^{(1)}$, i.e., $u^{(1)} = \{u^\alpha_i\}$, $\alpha = 1, \ldots, m$, $i = 1, \ldots, n$. In the same vein $u^{(2)} = \{u^\alpha_{ij}\}$, $\alpha = 1, \ldots, m$, $i, j = 1, \ldots, n$ and $u^{(3)} = \{u^\alpha_{ijk}\}$ and likewise $u^{(4)}$ etc. Since $u^\alpha_{ij} = u^\alpha_{ji}$, $u^{(2)}$ contains only $u^\alpha_{ij}$ for $i \leq j$.

Now consider a $k$th-order system of PDEs:

\[ G_{\alpha}(x, u, u^{(1)}, \ldots, u^{(k)}) = 0, \quad \alpha = 1, 2, \ldots, m. \]

The Euler-Lagrange operator, for every $\alpha$, is presented as

\[ \frac{\delta}{\delta u^\alpha}(\Omega_{\alpha} G_{\alpha}) = 0 \]

holds identically [22]. The determining equations for multipliers are obtained by taking the variational derivative of (3.6), namely

\[ \frac{\delta}{\delta u^\alpha}(\Omega_{\alpha} G_{\alpha}) = 0. \]

The moment multipliers are generated from (3.7), the conserved vectors can be derived systematically using (3.6) as the determining equation.

3.2. Construction of conservation laws for (1.7)

Conservation laws of 2D-gBK equation (1.7) are derived by utilizing second-order multiplier $\Omega(t, x, y, u, u_t, u_x, u_y, u_{xx}, u_{xy})$, in Eq (3.7), where $G$ is given as

\[ G \equiv u_{tx} + 6u_x u_{xx} + u_{xxx} + u_{xxy} + 3(u_x u_y)_x + u_{xx} + u_{yy}, \]

and the Euler operator $\delta/\delta u$ is expressed in this case as
Remark 3.1. We notice that this method assists in the construction of conservation laws of (1.7) despite the fact that it possesses no variational principle [51]. Moreover, the presence of arbitrary functions in the multiplier indicates that 2D-gBK (1.7) has infinite number of conserved vectors.

\[
\frac{\delta}{\delta u} = \frac{\partial}{\partial u} - D_t \frac{\partial}{\partial u_t} - D_x \frac{\partial}{\partial u_x} - D_y \frac{\partial}{\partial u_y} + D_t D_x \frac{\partial}{\partial u_{tx}} + D_t D_y \frac{\partial}{\partial u_{ty}} + D_x^2 \frac{\partial}{\partial u_{xx}} + D_y^2 \frac{\partial}{\partial u_{yy}} + D_x D_y \frac{\partial}{\partial u_{xy}}.
\]

Expansion of Eq (3.7) and splitting on diverse derivatives of dependent variable \(u\) gives

\[
\Omega_u = 0, \; \Omega_x = 0, \; \Omega_{yy} = 0, \; \Omega_{yyt} = 0, \; \Omega_{u_xu_y} = 0, \; \Omega_{u_yu_x} = 0, \; \Omega_{u_yu_y} = 0, \; \Omega_{u_y} = 0.
\]

(3.8)

Solution to the above system of equations gives \(\Omega(t, x, y, u, u_t, u_x, u_y, u_{xx}, u_{xy})\) as

\[
\Omega(t, x, y, u, u_t, u_x, u_y, u_{xx}, u_{xy}) = C_1 u_x + f_1(t)y + f_2(t),
\]

(3.9)

with \(C_1\) being an arbitrary constant and \(f_1(t), f_2(t)\) being arbitrary functions of \(t\). Using Eq (3.6), one obtains the following three conserved vectors of Eq (1.7) that correspond to the three multipliers \(u_x, f_1(t)\) and \(f_2(t)\):

**Case 1.** For the first multiplier \(Q_1 = u_x\), the corresponding conserved vector \((T'_1, T^x_1, T^y_1)\) is given by

\[
T'_1 = \frac{1}{2} u_x^2,
\]

\[
T^x_1 = \frac{1}{2} u_x^3 + 2u_x^2 - u_x u_{xx} - \frac{1}{2} u_{xx} u_{xy} + \frac{1}{2} u_x u_{xxy} + u_{xxx} u_x
\]

\[
+ \frac{1}{2} u_{xxx} u_y + \frac{1}{2} u_{xxy} u_{yy} + \frac{1}{2} u_{xyy} u_{x} + u_{xx} u_{xy} + 2u_x u_y + u_y^2,
\]

\[
T^y_1 = \frac{1}{2} u_y u_x - u_x u_{xx} - \frac{1}{2} u_{xx} u_{xy} - \frac{1}{2} u_{xxx} u_{xx},
\]

**Case 2.** For the second multiplier \(Q_2 = f_1(t)\), we obtain the corresponding conserved vector \((T'_2, T^x_2, T^y_2)\) as

\[
T'_2 = u_x f_1(t)y,
\]

\[
T^x_2 = 3y f_1(t) u_x^2 + 3y f_1(t) u_x u_y - y f'_1(t) u,
\]

\[
+ y f_1(t) u_x + y f_1(t) u_{xxx} + y f_1(t) u_{xxy},
\]

\[
T^y_2 = u_y f_1(t)y - u f_1(t).
\]

**Case 3.** Finally, for the third multiplier \(Q_3 = f_2(t)\), the corresponding conserved vector \((T'_3, T^x_3, T^y_3)\) is

\[
T'_3 = u_x f_2(t),
\]

\[
T^x_3 = 3u_x^2 f_2(t) + 3u_x u_y f_2(t) - uf'_2(t) + u_x f_2(t)
\]

\[
+ u_{xxx} f_2(t) + u_{xxy} f_2(t),
\]

\[
T^y_3 = u_y f_2(t).
\]
4. Conclusions

In this paper, we carried out a study on two-dimensional generalized Bogoyavlensky-Konopelchenko equation (1.7). We obtained solutions for Eq (1.7) with the use of Lie symmetry reductions, direct integration, Kudryashov’s and \((G'/G)\)-expansion techniques. We obtained solutions of (1.7) in the form of algebraic, rational, periodic, hyperbolic as well as trigonometric functions. Furthermore, we derived conservation laws of (1.7) by engaging the multiplier method were we obtained three local conserved vectors. We note here that the presence of the arbitrary functions \(f_1(t)\) and \(f_2(t)\) in the multipliers, tells us that one can generate unlimited number of conservation laws for the underlying equation.

Conflict of interest

The authors state no conflicts of interest.

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