Mathematics

## Research article

## The behavior of solutions of a parametric weighted $(p, q)$-Laplacian equation

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Abstract: We study the behavior of solutions for the parametric equation

$$
-\Delta_{p}^{a_{1}} u(z)-\Delta_{q}^{a_{2}} u(z)=\lambda|u(z)|^{q-2} u(z)+f(z, u(z)) \quad \text { in } \Omega, \lambda>0,
$$

under Dirichlet condition, where $\Omega \subseteq \mathbb{R}^{N}$ is a bounded domain with a $C^{2}$-boundary $\partial \Omega, a_{1}, a_{2} \in$ $L^{\infty}(\Omega)$ with $a_{1}(z), a_{2}(z)>0$ for a.a. $z \in \Omega, p, q \in(1, \infty)$ and $\Delta_{p}^{a_{1}}, \Delta_{q}^{a_{2}}$ are weighted versions of $p$ Laplacian and $q$-Laplacian. We prove existence and nonexistence of nontrivial solutions, when $f(z, x)$ asymptotically as $x \rightarrow \pm \infty$ can be resonant. In the studied cases, we adopt a variational approach and use truncation and comparison techniques. When $\lambda$ is large, we establish the existence of at least three nontrivial smooth solutions with sign information and ordered. Moreover, the critical parameter value is determined in terms of the spectrum of one of the differential operators.

Keywords: weighted ( $p, q$ )-Laplacian; resonant Carathéodory function; parametric power term; positive and negative solutions; nodal solutions
Mathematics Subject Classification: 35J20, 35J60

## 1. Introduction

Our goal here is to investigate the existence and nonexistence of nontrivial smooth solutions for the following parametric Dirichlet problem

$$
\left\{\begin{array}{l}
-\Delta_{p}^{a_{1}} u(z)-\Delta_{q}^{a_{2}} u(z)=\lambda|u(z)|^{q-2} u(z)+f(z, u(z)) \quad \text { in } \Omega, \\
\left.u\right|_{\partial \Omega}=0,1<q<p, \lambda>0,
\end{array}\right.
$$

where $\Omega \subseteq \mathbb{R}^{N}$ is a bounded domain with a $C^{2}$-boundary $\partial \Omega$. Given $r \in(1, \infty)$ and $a \in L^{\infty}(\Omega)$ with $a(z)>0$ for a.a. $z \in \Omega$, by $\Delta_{r}^{a}$ we mean the weighted $r$-Laplacian of the form $\Delta_{r}^{a} u=\operatorname{div}\left(a(z)|\nabla u|^{r-2} \nabla u\right)$ for all $u \in W_{0}^{1, r}(\Omega)$. Thus, $\left(P_{\lambda}\right)$ is driven by the operator $-\Delta_{p}^{a_{1}}-\Delta_{q}^{a_{2}}$, whose weights $a_{1}, a_{2}$ are Lipschitz continuous, positive and bounded away from zero. These conditions imply that the integrand corresponding to this differential operator, exhibits balanced growth. However, the fact that the two weights are different, does not allow the use of the nonlinear strong maximum principle (see Pucci and Serrin [25], pp. 111, 120). Instead we use a recent result of Papageorgiou et al. [20], together with an additional comparison argument, which allows us to conclude that the constant sign solutions of the problem satisfy the nonlinear Hopf's lemma. The right-side of $\left(P_{\lambda}\right)$ is the sum of the power term $\lambda|x|^{q-2} x$ and of the Carathéodory function $f(z, x)$. The $\lambda$-parametric term is $(p-1)$-sublinear (recall that $q<p)$, and $f(z, x)$ is $(p-1)$-linear as $x \rightarrow \pm \infty$ and can be resonant with respect to the first eigenvalue of $\left(-\Delta_{p}^{a_{1}}, W_{0}^{1, p}(\Omega)\right)$. We mention that the power of the parametric term (namely $q$ ) is the same with the exponent of the second differential operator $-\Delta_{q}^{a_{2}}$. This distinguishes $\left(P_{\lambda}\right)$ from problems with concave terms, where the power of the parametric term is strictly less than the exponents of all the differential operators in the left-side. Such concave problems, were studied recently by Gasiński and Papageorgiou [5], Gasiński et al. [6] ( $p$-equations), Marano et al. [12], Papageorgiou and Scapellato [19], Papageorgiou and Zhang [22] ( $p, 2$ )-equations), Papageorgiou et al. [17] (anisotropic equations), Papageorgiou and Winkert [21], Papageorgiou and Zhang [23, 24] ( $(p, q)$-equations) and Papageorgiou et al. [18] (nonhomogeneous Robin problems).

Let $\widehat{\lambda}_{1}\left(q, a_{2}\right)>0$ be the principal eigenvalue of $\left(-\Delta_{q}^{a_{2}}, W_{0}^{1, q}(\Omega)\right)$. Using variational tools from the critical point theory, truncation and comparison methods, then $\left(P_{\lambda}\right)$ (for all $\lambda>\bar{\lambda}_{1}\left(q, a_{2}\right)$ ) admits at least three nontrivial smooth solutions (positive, negative, nodal). Moreover, under an additional mild regularity for $f(z, \cdot)$, we get that $\left(P_{\lambda}\right)$ (for all $\lambda<\widehat{\lambda}_{1}\left(q, a_{2}\right)$ ) has no nontrivial solutions.

## 2. Preliminaries

A crucial point is to establish the appropriate spaces, where carrying out the study. Here, $\left(P_{\lambda}\right)$ is analyzed in $W_{0}^{1, p}(\Omega)$ (namely, Sobolev space) and in $C_{0}^{1}(\bar{\Omega})=\left\{u \in C^{1}(\bar{\Omega}):\left.u\right|_{\partial \Omega}=0\right\}$ (classical Banach space). Additionally, $\|\cdot\|$ means the norm of $W_{0}^{1, p}(\Omega)$ with

$$
\|u\|=\|\nabla u\|_{p} \quad \text { for all } u \in W_{0}^{1, p}(\Omega) \quad \text { (by Poincaré inequality). }
$$

$C_{0}^{1}(\bar{\Omega})$ is ordered, with positive (order) cone $C_{+}=\left\{u \in C_{0}^{1}(\bar{\Omega}): u(z) \geq 0\right.$ for all $\left.z \in \bar{\Omega}\right\}$. Now, $C_{+}$has the nonempty interior

$$
\operatorname{int} C_{+}=\left\{u \in C_{+}: u(z)>0 \text { for all } z \in \Omega,\left.\quad \frac{\partial u}{\partial n}\right|_{\partial \Omega}<0\right\},
$$

with $n(\cdot)$ being the outward unit normal on $\partial \Omega$. Let $r \in(1, \infty)$ and $a \in C^{0,1}(\bar{\Omega})$ (that is, $a(\cdot)$ is Lipschitz continuous on $\bar{\Omega}$ ) with $a(z) \geq \widehat{c}_{0}>0$ for all $z \in \bar{\Omega}$.

By $A_{r}^{a}: W_{0}^{1, r}(\Omega) \rightarrow W^{-1, r^{\prime}}(\Omega)=W_{0}^{1, r}(\Omega)^{*}\left(\frac{1}{r}+\frac{1}{r^{\prime}}=1\right)$, we denote the operator

$$
\left\langle A_{r}^{a}(u), h\right\rangle=\int_{\Omega} a(z)|\nabla u|^{r-2}(\nabla u, \nabla h)_{\mathbb{R}^{v}} d z \quad \text { for all } u, h \in W_{0}^{1, r}(\Omega) .
$$

We recall some features of $A_{r}^{a}(\cdot)$ as follows:

- $A_{r}^{a}(\cdot)$ is bounded and continuous;
- $A_{r}^{a}(\cdot)$ is strictly monotone, and hence maximal monotone;
- $A_{r}^{a}(\cdot)$ is of type $(S)_{+}$. It means that, if $u_{n} \xrightarrow{w} u$ in $W_{0}^{1, r}(\Omega)$ and $\limsup _{n \rightarrow \infty}\left\langle A_{r}^{a}\left(u_{n}\right), u_{n}-u\right\rangle \leq 0$, then $u_{n} \rightarrow u$ in $W_{0}^{1, r}(\Omega)$.

Given the eigenvalue problem

$$
-\Delta_{r}^{a} u(z)=\widehat{\lambda}|u(z)|^{r-2} u(z) \text { in } \Omega,\left.u\right|_{\partial \Omega}=0,
$$

we say that $\widehat{\lambda} \in \mathbb{R}$ is an eigenvalue of $\left(-\Delta_{r}^{a}, W_{0}^{1, r}(\Omega)\right)$, if the above problem admits a nontrivial solution $\widehat{u} \in W_{0}^{1, r}(\Omega)$ (namely, eigenfunction of $\widehat{\lambda}$ ). There is a smallest eigenvalue $\widehat{\lambda}_{1}(r, a)>0$. Indeed, consider

$$
\begin{align*}
0 \leq \widehat{\lambda}_{1}(r, a) & =\inf \left[\frac{\int_{\Omega} a(z)|\nabla u|^{r} d z}{\|u\|_{r}^{r}}: u \in W_{0}^{1, r}(\Omega), u \neq 0\right] \\
& =\inf \left[\int_{\Omega} a(z)|\nabla u|^{r} d z: u \in W_{0}^{1, r}(\Omega),\|u\|_{r}=1\right] . \tag{2.1}
\end{align*}
$$

We claim that the infimum in (2.1) is attained. To see this consider a sequence $\left\{u_{n}\right\}_{n \in \mathbb{N}} \subseteq W_{0}^{1, r}(\Omega)$ satisfying $\left\|u_{n}\right\|_{r}=1$ for all $n \in \mathbb{N}$, and $\int_{\Omega} a(z)\left|\nabla u_{n}\right|^{r} d z \downarrow \widehat{\lambda}_{1}(r, a)$. From the boundedness of $\left\{u_{n}\right\}_{n \in \mathbb{N}} \subseteq$ $W_{0}^{1, r}(\Omega)$, it is possible to suppose

$$
\begin{equation*}
u_{n} \xrightarrow{w} \widehat{u}_{1} \text { in } W_{0}^{1, r}(\Omega), u_{n} \rightarrow \widehat{u}_{1} \text { in } L^{r}(\Omega) . \tag{2.2}
\end{equation*}
$$

On account of our hypothesis on the weight $a(\cdot)$, on $L^{r}\left(\Omega, \mathbb{R}^{N}\right) y \rightarrow\left[\int_{\Omega} a(z)|y|^{r} d z\right]^{1 / r}$ is an equivalent norm. From (2.2), since the norm (in Banach space) is weakly lower semicontinuous, also using the Lagrange multiplier rule (Papageorgiou and Kyritsi-Yiallourou [14], p. 76) and the nonlinear regularity theory, after standard calculations we get $\widehat{u}_{1} \in C_{0}^{1}(\bar{\Omega}) \backslash\{0\}$. Additionally, it is clear from (2.1) that we may assume that $\widehat{u}_{1} \in C_{+} \backslash\{0\}$ (just replace $\widehat{u}_{1}$ by $\left|\widehat{u}_{1}\right|$ ). Then the nonlinear Hopf's lemma (Pucci and Serrin [25], pp. 111, 120), gives us $\widehat{u}_{1}=\widehat{u}_{1}(r, a) \in \operatorname{int} C_{+}$. From Jaros [9, Theorem 3.3], we know that $\widehat{\lambda}_{1}(r, a)$ is simple, i.e., if $\widehat{u}_{1}, \widehat{v}_{1}$ are eigenfunctions corresponding to $\widehat{\lambda}_{1}$, then $\widehat{u}_{1}=\vartheta \widehat{v}_{1}$ for certain $\vartheta \in \mathbb{R} \backslash\{0\}$. Also $\hat{\lambda}_{1}(r, a)>0$ is isolated in the spectrum $\sigma(r, a)$ of $\left(-\Delta_{r}^{a}, W_{0}^{1, r}(\Omega)\right)$. For this purpose, let us consider eigenvalues $\left\{\widehat{\lambda}_{n}\right\}_{n \in \mathbb{N}} \subseteq \sigma(r, a)$ satisfying $\widehat{\lambda}_{1}(r, a)<\widehat{\lambda}_{n}$ for all $n \in \mathbb{N}$, and $\widehat{\lambda}_{n} \downarrow \widehat{\lambda}_{1}(r, a)$. So, we can find $\widehat{u}_{n} \in W_{0}^{1, r}(\Omega), \widehat{u}_{n} \neq 0$ such that

$$
-\Delta_{r}^{a} \widehat{u}_{n}=\widehat{\lambda}_{n}\left|\widehat{u}_{n}\right|^{r-2} \widehat{u}_{n} \text { in } \Omega,\left.\widehat{u}_{n}\right|_{\partial \Omega}=0, n \in \mathbb{N} .
$$

By homogeneity we can always assume that $\left\|\widehat{u}_{n}\right\|_{r}=1$ for all $n \in \mathbb{N}$. The nonlinear regularity theory (see Lieberman [11]), implies that there exist $\alpha \in(0,1)$ and $c_{0}>0$ such that

$$
\begin{equation*}
\widehat{u}_{n} \in C_{0}^{1, \alpha}(\bar{\Omega}),\left\|\widehat{u}_{n}\right\|_{C_{0}^{1, \alpha}(\bar{\Omega})} \leq c_{0} \quad \text { for all } n \in \mathbb{N} . \tag{2.3}
\end{equation*}
$$

The compact embedding $C_{0}^{1, \alpha}(\bar{\Omega}) \hookrightarrow C_{0}^{1}(\bar{\Omega})$ and (2.3), ensure one can suppose

$$
u_{n} \rightarrow \widetilde{u} \text { in } C_{0}^{1}(\bar{\Omega}),\|\widetilde{u}\|_{r}=1,
$$

$$
\begin{aligned}
& \Rightarrow \quad-\Delta_{r}^{a} \widetilde{u}=\widehat{\lambda}_{1}(r, a)|\widetilde{u}|^{r-2} \widetilde{u} \text { in } \Omega,\left.\widetilde{u}\right|_{\partial \Omega}=0, \\
& \Rightarrow \quad \widetilde{u}=\vartheta \widehat{u}_{1} \in \operatorname{int} C_{+} \quad \text { for some } \vartheta>0,
\end{aligned}
$$

and hence $\widehat{u}_{n} \in \operatorname{int} C_{+}$for all $n \geq n_{0}$, which leads to contradiction with Jaros [9, Corollary 3.2]. This proves that $\bar{\lambda}_{1}(r, a)>0$ is isolated. The Ljusternik-Schnirelmann minimax scheme (see, for example, Gasiński and Papageorgiou [3]), ensures a whole strictly increasing sequence of distinct eigenvalues $\left\{\bar{\lambda}_{n}\right\}_{n \in \mathbb{N}}$ such that $\widehat{\lambda}_{n} \rightarrow+\infty$. If $r=2$, then these eigenvalues exhaust the spectrum. If $r \neq 2$, then it is not known if the LS-eigenvalues fully describe $\sigma(r, a)$. Moreover, every $\widehat{\lambda} \in \sigma(r, a) \backslash\left\{\widehat{\lambda}_{1}(r, a)\right\}$ has eigenfunctions which are nodal functions (that is, sign-changing functions), see again Jaros [9, Corollary 3.2]. We can easily check that $\sigma(r, a) \subseteq\left[\widehat{\lambda}_{1}(r, a),+\infty\right)$ is closed. So, we can define the second eigenvalue of $\left(-\Delta_{r}^{a}, W_{0}^{1, r}(\Omega)\right.$ ) by

$$
\widehat{\lambda}_{2}(r, a)=\inf \left[\widehat{\lambda} \in \sigma(r, a): \widehat{\lambda}_{1}(r, a)<\widehat{\lambda}\right] .
$$

Reasoning as in Cuesta et al. [1], one can show that $\widehat{\lambda}_{2}(r, a)$ corresponds to the second LS-eigenvalue and

$$
\begin{equation*}
\widehat{\lambda}_{2}(r, a)=\inf _{\widehat{\gamma} \bar{\Gamma}-1 \leq t \leq 1} \max _{\Omega} a(z)|\nabla \widehat{\gamma}(t)|^{r} d z \tag{2.4}
\end{equation*}
$$

where $\widehat{\Gamma}=\left\{\widehat{\gamma} \in C([-1,1], M): \widehat{\gamma}(-1)=-\widehat{u}_{1}(r, a), \widehat{\gamma}(1)=\widehat{u}_{1}(r, a)\right\}$ with $M=W_{0}^{1, r}(\Omega) \cap \partial B_{1}^{L^{r}}\left(\partial B_{1}^{L^{r}}=\right.$ $\left.\left\{u \in L^{r}(\Omega):\|u\|_{r}=1\right\}\right)$ and $\widehat{u}_{1}(r, a)$ is the positive, $L^{r}$-normalized eigenfunction (i.e., $\left\|\widehat{u}_{1}(r, a)\right\|_{r}=1$ ) corresponding to $\widehat{\lambda}_{1}(r, a)>0$. Recall that $\widehat{u}_{1}=\widehat{u}_{1}(r, a) \in \operatorname{int} C_{+}$.

The above features lead to the following proposition.
Proposition 2.1. Let $\eta \in L^{\infty}(\Omega), \eta(z) \leq \widehat{\lambda}_{1}(r, a)$ for a.a. $z \in \Omega$ and the inequality be strict on $a$ set of positive Lebesgue measure. Then, $\int_{\Omega} a(z)|\nabla u|^{p} d z-\int_{\Omega} \eta(z)|u|^{p} d z \geq \widehat{c}\|\nabla u\|^{p}$ for some $\widehat{c}>0$, all $u \in W_{0}^{1, p}(\Omega)$.

If $u: \Omega \rightarrow \mathbb{R}$ is measurable, let $u^{ \pm}(z)=\max \{ \pm u(z), 0\}$ for all $z \in \Omega$. If $u \in W_{0}^{1, p}(\Omega)$, then $u^{ \pm} \in W_{0}^{1, p}(\Omega)$ and $u=u^{+}-u^{-},|u|=u^{+}+u^{-}$. Also, if $u, v: \Omega \rightarrow \mathbb{R}$ are measurable with $u(z) \leq v(z)$ for a.a. $z \in \Omega$, then we set:

$$
[u, v]=\left\{h \in W_{0}^{1, p}(\Omega): u(z) \leq h(z) \leq v(z) \text { for a.a. } z \in \Omega\right\} .
$$

Now, $\operatorname{int}_{C_{0}^{1}(\bar{\Omega})}[u, v]$ means the interior in $C_{0}^{1}(\bar{\Omega})$ of $[u, v] \cap C_{0}^{1}(\bar{\Omega})$. For a Banach space $X$ and $\varphi \in$ $C^{1}(X)$, let $K_{\varphi}=\left\{u \in X: \varphi^{\prime}(u)=0\right\}$ (namely, critical set of $\varphi$ ). For $c \in \mathbb{R}$, let $\varphi^{c}=\{u \in X: \varphi(u) \leq c\}$, $K_{\varphi}^{c}=\left\{u \in K_{\varphi}: \varphi(u)=c\right\}$.

For a measurable function $g: \Omega \rightarrow \mathbb{R}$, then $0 \leq g$ if and only if for every $K \subseteq \Omega$ compact, one has $0<c_{K} \leq g(z)$ for a.a. $z \in K$. When $g \in C(\Omega)$ and $g(z)>0$ for all $z \in \Omega$, clearly $0 \leq g$.

In the study of $\left(P_{\lambda}\right)$, we use the assumption $H_{0}$ stated as follows:
$H_{0}: a_{1}, a_{2} \in C^{0,1}(\bar{\Omega})$ and $0<c_{1} \leq a_{1}(z), a_{2}(z)$ for all $z \in \bar{\Omega}$.
Remark 2.1. If $\widehat{a}(z, y)=a_{1}(z)|y|^{p-2} y+a_{2}(z)|y|^{q-2} y$ for all $(z, y) \in \Omega \times \mathbb{R}^{N}$, then we see that $\operatorname{div} a(z, \nabla u)=$ $\Delta_{p}^{a_{1}} u+\Delta_{q}^{a_{2}} u$ for all $u \in W_{0}^{1, p}(\Omega)$. The primitive of $\widehat{a}(z, y)$ is the function $\widehat{G}(z, y)=\frac{a_{1}(z)}{p}|y|^{p}+\frac{a_{2}(z)}{q}|y|^{q}$ for all $(z, y) \in \Omega \times \mathbb{R}^{N}$. On account of $H_{0}$, we see that $\widehat{G}(\cdot, \cdot)$ exhibits balanced growth, namely

$$
\frac{c_{1}}{p}|y|^{p} \leq \widehat{G}(z, y) \leq c_{2}\left[1+|y|^{p}\right] \quad \text { for some } c_{2}>0 \text { and all }(z, y) \in \Omega \times \mathbb{R}^{N} .
$$

We also consider the following set of assumptions on the data:
$H_{1}: f: \Omega \times \mathbb{R} \rightarrow \mathbb{R}$ is Carathéodory with $f(z, 0)=0$ for a.a. $z \in \Omega$, and
(i) for every $\rho>0$, there exists $a_{\rho} \in L^{\infty}(\Omega)$ with $|f(z, x)| \leq a_{\rho}(z)$ for a.a. $z \in \Omega$, all $|x| \leq \rho$;
(ii) $\limsup _{x \rightarrow \pm \infty} \frac{f(z, x)}{|x|^{p-2} x} \leq \widehat{\lambda}_{1}\left(p, a_{1}\right)$ uniformly for a.a. $z \in \Omega$;
(iii) If $F(z, x)=\int_{0}^{x} f(z, s) d s$, there is $\tau \in(q, p)$ with $\lim _{x \rightarrow \pm \infty} \frac{f(z, x) x-p F(z, x)}{|x|^{\tau}}=+\infty$ uniformly for a.a. $z \in \Omega$;
(iv) $\lim _{x \rightarrow 0} \frac{f(z, x)}{|x|^{q-2} x}=0$ uniformly for a.a. $z \in \Omega$;
(v) for every $s>0$, there exists $m_{s}>0$ with $m_{s} \leq f(z, x) x$ for a.a. $z \in \Omega$, all $|x| \geq s$.

Remark 2.2. According to $H_{1}\left(\right.$ ii), we can have resonance of $\left(P_{\lambda}\right)$ with respect to $\widehat{\lambda}_{1}\left(p, a_{1}\right)>0$. By the proof of Proposition 3.1, we will see that this phenomenon originates from the left of $\hat{\lambda}_{1}\left(p, a_{1}\right)$ in the sense that

$$
\lim _{x \rightarrow \pm \infty}\left[\widehat{\lambda}_{1}\left(p, a_{1}\right)|x|^{p}-p F(z, x)\right]=+\infty \quad \text { uniformly for a.a. } z \in \Omega .
$$

We stress that this ensures the coercivity of the corresponding energy functional. Therefore, we can use classical tools of the calculus of variations. Assumption $H_{1}(i v)$ does not permit the presence of a concave term and this changes the geometry of our problem compared to those of the "concave" works mentioned in the Introduction. Finally we mention that assumptions $H_{1}$ imply that

$$
\begin{equation*}
|f(z, x)| \leq a(z)\left[1+|x|^{p-1}\right] \quad \text { for a.a. } z \in \Omega \text {, all } x \in \mathbb{R}, a \in L^{\infty}(\Omega)_{+} . \tag{2.5}
\end{equation*}
$$

When $q=2$, we improve our conclusion about the nodal solution, provided we add a perturbed monotonicity assumption for $f(z, \cdot)$, as follows
$H_{1}^{\prime}: H_{1}$ hold (with $q=2$ ) and
(vi) for every $\rho>0$, there exists $\widehat{\xi}_{\rho}>0$ such that for a.a. $z \in \Omega$, the function $x \rightarrow f(z, x)+\widehat{\xi}_{\rho}|x|^{p-2} x$ is nondecreasing on $[-\rho, \rho]$.

Finally, we can have a nonexistence result for $\left(P_{\lambda}\right)$ provided we add a growth restriction for $f(z, \cdot)$, as follows
$H_{1}^{\prime \prime}: H_{1}$ hold and
(vi) $f(z, x) x \leq \widehat{\lambda}_{1}\left(p, a_{1}\right)|x|^{p}$ for a.a. $z \in \Omega$, all $x \in \mathbb{R}$.

## 3. Positive and negative solutions

The existence of positive and negative solutions for $\left(P_{\lambda}\right)$ is established in the case $\lambda>\bar{\lambda}_{1}\left(q, a_{2}\right)$. We obtain smallest positive and biggest negative solutions. These solutions of $\left(P_{\lambda}\right)$ (namely, extremal constant sign solutions) play a crucial role in Section 4 to generate a nodal solution.

Proposition 3.1. Let $H_{0}, H_{1}$ be satisfied, and $\lambda>\widehat{\lambda}_{1}\left(q, a_{2}\right)$. Then $\left(P_{\lambda}\right)$ admits solutions $u_{\lambda} \in \operatorname{int} C_{+}$, $v_{\lambda} \in-$ int $C_{+}$.

Proof. Let $\varphi_{\lambda}^{+}: W_{0}^{1, p}(\Omega) \rightarrow \mathbb{R}$ be a $C^{1}$-functional given as

$$
\varphi_{\lambda}^{+}(u)=\frac{1}{p} \int_{\Omega} a_{1}(z)|\nabla u|^{p} d z+\frac{1}{q} \int_{\Omega} a_{2}(z)|\nabla u|^{q} d z-\frac{\lambda}{q}\left\|u^{+}\right\|_{q}^{q}-\int_{\Omega} F\left(z, u^{+}\right) d z
$$

for all $u \in W_{0}^{1, p}(\Omega)$. We discuss the properties of $\varphi_{\lambda}^{+}(\cdot)$ to obtain a positive solution of $\left(P_{\lambda}\right)$. As already mentioned the coercivity of functionals is a crucial key to apply the direct methods of calculus of variations.

Claim: $\varphi_{\lambda}^{+}(\cdot)$ is coercive.
Arguing by contradiction, suppose that there is $\left\{u_{n}\right\}_{n \in \mathbb{N}} \subseteq W_{0}^{1, p}(\Omega)$ satisfying

$$
\begin{align*}
& \varphi_{\lambda}^{+}\left(u_{n}\right) \leq c_{3} \quad \text { for some } c_{3}>0, \text { all } n \in \mathbb{N},  \tag{3.1}\\
& \left\|u_{n}\right\| \rightarrow \infty \quad \text { as } n \rightarrow \infty \tag{3.2}
\end{align*}
$$

If $\left\{u_{n}^{+}\right\}_{n \in \mathbb{N}} \subseteq W_{0}^{1, p}(\Omega)$ is bounded, then from (3.1) we deduce the boundedness of $\left\{u_{n}^{-}\right\}_{n \in \mathbb{N}} \subseteq W_{0}^{1, p}(\Omega)$. Consequently, we get the boundedness of $\left\{u_{n}\right\}_{n \in \mathbb{N}} \subseteq W_{0}^{1, p}(\Omega)$, which contradicts (3.2). Therefore, one can suppose

$$
\begin{equation*}
\left\|u_{n}^{+}\right\| \rightarrow \infty \quad \text { as } n \rightarrow \infty . \tag{3.3}
\end{equation*}
$$

We set $y_{n}=\frac{u_{n}^{+}}{\left\|u_{n}^{+}\right\|}, n \in \mathbb{N}$. Then $\left\|y_{n}\right\|=1, y_{n} \geq 0$ for all $n \in \mathbb{N}$. So, we suppose

$$
\begin{equation*}
y_{n} \xrightarrow{w} y \text { in } W_{0}^{1, p}(\Omega) \text { and } y_{n} \rightarrow y \text { in } L^{p}(\Omega), y \geq 0 . \tag{3.4}
\end{equation*}
$$

From (3.1) we have

$$
\begin{align*}
& \frac{1}{p} \int_{\Omega} a_{1}(z)\left|\nabla u_{n}\right|^{p} d z+\frac{1}{q} \int_{\Omega} a_{2}(z)\left|\nabla u_{n}\right|^{q} d z \leq c_{3}+\frac{\lambda}{q}\left\|u_{n}^{+}\right\|_{q}^{q}+\int_{\Omega} F\left(z, u_{n}^{+}\right) d z  \tag{3.5}\\
\Rightarrow \quad & \frac{1}{p} \int_{\Omega} a_{1}(z)\left|\nabla y_{n}\right|^{p} d z+\frac{1}{q\left\|u_{n}^{+}\right\|^{p-q}} \int_{\Omega} a_{2}(z)\left|\nabla y_{n}\right|^{q} d z \\
& \leq \frac{c_{3}}{\left\|u_{n}^{+}\right\|^{p}}+\frac{\lambda}{q\left\|u_{n}^{+}\right\|^{p-q}}\left\|y_{n}\right\|_{q}^{q}+\int_{\Omega} \frac{F\left(z, u_{n}^{+}\right)}{\left\|u_{n}^{+}\right\|^{p}} d z \quad \text { for all } n \in \mathbb{N} . \tag{3.6}
\end{align*}
$$

Assumption $H_{1}(i i)$ leads to

$$
\begin{align*}
& \frac{F\left(\cdot, u_{n}^{+}(\cdot)\right)}{\left\|u_{n}^{+}\right\|^{p}} \xrightarrow{w} \frac{1}{p} \eta y^{p} \quad \text { in } L^{1}(\Omega)  \tag{3.7}\\
& \text { with } \eta \in L^{\infty}(\Omega) \text { satisfying } \eta(z) \leq \widehat{\lambda}_{1}\left(p, a_{1}\right) \text { for a.a. } z \in \Omega \tag{3.8}
\end{align*}
$$

Letting $n \rightarrow \infty$ in (3.6), by (3.3), (3.4), (3.7) and the fact that $q<p$, we deduce that

$$
\begin{equation*}
\int_{\Omega} a_{1}(z)|\nabla y|^{p} d z \leq \int_{\Omega} \eta(z) y^{p} d z \tag{3.9}
\end{equation*}
$$

If $\eta \not \equiv \widehat{\lambda}_{1}\left(p, a_{1}\right)$ (see (3.8)), then from (3.9) one has $\widehat{c}\|y\|^{p} \leq 0$ (see Proposition 2.1), and hence $y=0$. From (3.4) and (3.6), we see that $\left\|\nabla y_{n}\right\|_{p} \rightarrow 0$, which leads to contradiction with $\left\|y_{n}\right\|=1$ for all $n \in \mathbb{N}$.

If $\eta(z)=\widehat{\lambda}_{1}\left(p, a_{1}\right)$ for a.a. $z \in \Omega$, again from (3.9) one has $\int_{\Omega} a_{1}(z)|\nabla y|^{p} d z=\widehat{\lambda}_{1}\left(p, a_{1}\right)\|y\|_{p}^{p}$, and hence $y=\vartheta \widehat{u}_{1}\left(p, a_{1}\right)$ for some $\vartheta \geq 0$.

If $\vartheta=0$, then $y=0$ which leads to contradiction with $\left\|y_{n}\right\|=1$ for all $n \in \mathbb{N}$.
If $\vartheta>0$, then $y \in \operatorname{int} C_{+}$and so we have $u_{n}^{+}(z) \rightarrow+\infty$ for a.a. $z \in \Omega$. By $H_{1}(i i i)$ given $\xi>0$, there is $M=M(\xi)>0$ satisfying

$$
\begin{equation*}
f(z, x) x-p F(z, x) \geq \xi|x|^{\tau} \quad \text { for a.a. } z \in \Omega \text {, all }|x| \geq M . \tag{3.10}
\end{equation*}
$$

Additionally

$$
\begin{aligned}
& \frac{d}{d x}\left[\frac{F(z, x)}{|x|^{p}}\right]=\frac{f(z, x)|x|^{p}-p|x|^{p-2} x F(z, x)}{|x|^{2 p}} \\
&=\frac{f(z, x) x-p F(z, x)}{|x|^{p} x}\left\{\begin{array}{ll}
\geq \frac{\xi}{x^{p-\tau+1}} & \text { if } x \geq M \\
\leq \frac{\xi}{\mid x x^{p-x}} & \text { if } x \leq-M
\end{array} \quad(\text { see (3.10)), }\right. \\
& \Rightarrow \quad \frac{F(z, y)}{|y|^{p}}-\frac{F(z, x)}{|x|^{p}} \geq-\frac{\xi}{p-\tau}\left[\frac{1}{|y|^{p-\tau}}-\frac{1}{|x|^{p-\tau}}\right]
\end{aligned}
$$

for a.a. $z \in \Omega$, for all $|y| \geq|x| \geq M$. Letting $|y| \rightarrow \infty$, by $H_{1}(i i)$ we deduce that

$$
\begin{align*}
& \frac{\widehat{\lambda}_{1}\left(p, a_{1}\right)}{p}-\frac{F(z, x)}{|x|^{p}} \geq \frac{\xi}{p-\tau} \frac{1}{|x|^{p-\tau}}, \\
\Rightarrow & \frac{\widehat{\lambda}_{1}\left(p, a_{1}\right) \mid x x^{p}-p F(z, x)}{|x|^{\tau}} \geq \frac{\xi p}{p-\tau} \quad \text { for a.a. } z \in \Omega, \text { all }|x| \geq M, \\
\Rightarrow & \lim _{x \rightarrow \pm \infty} \frac{\widehat{\lambda}_{1}\left(p, a_{1}\right)|x|^{p}-p F(z, x)}{|x|^{\tau}}=+\infty \quad \text { uniformly for a.a. } z \in \Omega . \tag{3.11}
\end{align*}
$$

Now, (3.5) gives us

$$
\begin{align*}
& \frac{1}{p} \int_{\Omega}\left[\widehat{\lambda}_{1}\left(p, a_{1}\right)\left(u_{n}^{+}\right)^{p}-p F\left(z, u_{n}^{+}\right)\right] d z \leq c_{3}+\lambda\left\|u_{n}^{+}\right\|_{q}^{q}, \\
\Rightarrow & \frac{1}{p} \int_{\Omega} \frac{\widehat{\lambda}_{1}\left(p, a_{1}\right)\left(u_{n}^{+}\right)^{p}-p F\left(z, u_{n}^{+}\right)}{\left(u_{n}^{+}\right)^{\tau}} y_{n}^{\tau} d z \leq \frac{c_{3}}{\left\|u_{n}^{+}\right\|^{\tau}}+\frac{\lambda c_{4}}{\left\|u_{n}^{+}\right\|^{\tau-q}}, \tag{3.12}
\end{align*}
$$

for some $c_{4}>0$, for all $n \in \mathbb{N}$. For $n \rightarrow \infty$ in (3.12) combining (3.3), (3.11), Fatou's lemma and recalling that $\tau>q$, we leads to contradiction. The boundedness of $\left\{u_{n}^{+}\right\}_{n \in \mathbb{N}} \subseteq W_{0}^{1, p}(\Omega)$ is so established. This implies the boundedness of $\left\{u_{n}\right\}_{n \in \mathbb{N}} \subseteq W_{0}^{1, p}(\Omega)$ (see (3.1)), which contradicts (3.2). This argument establishes the coercivity of $\varphi_{\lambda}^{+}(\cdot)$, as stated in the Claim. Next, we observe that $\varphi_{\lambda}^{+}(\cdot)$ is sequentially weakly lower semicontinuous (by Sobolev embedding theorem). This fact, the Claim and the Weierstrass-Tonelli theorem, lead to the existence of a $u_{\lambda} \in W_{0}^{1, p}(\Omega)$ satisfying

$$
\begin{equation*}
\varphi_{\lambda}^{+}\left(u_{\lambda}\right)=\inf \left[\varphi_{\lambda}^{+}(u): u \in W_{0}^{1, p}(\Omega)\right] . \tag{3.13}
\end{equation*}
$$

So, $H_{1}(i v)$ for fixed $\varepsilon>0$, gives us $\delta=\delta(\varepsilon)>0$ with

$$
\begin{equation*}
|F(z, x)| \leq \frac{\varepsilon}{q}|x|^{q} \quad \text { for a.a. } z \in \Omega, \text { all }|x| \leq \delta \tag{3.14}
\end{equation*}
$$

But $\widehat{u}_{1}\left(q, a_{2}\right) \in \operatorname{int} C_{+}$(Section 2) ensures there exists $t \in(0,1)$ small enough to get

$$
\begin{equation*}
0 \leq \widehat{t u}_{1}\left(q, a_{2}\right)(z) \leq \delta \quad \text { for all } z \in \bar{\Omega} \tag{3.15}
\end{equation*}
$$

Therefore,

$$
\varphi_{\lambda}^{+}\left(\widehat{t u}_{1}\left(q, a_{2}\right)\right) \leq \frac{t^{p}}{p} \int_{\Omega} a_{1}(z)\left|\nabla \widehat{u}_{1}\left(q, a_{2}\right)\right|^{p} d z+\frac{t^{q}}{q}\left[\widehat{\lambda}_{1}\left(q, a_{2}\right)+\varepsilon-\lambda\right]
$$

(see (3.14), (3.15), recall $\left.\left\|\widehat{u}_{1}\left(q, a_{2}\right)\right\|_{q}=1\right)$. If we choose $\varepsilon \in\left(0, \lambda-\widehat{\lambda}_{1}\left(q, a_{2}\right)\right)$, then

$$
\begin{equation*}
\varphi_{\lambda}^{+}\left(\widehat{u}_{1}\left(q, a_{2}\right)\right) \leq c_{5} t^{p}-c_{6} t^{q} \quad \text { for some } c_{5}, c_{6}>0 \tag{3.16}
\end{equation*}
$$

As $p>q$, we choose $t \in(0,1)$ appropriately (i.e., even smaller if necessary), then from (3.16) we get

$$
\begin{aligned}
& \varphi_{\lambda}^{+}\left(\widehat{t u}_{1}\left(q, a_{2}\right)\right)<0, \\
\Rightarrow \quad & \varphi_{\lambda}^{+}\left(u_{\lambda}\right)<0=\varphi_{\lambda}^{+}(0) \quad(\text { recall }(3.13)),
\end{aligned}
$$

and so $u_{\lambda} \neq 0$. Again (3.13) leads to $\left(\varphi_{\lambda}^{+}\right)^{\prime}\left(u_{\lambda}\right)=0$, which implies

$$
\begin{equation*}
\left\langle A_{p}^{a_{1}}\left(u_{\lambda}\right), h\right\rangle+\left\langle A_{q}^{a_{2}}\left(u_{\lambda}\right), h\right\rangle=\int_{\Omega}\left[\lambda\left(u_{\lambda}^{+}\right)^{q-1}+f\left(z, u_{\lambda}^{+}\right)\right] h d z \quad \text { for all } h \in W_{0}^{1, p}(\Omega) \tag{3.17}
\end{equation*}
$$

Equation (3.17) for the test function $h=-u_{\lambda}^{-} \in W_{0}^{1, p}(\Omega)$, leads to the inequality $c_{1}\left\|\nabla u_{\lambda}^{-}\right\|_{p}^{p} \leq 0$ (see $H_{0}$ ), and hence $u_{\lambda} \geq 0, u_{\lambda} \neq 0$. Thus $u_{\lambda}$ is a positive solution of ( $P_{\lambda}$ ) (see (3.17)). Ladyzhenskaya and Ural'tseva [10, Theorem 7.1] ensures that $u_{\lambda} \in L^{\infty}(\Omega)$. Consequently, the regularity theory of Lieberman [11] implies $u_{\lambda} \in C_{+} \backslash\{0\}$. Now, Papageorgiou et al. [20, Proposition 2.2] gives us

$$
\begin{equation*}
0<u_{\lambda}(z) \quad \text { for all } z \in \Omega . \tag{3.18}
\end{equation*}
$$

We can continue the proof of [20, Proposition 2.2], since now we have more regularity (namely now $u_{\lambda} \in C_{+} \backslash\{0\}$ ). So, let $z_{1} \in \partial \Omega$ and set $z_{2}=z_{1}-2 \rho n$ with $\rho \in(0,1)$ small and $n=n\left(z_{1}\right)$ is the outward unit normal at $z_{1}$. As in [20], we consider the annulus $D=\left\{z \in \Omega: \rho<\left|z-z_{2}\right|<2 \rho\right\}$ and let $m=\min \left\{u(z): z \in \partial B_{\rho}\left(z_{2}\right)>0\right\}$ (see (3.18)). From the proof in [20], for $\vartheta \in(0, m)$ small, there is $y \in C^{1}(\bar{D}) \cap C^{2}(D)$ satisfying the inequality $-\Delta_{p}^{a_{1}} y-\Delta_{q}^{a_{2}} y \leq 0$ in $D$ with $y\left(z_{1}\right)=0, \frac{\partial y}{\partial n}\left(z_{1}\right)<0$. We know that $-\Delta_{p}^{a_{1}} u_{\lambda}-\Delta_{q}^{a_{2}} u_{\lambda} \geq 0$ in $\Omega$. So, from the weak comparison principle (Pucci and Serrin [25], p. 61), one has $y(z) \leq u_{\lambda}(z)$ for all $z \in D$. It follows that $\frac{\partial u_{\lambda}}{\partial n}\left(z_{1}\right) \leq \frac{\partial y}{\partial n}\left(z_{1}\right)<0$, and so $u_{\lambda} \in \operatorname{int} C_{+}$. Similarly working with $\varphi_{\lambda}^{-}: W_{0}^{1, p}(\Omega) \rightarrow \mathbb{R}$ of the form

$$
\varphi_{\lambda}^{-}(u)=\frac{1}{p} \int_{\Omega} a_{1}(z)|\nabla u|^{p} d z+\frac{1}{q} \int_{\Omega} a_{2}(z)|\nabla u|^{q} d z-\frac{\lambda}{q}\left\|u^{-}\right\|_{q}^{q}-\int_{\Omega} F\left(z,-u^{-}\right) d z
$$

for all $u \in W_{0}^{1, p}(\Omega)$, we get a negative solution $v_{\lambda} \in-\operatorname{int} C_{+}$for problem $\left(P_{\lambda}\right)\left(\lambda>\widehat{\lambda}_{1}\left(q, a_{2}\right)\right)$.

Remark 3.1. An alternative way to show that $u_{\lambda} \in \operatorname{int} C_{+}$, is the following one. Let $\widehat{d}(z)=d(z, \partial \Omega)$ for all $z \in \bar{\Omega}$. By Gilbarg and Trudinger [7, Lemma 14.16], we can find $\delta_{0}>0$ such that $\widehat{d} \in C^{2}\left(\bar{\Omega}_{0}\right)$, where $\Omega_{0}=\left\{z \in \bar{\Omega}: \widehat{d}(z) \leq \delta_{0}\right\}$. It follows that $\widehat{d} \in \operatorname{int} C_{+}$. From Rademacher's theorem (see Gasiński and Papageorgiou [3], p. 56), we know that $a_{1}, a_{2}$ are both differentiable a.e. in $\Omega$. So, by taking $\delta_{0}>0$ even smaller if necessary we can have $\left.\frac{\partial a_{1}}{\partial n}\right|_{\Omega_{0}},\left.\frac{\partial a_{2}}{\partial n}\right|_{\Omega_{0}} \leq 0$. On account of (3.18), we can find $t \in(0,1)$ small such that $w=t \widehat{d} \leq \bar{u}_{\lambda}$ on $\partial \Omega_{0}$. Additionally, [7, Lemma 14.17] leads to

$$
-\Delta_{p}^{a_{1}} w-\Delta_{q}^{a_{2}} w \leq 0 \leq-\Delta_{p}^{a_{1}} u_{\lambda}-\Delta_{q}^{a_{2}} u_{\lambda} \quad \text { in } \Omega_{0}\left(\text { see } H_{1}(v)\right), \quad w \leq u_{\lambda} \quad \text { on } \partial \Omega_{0} .
$$

Then the weak comparison principle (see Pucci and Serrin [25], p. 61), gives us $w \leq u_{\lambda}$ in $\Omega_{0}$. Hence for a certain $\widehat{t} \in(0,1)$ small satisfying $\widehat{t d} \leq u_{\lambda}$ in $\Omega$, we get $u_{\lambda} \in \operatorname{int} C_{+}$.

We now establish the existence of smallest positive and biggest negative solutions. From $H_{1}(i v)$ and (2.5), fixed $\varepsilon>0$, there exists a constant $c_{7}=c_{7}(\varepsilon)>0$ satisfying

$$
\begin{align*}
& f(z, x) x \geq-\varepsilon|x|^{q}-c_{7}|x|^{p} \quad \text { for a.a. } z \in \Omega \text {, all } x \in \mathbb{R}, \\
\Rightarrow \quad & \lambda|x|^{q}+f(z, x) x \geq[\lambda-\varepsilon]|x|^{q}-c_{7}|x|^{p} \quad \text { for a.a. } z \in \Omega, \text { all } x \in \mathbb{R} . \tag{3.19}
\end{align*}
$$

Observe that (3.19) leads to the following auxiliary Dirichlet problem

$$
\left\{\begin{array}{l}
-\Delta_{p}^{a_{1}} u(z)-\Delta_{q}^{a_{2}} u(z)=[\lambda-\varepsilon]|u(z)|^{q-2} u(z)-c_{7}|u(z)|^{p-2} u(z) \quad \text { in } \Omega,  \tag{3.20}\\
\left.u\right|_{\partial \Omega}=0, \lambda>0 .
\end{array}\right.
$$

Proposition 3.2. Let $H_{0}$ be satisfied, $\lambda>\widehat{\lambda}_{1}\left(q, a_{2}\right)$ and $\varepsilon \in\left(0, \lambda-\widehat{\lambda}_{1}\left(q, a_{2}\right)\right)$. Then (3.20) admits a unique positive solution $\bar{u}_{\lambda} \in \operatorname{int} C_{+}$. Additionally, as (3.20) is odd, then it admits a unique negative solution $\bar{v}_{\lambda}=-\bar{u}_{\lambda} \in-\operatorname{int} C_{+}$.

Proof. We start discussing the existence of a positive solution for problem (3.20). To this end let $\psi_{\lambda}^{+}: W_{0}^{1, p}(\Omega) \rightarrow \mathbb{R}$ defined by

$$
\psi_{\lambda}^{+}(u)=\frac{1}{p} \int_{\Omega} a_{1}(z)|\nabla u|^{p} d z+\frac{1}{q} \int_{\Omega} a_{2}(z)|\nabla u|^{q} d z-\frac{\lambda-\varepsilon}{q}\left\|u^{+}\right\|_{q}^{q}+\frac{c_{7}}{p}\left\|u^{+}\right\|_{p}^{p}
$$

for all $u \in W_{0}^{1, p}(\Omega)$. Since $q<p$, we see that $\psi_{\lambda}^{+}(\cdot)$ is coercive. Also, it is sequentially weakly lower semicontinuous. By using the similar arguments to the ones in the proof of Proposition 3.1, one can find $\bar{u}_{\lambda} \in W_{0}^{1, p}(\Omega)$ positive solution to (3.20) (i.e., $\bar{u}_{\lambda} \geq 0, \bar{u}_{\lambda} \neq 0$ ) and also $\bar{u}_{\lambda} \in \operatorname{int} C_{+}$. To establish the uniqueness of $\bar{u}_{\lambda}$, we need the functional $j: L^{1}(\Omega) \rightarrow \overline{\mathbb{R}}=\mathbb{R} \cup\{+\infty\}$ of the form

$$
j(u)= \begin{cases}\frac{1}{p} \int_{\Omega} a_{1}(z)\left|\nabla u^{1 / q}\right|^{p} d z+\frac{1}{q} \int_{\Omega} a_{2}(z)\left|\nabla u^{1 / q}\right|^{q} d z & \text { if } u \geq 0, u^{1 / q} \in W_{0}^{1, p}(\Omega)  \tag{3.21}\\ +\infty & \text { otherwise. }\end{cases}
$$

The convexity of (3.21) follows from Díaz and Saá [2, Lemma 1]. We introduce dom $j=\left\{u \in L^{1}(\Omega)\right.$ : $j(u)<+\infty\}$ and argue by contradiction. Suppose that $\bar{w}_{\lambda}$ is another positive solution of (3.20). Of course, $\bar{w}_{\lambda} \in \operatorname{int} C_{+}$and Papageorgiou et al. [16, Proposition 4.1.22] give us $\frac{\bar{u}_{\lambda}}{\bar{w}_{\lambda}} \in L^{\infty}(\Omega)$ and $\frac{\bar{w}_{\lambda}}{\bar{u}_{\lambda}} \in$
$L^{\infty}(\Omega)$. Hence if $h=\bar{u}_{\lambda}^{q}-\bar{w}_{\lambda}^{q}$, a sufficiently small $|t|<1$ leads to $\bar{u}_{\lambda}^{q}+t h \in \operatorname{dom} j, \bar{w}_{\lambda}^{q}+t h \in \operatorname{dom} j$. Since (3.21) is convex, we have that it is also Gateaux differentiable (in the direction $h$ ) at $\bar{u}_{\lambda}^{q}$ and at $\bar{w}_{\lambda}^{q}$. Using chain rule together with nonlinear Green's identity ([16], p. 35), one has

$$
\begin{aligned}
& j^{\prime}\left(\bar{u}_{\lambda}\right)(h)=\frac{1}{q} \int_{\Omega} \frac{-\Delta_{p}^{a_{1}} \bar{u}_{\lambda}-\Delta_{q}^{a_{2}} \bar{u}_{\lambda}}{\bar{u}_{\lambda}^{q-1}} h d z=\int_{\Omega}\left([\lambda-\varepsilon]-c_{7} \bar{u}_{\lambda}^{p-q}\right) h d z, \\
& j^{\prime}\left(\bar{w}_{\lambda}\right)(h)=\frac{1}{q} \int_{\Omega} \frac{-\Delta_{p}^{a_{1}} \bar{w}_{\lambda}-\Delta_{q}^{a_{2}} \bar{w}_{\lambda}}{\bar{w}_{\lambda}^{q-1}} h d z=\int_{\Omega}\left([\lambda-\varepsilon]-c_{7} \bar{w}_{\lambda}^{p-q}\right) h d z .
\end{aligned}
$$

Since (3.21) is convex, then $j^{\prime}(\cdot)$ is monotone, and so

$$
0 \leq \int_{\Omega} c_{7}\left[\bar{w}_{\lambda}^{p-q}-\bar{u}_{\lambda}^{p-q}\right]\left(\bar{u}_{\lambda}^{q}-\bar{w}_{\lambda}^{q}\right) d z \leq 0,
$$

which implies that $\bar{u}_{\lambda}=\bar{w}_{\lambda}$. We conclude that (3.20) admits a unique positive solution $\bar{u}_{\lambda} \in \operatorname{int} C_{+}$. By oddness of (3.20), we deduce that it admits a unique negative solution $\bar{v}_{\lambda}=-\bar{u}_{\lambda} \in-$ int $C_{+}$.

In the sequel, we will work with:

$$
\begin{aligned}
& \mathcal{S}_{\lambda}^{+}=\left\{\text {set of positive solutions to }\left(P_{\lambda}\right)\right\}, \\
& \mathcal{S}_{\lambda}^{-}=\left\{\text {set of negative solutions to }\left(P_{\lambda}\right)\right\} .
\end{aligned}
$$

Observe (by Proposition 3.1) that if $\lambda>\widehat{\lambda}_{1}\left(q, a_{2}\right)$, then $\emptyset \neq \mathcal{S}_{\lambda}^{+} \subseteq \operatorname{int} C_{+}$and $\emptyset \neq \mathcal{S}_{\lambda}^{-} \subseteq-\operatorname{int} C_{+}$. We also mention that the unique constant sign solutions of (3.20) provide bounds for the elements of these two solution sets.

Proposition 3.3. Let $H_{0}, H_{1}$ be satisfied, and $\lambda>\widehat{\lambda}_{1}\left(q, a_{2}\right)$. Then $\bar{u}_{\lambda} \leq u$ for all $u \in \mathcal{S}_{\lambda}^{+}$and $v \leq \bar{v}_{\lambda}$ for all $v \in \mathcal{S}_{\lambda}^{-}$.

Proof. For $u \in \mathcal{S}_{\lambda}^{+} \subseteq \operatorname{int} C_{+}$and $\varepsilon \in\left(0, \lambda-\widehat{\lambda}_{1}\left(q, a_{2}\right)\right)$, we introduce a Carathéodory function $k_{\lambda}^{+}$: $\Omega \times \mathbb{R} \rightarrow \mathbb{R}$ defined by

$$
k_{\lambda}^{+}(z, x)= \begin{cases}{[\lambda-\varepsilon]\left(x^{+}\right)^{q-1}-c_{7}\left(x^{+}\right)^{p-1}} & \text { if } x \leq u(z)  \tag{3.22}\\ {[\lambda-\varepsilon] u(z)^{q-1}-c_{7} u(z)^{p-1}} & \text { if } u(z)<x\end{cases}
$$

Let $K_{\lambda}^{+}(z, x)=\int_{0}^{x} k_{\lambda}^{+}(z, s) d s$ and $\beta_{\lambda}^{+}: W_{0}^{1, p}(\Omega) \rightarrow \mathbb{R}$ be the $C^{1}$-functional

$$
\beta_{\lambda}^{+}(u)=\frac{1}{p} \int_{\Omega} a_{1}(z)|\nabla u|^{p} d z+\frac{1}{q} \int_{\Omega} a_{2}(z)|\nabla u|^{q} d z-\int_{\Omega} K_{\lambda}^{+}(z, u) d z
$$

for all $u \in W_{0}^{1, p}(\Omega)$. Now (3.22) ensures the coercivity of $\beta_{\lambda}^{+}(\cdot)$; additionally, $\beta_{\lambda}^{+}(\cdot)$ is sequentially weakly lower semicontinuous. By using the similar arguments to the ones in the proof of Proposition 3.1, one can deduce that there exists $\widetilde{u}_{\lambda} \in W_{0}^{1, p}(\Omega)$ with

$$
\begin{equation*}
\left\langle A_{p}^{a_{1}}\left(\widetilde{u}_{\lambda}\right), h\right\rangle+\left\langle A_{q}^{a_{2}}\left(\widetilde{u}_{\lambda}\right), h\right\rangle=\int_{\Omega} k_{\lambda}^{+}\left(z, \widetilde{u}_{\lambda}\right) h d z \quad \text { for all } h \in W_{0}^{1, p}(\Omega) . \tag{3.23}
\end{equation*}
$$

In (3.23) first we use $h=-\widetilde{u}_{\lambda}^{-} \in W_{0}^{1, p}(\Omega)$ leading to $\widetilde{u}_{\lambda} \geq 0, \widetilde{u}_{\lambda} \neq 0$. Next taking $h=\left(\widetilde{u}_{\lambda}-u\right)^{+} \in$ $W_{0}^{1, p}(\Omega)$, we have

$$
\begin{aligned}
& \left.\left.\left\langle A_{p}^{a_{1}} \widetilde{u}_{\lambda}\right),\left(\widetilde{u}_{\lambda}-u\right)^{+}\right\rangle+\left\langle A_{q}^{a_{2}}\left(\widetilde{u}_{\lambda}\right), \widetilde{u}_{\lambda}-u\right)^{+}\right\rangle \\
& =\int_{\Omega}\left([\lambda-\varepsilon] u^{q-1}-c_{7} u^{p-1}\right)\left(\widetilde{u}_{\lambda}-u\right)^{+} d z \quad(\text { see }(3.22)) \\
& \leq \int_{\Omega}\left(\lambda u^{q-1}+f(z, u)\right)\left(\widetilde{u}_{\lambda}-u\right)^{+} d z \quad(\text { see }(3.19)) \\
& =\left\langle A_{p}^{a_{1}}(u),\left(\widetilde{u}_{\lambda}-u\right)^{+}\right\rangle+\left\langle A_{q}^{a_{2}}(u),\left(\widetilde{u}_{\lambda}-u\right)^{+}\right\rangle \quad\left(\text { since } u \in \mathcal{S}_{\lambda}^{+}\right),
\end{aligned}
$$

which implies $\widetilde{u}_{\lambda} \leq u$. Summarizing

$$
\begin{equation*}
\widetilde{u}_{\lambda} \in[0, u], \quad \widetilde{u}_{\lambda} \neq 0 . \tag{3.24}
\end{equation*}
$$

Using (3.22), (3.24), (3.23), then $\widetilde{u}_{\lambda}$ is positive solution of (3.20). So, on account of Proposition 3.2, we have $\bar{u}_{\lambda}=\bar{u}_{\lambda}$. Therefore $\bar{u}_{\lambda} \leq u$ for all $u \in \mathcal{S}_{\lambda}^{+}$(see (3.24)). Clearly, on the similar lines, one can establish that $v \leq \bar{v}_{\lambda}$ for all $v \in \mathcal{S}_{\lambda}^{-}$.

The extremal constant sign solutions to $\left(P_{\lambda}\right)\left(\lambda>\bar{\lambda}_{1}\left(q, a_{2}\right)\right)$ are obtained as follows.
Proposition 3.4. Let $H_{0}, H_{1}$ be satisfied, and $\lambda>\widehat{\lambda}_{1}\left(q, a_{2}\right)$. Then there exist $u_{\lambda}^{*} \in \mathcal{S}_{\lambda}^{+}$and $v_{\lambda}^{*} \in \mathcal{S}_{\lambda}^{-}$ where $u_{\lambda}^{*} \leq u$ for all $u \in \mathcal{S}_{\lambda}^{+}, v \leq v_{\lambda}^{*}$ for all $v \in \mathcal{S}_{\lambda}^{-}$.

Proof. We mention that Papageorgiou et al. [15, Proposition 7] ensures that $S_{\lambda}^{+}$is downward directed (i.e., if $u_{1}, u_{2} \in \mathcal{S}_{\lambda}^{+}$, then there exists $u \in \mathcal{S}_{\lambda}^{+}$with $u \leq u_{1}, u \leq u_{2}$ ). Moreover, Hu and Papageorgiou [8, Lemma 3.10] help us to find $\left\{u_{n}\right\}_{n \in \mathbb{N}} \subseteq \mathcal{S}_{\lambda}^{+} \subseteq \operatorname{int} C_{+}$decreasing and satisfying

$$
\begin{equation*}
\inf _{n \in \mathbb{N}} u_{n}=\inf \mathcal{S}_{\lambda}^{+}, \quad \bar{u}_{\lambda} \leq u_{n} \leq u_{1} \quad \text { for all } n \in \mathbb{N} \text { (see Proposition 3.3). } \tag{3.25}
\end{equation*}
$$

Starting from

$$
\begin{equation*}
\left\langle A_{p}^{a_{1}}\left(u_{n}\right), h\right\rangle+\left\langle A_{q}^{a_{2}}\left(u_{n}\right), h\right\rangle=\int_{\Omega}\left[\lambda u_{n}^{q-1}+f\left(z, u_{n}\right)\right] h d z \quad \text { for all } h \in W_{0}^{1, p}(\Omega), \tag{3.26}
\end{equation*}
$$

and taking $h=u_{n} \in W_{0}^{1, p}(\Omega)$, then (3.25) and $H_{0}$ give us $c_{1}\left\|\nabla u_{n}\right\|_{p}^{p} \leq c_{8}$ for some $c_{8}>0$, for all $n \in \mathbb{N}$, and hence $\left\{u_{n}\right\}_{n \in \mathbb{N}} \subseteq W_{0}^{1, p}(\Omega)$ is bounded. Therefore, it is possible to suppose

$$
\begin{equation*}
u_{n} \xrightarrow{w} u_{\lambda}^{*} \text { in } W_{0}^{1, p}(\Omega), u_{n} \rightarrow u_{\lambda}^{*} \text { in } L^{p}(\Omega) . \tag{3.27}
\end{equation*}
$$

Before taking $n \rightarrow \infty$ in (3.26), we use $h=u_{n}-u_{\lambda}^{*} \in W_{0}^{1, p}(\Omega)$, and by (3.27) we get

$$
\begin{align*}
& \lim _{n \rightarrow \infty}\left[\left\langle A_{p}^{a_{1}}\left(u_{n}\right), u_{n}-u_{\lambda}^{*}\right\rangle+\left\langle A_{q}^{a_{2}}\left(u_{n}\right), u_{n}-u_{\lambda}^{*}\right\rangle\right]=0, \\
\Rightarrow \quad & \limsup _{n \rightarrow \infty}\left[\left\langle A_{p}^{a_{1}}\left(u_{n}\right), u_{n}-u_{\lambda}^{*}\right\rangle+\left\langle A_{q}^{a_{2}}\left(u_{\lambda}^{*}\right), u_{n}-u_{\lambda}^{*}\right\rangle\right] \leq 0 \quad\left(\text { since } A_{q}^{a_{2}}(\cdot) \text { is monotone }\right), \\
\Rightarrow \quad & \limsup _{n \rightarrow \infty}\left\langle A_{p}^{a_{1}}\left(u_{n}\right), u_{n}-u_{\lambda}^{*}\right\rangle \leq 0 \quad(\operatorname{see}(3.27)), \\
\Rightarrow \quad & u_{n} \rightarrow u_{\lambda}^{*} \text { in } W_{0}^{1, p}(\Omega)\left(A_{p}^{a_{1}} \text { is of type }(S)_{+}\right) . \tag{3.28}
\end{align*}
$$

Returning to Eq (3.26) and letting again $n \rightarrow \infty$, (3.28) and (3.25) lead to

$$
\begin{aligned}
& \left\langle A_{p}^{a_{1}}\left(u_{\lambda}^{*}\right), h\right\rangle+\left\langle A_{q}^{a_{2}}\left(u_{\lambda}^{*}\right), h\right\rangle=\int_{\Omega}\left[\lambda\left(u_{\lambda}^{*}\right)^{q-1}+f\left(z, u_{\lambda}^{*}\right)\right] h d z \quad \text { for all } h \in W_{0}^{1, p}(\Omega), \\
& \bar{u}_{\lambda} \leq u_{\lambda}^{*} .
\end{aligned}
$$

We arrive to the conclusion that $u_{\lambda}^{*} \in \mathcal{S}_{\lambda}^{+}$and $u_{\lambda}^{*}=\inf \mathcal{S}_{\lambda}^{+}$. Similarly, we produce $v_{\lambda}^{*} \in \mathcal{S}_{\lambda}^{-}, v_{\lambda}^{*}=$ $\sup \mathcal{S}_{\lambda}^{-}$, where $\mathcal{S}_{\lambda}^{-}$is upward directed (i.e., if $v_{1}, v_{2} \in \mathcal{S}_{\lambda}^{-}$, then there exists $v \in \mathcal{S}_{\lambda}^{-}$with $v_{1} \leq v$, $v_{2} \leq v$ ).

## 4. Nodal solutions

We implement a simple idea: we will use truncations to work over the order interval $\left[v_{\lambda}^{*}, u_{\lambda}^{*}\right]$. Any nontrivial solution ( $\equiv u_{\lambda}^{*}, v_{\lambda}^{*}$ ) of $\left(P_{\lambda}\right)$ there, will be nodal. The key ingredient is the minimax characterization of $\widehat{\lambda}_{2}\left(q, a_{2}\right)$ (see (2.4)). From Section 3 we have $u_{\lambda}^{*} \in \operatorname{int} C_{+}$and $v_{\lambda}^{*} \in-\operatorname{int} C_{+}$solving $\left(P_{\lambda}\right)\left(\lambda>\widehat{\lambda}_{1}\left(q, a_{2}\right)\right)$. Then we introduce

$$
\mu_{\lambda}(z, x)= \begin{cases}\lambda\left|v_{\lambda}^{*}(z)\right|^{q-2} v_{\lambda}^{*}(z)+f\left(z, v_{\lambda}^{*}(z)\right) & \text { if } x<v_{\lambda}^{*}(z),  \tag{4.1}\\ \lambda|x|^{q-2} x+f(z, x) & \text { if } v_{\lambda}^{*}(z) \leq x \leq u_{\lambda}^{*}(z), \\ \lambda u_{\lambda}^{*}(z)^{q-1}+f\left(z, u_{\lambda}^{*}(z)\right) & \text { if } u_{\lambda}^{*}(z)<x .\end{cases}
$$

Evidently $\mu_{\lambda}(\cdot, \cdot)$ is of Carathéodory. Additionally, we need

$$
\begin{equation*}
\mu_{\lambda}^{ \pm}(z, x)=\mu_{\lambda}\left(z, \pm x^{ \pm}\right) . \tag{4.2}
\end{equation*}
$$

Putting $M_{\lambda}(z, x)=\int_{0}^{x} \mu_{\lambda}(z, s) d s, M_{\lambda}^{ \pm}(z, x)=\int_{0}^{x} \mu_{\lambda}^{ \pm}(z, s) d s$, one can define the $C^{1}$-functionals $\widehat{\psi}_{\lambda}, \widehat{\psi}_{\lambda}^{ \pm}$: $W_{0}^{1, p}(\Omega) \rightarrow \mathbb{R}$ as

$$
\begin{aligned}
& \widehat{\psi}_{\lambda}(u)=\frac{1}{p} \int_{\Omega} a_{1}(z)|\nabla u|^{p} d z+\frac{1}{q} \int_{\Omega} a_{2}(z)|\nabla u|^{q} d z-\int_{\Omega} M_{\lambda}(z, u) d z, \\
& \widehat{\psi}_{\lambda}^{ \pm}(u)=\frac{1}{p} \int_{\Omega} a_{1}(z)|\nabla u|^{p} d z+\frac{1}{q} \int_{\Omega} a_{2}(z)|\nabla u|^{q} d z-\int_{\Omega} M_{\lambda}^{ \pm}(z, u) d z
\end{aligned}
$$

for all $u \in W_{0}^{1, p}(\Omega)$. From (4.1), (4.2), the nonlinear regularity theory and the extremality of $u_{\lambda}^{*}$ and $v_{\lambda}^{*}$, we infer easily the following result.

Proposition 4.1. Let $H_{0}, H_{1}$ be satisfied, and $\lambda>\widehat{\lambda}_{1}\left(q, a_{2}\right)$. Then, $K_{\widehat{\psi}_{\lambda}} \subseteq\left[v_{\lambda}^{*}, u_{\lambda}^{*}\right] \cap C_{0}^{1}(\bar{\Omega}), K_{\widetilde{\psi}_{\lambda}^{+}}=\left\{0, u_{\lambda}^{*}\right\}$, $K_{\widehat{\psi}_{\lambda}^{-}}=\left\{0, v_{\lambda}^{*}\right\}$.

We establish the following auxiliary proposition.
Proposition 4.2. Let $H_{0}, H_{1}$ be satisfied, and $\lambda>\widehat{\lambda}_{1}\left(q, a_{2}\right)$. Then, $u_{\lambda}^{*} \in \operatorname{int} C_{+}$and $v_{\lambda}^{*} \in-\operatorname{int} C_{+}$are local minimizers of $\widehat{\psi}_{\lambda}(\cdot)$.
Proof. Definitions (4.1) and (4.2) give us the coercivity of $\widehat{\psi}_{\lambda}^{ \pm}(\cdot)$, which are sequentially weakly lower semicontinuous too. Similarly to the proofs of previous propositions but involving $\bar{\psi}_{\lambda}^{+}(\cdot)$ this time,
there exists a certain $\widetilde{u}_{\lambda}^{*} \in W_{0}^{1, p}(\Omega)$ with $\widetilde{u}_{\lambda}^{*} \neq 0$. As $\widetilde{u}_{\lambda}^{*} \in K_{\widetilde{\psi}_{\lambda}^{+}} \backslash\{0\}$, from Proposition 4.1, we get $\widetilde{u}_{\lambda}^{*}=u_{\lambda}^{*} \in \operatorname{int} C_{+}$. Observe $\left.\widehat{\psi}_{\lambda}\right|_{C_{+}}=\left.\widehat{\psi}_{\lambda}^{+}\right|_{C_{+}}($see (4.1), (4.2)), and hence we have

$$
\begin{aligned}
& u_{\lambda}^{*} \text { is a local } C_{0}^{1}(\bar{\Omega}) \text {-minimizer of } \widehat{\psi}_{\lambda}(\cdot), \\
\Rightarrow \quad & u_{\lambda}^{*} \text { is a local } W_{0}^{1, p}(\Omega) \text {-minimizer of } \widehat{\psi}_{\lambda}(\cdot)(\text { refer to }[4]) .
\end{aligned}
$$

Involving in a similar way $\widehat{\psi}_{\lambda}^{-}(\cdot)$, we complete the proof for $v_{\lambda}^{*} \in-\operatorname{int} C_{+}$.
Using the method outlined in the beginning of this section, we establish the following.
Proposition 4.3. Let $H_{0}, H_{1}$ be satisfied, and $\lambda>\widehat{\lambda}_{2}\left(q, a_{2}\right)$. Then, $\left(P_{\lambda}\right)$ admits a nodal solution $y_{\lambda} \in\left[v_{\lambda}^{*}, u_{\lambda}^{*}\right] \cap C_{0}^{1}(\bar{\Omega})$.
Proof. To develop the reasoning here, we start from the inequality

$$
\begin{equation*}
\widehat{\psi}_{\lambda}\left(v_{\lambda}^{*}\right) \leq \widehat{\psi}_{\lambda}\left(u_{\lambda}^{*}\right), \tag{4.3}
\end{equation*}
$$

but of course we could assume equivalently $\widehat{\psi}_{\lambda}\left(v_{\lambda}^{*}\right) \geq \widehat{\psi}_{\lambda}\left(u_{\lambda}^{*}\right)$. On account of Proposition 4.1 and without any restriction, let $K_{\widehat{\psi}_{\lambda}}$ be finite (otherwise we already have an infinity of nodal smooth solutions). Proposition 4.2, (4.3) and Papageorgiou et al. [16, Theorem 5.7.6], ensure us that there is $\rho \in(0,1)$ small with

$$
\begin{equation*}
\widehat{\psi}_{\lambda}\left(v_{\lambda}^{*}\right) \leq \widehat{\psi}_{\lambda}\left(u_{\lambda}^{*}\right)<\inf \left[\widehat{\psi}_{\lambda}(u):\left\|u-u_{\lambda}^{*}\right\|=\rho\right]=\widehat{m}_{\lambda}, \quad \rho<\left\|v_{\lambda}^{*}-u_{\lambda}^{*}\right\| \quad(\text { see }(4.3)) . \tag{4.4}
\end{equation*}
$$

Again definition (4.1) gives us the coercivity of $\widehat{\psi}_{\lambda}(\cdot)$, which hence satisfies the Palais-Smale condition ([16], p. 369). This fact and (4.4) lead to a mountain pass geometry, which ensures the existence of $y_{\lambda} \in W_{0}^{1, p}(\Omega)$ with

$$
\begin{equation*}
y_{\lambda} \in K_{\widehat{\psi}_{\lambda}} \subseteq\left[v_{\lambda}^{*}, u_{\lambda}^{*}\right] \cap C_{0}^{1}(\bar{\Omega}) \quad\left(\text { see Proposition 4.1), } \quad \widehat{m}_{\lambda} \leq \widehat{\psi}_{\lambda}\left(y_{\lambda}\right) .\right. \tag{4.5}
\end{equation*}
$$

From (4.5) and (4.1) it follows that $y_{\lambda} \in C_{0}^{1}(\bar{\Omega})$ solves $\left(P_{\lambda}\right)$ and it is distinct from $u_{\lambda}^{*}, v_{\lambda}^{*}$. To conclude, it remains to prove that $y_{\lambda} \neq 0$. Mountain pass theorem ensures that

$$
\widehat{\psi}_{\lambda}\left(y_{n}\right)=\inf _{\gamma \in \Gamma} \max _{-1 \leq t \leq 1} \widehat{\psi}_{\lambda}(\gamma(t)),
$$

with $\Gamma=\left\{\gamma \in C\left([-1,1], W_{0}^{1, p}(\Omega)\right): \gamma(-1)=v_{\lambda}^{*}, \gamma(1)=u_{\lambda}^{*}\right\}$. We consider the following Banach manifolds $M=W_{0}^{1, p}(\Omega) \cap \partial B_{1}^{L^{q}}, M_{c}=M \cap C_{0}^{1}(\bar{\Omega})$, where $\partial B_{1}^{L^{q}}=\left\{u \in L^{q}(\Omega):\|u\|_{q}=1\right\}$ and we introduce the sets of paths:

$$
\begin{aligned}
& \widehat{\Gamma}=\left\{\widehat{\gamma} \in C([-1,1], M): \widehat{\gamma}(-1)=-\widehat{u}_{1}\left(q, a_{2}\right), \widehat{\gamma}(1)=\widehat{u}_{1}\left(q, a_{2}\right)\right\}, \\
& \widehat{\Gamma}_{c}=\left\{\widehat{\gamma} \in C\left([-1,1], M_{c}\right): \widehat{\gamma}(-1)=-\widehat{u}_{1}\left(q, a_{2}\right), \widehat{\gamma}(1)=\widehat{u}_{1}\left(q, a_{2}\right)\right\} .
\end{aligned}
$$

Claim: $\widehat{\Gamma}_{c}$ is dense in $\widehat{\Gamma}$.
Given $\widehat{\gamma} \in \widehat{\Gamma}$ and $\varepsilon \in(0,1)$, we introduce $\widehat{K}_{\varepsilon}:[-1,1] \rightarrow 2^{C_{0}^{1}(\bar{\Omega})}$ of the form

$$
\widehat{K}_{\varepsilon}(t)= \begin{cases}\left\{u \in C^{1}(\bar{\Omega}):\|u-\widehat{\gamma}(t)\|<\varepsilon\right\} & \text { if }-1<t<1, \\ \left\{ \pm \widehat{u}_{1}\left(q, a_{2}\right)\right\} & \text { if } t= \pm 1 .\end{cases}
$$

This multifunction has nonempty and convex values. Additionally, for $t \in(-1,1) \widehat{K}_{\varepsilon}(t)$ is open, while the sets $\widehat{K}_{\varepsilon}(1), \widehat{K}_{\varepsilon}(-1)$ are singletons. Now, Hu and Papageorgiou [8, Proposition 2.6], implies that $\widehat{K}_{\varepsilon}(\cdot)$ is lsc, and hence Michael [13, Theorem 3.1 $1^{\prime \prime \prime}$ ] ensures the existence of a continuous map $\widehat{\gamma}_{\varepsilon}:[-1,1] \rightarrow C_{0}^{1}(\bar{\Omega})$ with $\widehat{\gamma}_{\varepsilon}(t) \in \widehat{K}_{\varepsilon}(t)$ for all $t \in[-1,1]$.

Put $\varepsilon=n^{-1}, n \in \mathbb{N}$ and let $\widehat{\gamma}_{n}=\widehat{\gamma}_{\frac{1}{n}}$ be the continuous selection of the multifunction $\widehat{K}_{\frac{1}{n}}(\cdot)$ produced above. The inequality

$$
\begin{equation*}
\left\|\widehat{\gamma}_{n}(t)-\widehat{\gamma}(t)\right\|<\frac{1}{n} \quad \text { for all } t \in[-1,1], \tag{4.6}
\end{equation*}
$$

holds and since $\widehat{\gamma} \in \widehat{\Gamma}$, we see that $\|\widehat{\gamma}(t)\| \geq m>0$ for all $t \in[-1,1]$. Hence (4.6) leads us to suppose $\left\|\widehat{\gamma}_{n}(t)\right\| \neq 0$ for all $t \in[-1,1]$, all $n \in \mathbb{N}$. We set $\widetilde{\gamma}_{n}(t)=\frac{\widehat{\gamma}_{n}(t)}{\left\|\widehat{\gamma}_{n}(t)\right\|_{q}}$ for all $t \in[-1,1]$, all $n \in \mathbb{N}$. Then we have $\widetilde{\gamma}_{n} \in C\left([-1,1], M_{c}\right), \widetilde{\gamma}_{n}( \pm 1)= \pm \widehat{u}_{1}\left(q, a_{2}\right)$. Moreover,

$$
\begin{align*}
\left\|\widehat{\gamma}_{n}(t)-\widehat{\gamma}(t)\right\| & \leq\left\|\widehat{\gamma}_{n}(t)-\widehat{\gamma}_{n}(t)\right\|+\left\|\widehat{\gamma}_{n}(t)-\widehat{\gamma}(t)\right\| \\
& \leq \frac{\left|1-\left\|\widehat{\gamma}_{n}(t)\right\|_{q}\right|}{\left\|\widehat{\gamma}_{n}(t)\right\|_{q}}\left\|\widehat{\gamma}_{n}(t)\right\|+\frac{1}{n} \quad \text { for all } t \in[-1,1], \text { all } n \in \mathbb{N} \text { (see (4.6)). } \tag{4.7}
\end{align*}
$$

Note that

$$
\begin{aligned}
\max _{-1 \leq t \leq 1}\left|1-\left\|\widehat{\gamma}_{n}(t)\right\|_{q}\right| & =\max _{-1 \leq \leq \leq 1}\|\widehat{\gamma}(t)\|_{q}-\left\|\widehat{\gamma}_{n}(t)\right\|_{q} \mid \quad(\text { since } \widehat{\gamma} \in \widehat{\Gamma}) \\
& \leq \max _{-1 \leq \leq 1}\left\|\widehat{\gamma}(t)-\widehat{\gamma}_{n}(t)\right\|_{q} \\
& \leq c_{9} \max _{-1 \leq \leq 1} \mid \widehat{\gamma}(t)-\widehat{\gamma}_{n}(t) \| \quad \text { for some } c_{9}>0\left(W_{0}^{1, q}(\Omega) \hookrightarrow L^{q}(\Omega)\right) \\
& \leq \frac{c_{9}}{n} \quad(\operatorname{see}(4.6)) .
\end{aligned}
$$

We use this estimate in (4.7), together with (4.6) and the fact that $W_{0}^{1, p}(\Omega) \hookrightarrow L^{q}(\Omega)$. We obtain

$$
\left\|\widetilde{\gamma}_{n}(t)-\widehat{\gamma}(t)\right\| \leq \frac{c_{9}}{n c_{10}-1}\left[1+\frac{1}{n}\right]+\frac{1}{n} \quad \text { for some } c_{10}>0, \text { all } n \in \mathbb{N},
$$

which implies that $\widehat{\Gamma}_{c}$ is dense in $\widehat{\Gamma}$. Using this and (2.4), one can find $\widehat{\gamma} \in \widehat{\Gamma}_{c}$ satisfying

$$
\int_{\Omega} a_{2}(z)|\nabla \widehat{\gamma}(t)|^{q} d z<\widehat{\lambda}_{2}\left(q, a_{2}\right)+\vartheta \quad \text { for all } t \in[-1,1], \text { with } 0<\vartheta<\frac{1}{2}\left(\lambda-\widehat{\lambda}_{2}\left(q, a_{2}\right)\right) .
$$

Next, $H_{1}(i v)$ ensures the existence of $\delta>0$ satisfying

$$
\begin{equation*}
F(z, x) \geq-\frac{\vartheta}{q}|x|^{q} \quad \text { for a.a. } z \in \Omega \text {, all }|x| \leq \delta . \tag{4.8}
\end{equation*}
$$

We have the compactness of $\widehat{\gamma}([-1,1]) \subseteq M_{c}$, and we know that $u_{\lambda}^{*} \in \operatorname{int} C_{+}$and $v_{\lambda}^{*} \in-\operatorname{int} C_{+}$. Now, by Papageorgiou et al. [16, Proposition 4.1.24], we can find $\xi \in(0,1)$ small with

$$
\begin{align*}
& \xi \widehat{\gamma}(t) \in\left[v_{\lambda}^{*}, u_{\lambda}^{*}\right] \cap C_{0}^{1}(\bar{\Omega}) \quad \text { for all } t \in[-1,1], \\
& |\xi \widehat{\gamma}(t)(z)| \leq \delta \quad \text { for all } t \in[-1,1], \text { all } z \in \bar{\Omega} . \tag{4.9}
\end{align*}
$$

Consider $u \in \xi \widehat{\gamma}([-1,1])$. Therefore $u=\xi \widehat{u}$ with $\widehat{u} \in \widehat{\gamma}([-1,1])$. We have

$$
\begin{aligned}
\widehat{\psi}_{\lambda}(u) & \leq \frac{\xi^{p}}{p} \int_{\Omega} a_{1}(z)|\nabla \hat{u}|^{p} d z+\frac{\xi^{q}}{q}\left[\int_{\Omega} a_{2}(z)|\nabla \hat{u}|^{q} d z-(\lambda-\vartheta)\right] \\
& \left.\quad \text { (see (4.8), (4.9) and recall }\|\widehat{\gamma}(t)\|_{q}=1\right) \\
& \leq \frac{\xi^{p}}{p} \int_{\Omega} a_{1}(z)|\nabla \widehat{u}|^{p} d z-\frac{\xi^{q}}{q}\left[\lambda-\left(\widehat{\lambda}_{2}\left(q, a_{2}\right)+2 \vartheta\right)\right] \quad \text { (see again (4.8), (4.9)) } \\
& \left.\leq c_{11} \xi^{p}-c_{12} \xi^{q} \quad \text { for some } c_{11}, c_{12}>0 \text { (recall the choice of } \vartheta\right) .
\end{aligned}
$$

Then choosing $\xi \in(0,1)$ (smaller enough), one has

$$
\begin{equation*}
\left.\widehat{\psi}_{\lambda}\right|_{\gamma_{0}}<0 \quad \text { where } \gamma_{0}=\xi \widehat{\gamma} . \tag{4.10}
\end{equation*}
$$

Let $a=\widehat{\psi}_{\lambda}^{+}\left(u_{\lambda}^{*}\right)=\widehat{\psi}_{\lambda}\left(u_{\lambda}^{*}\right)$ and $b=0=\widehat{\psi}_{\lambda}^{+}(0)=\widehat{\psi}_{\lambda}(0)$. From the proof of Proposition 4.2, we know that $a<b=0$. Moreover on account of Proposition 4.1 and since $u_{\lambda}^{*}$ is the global minimizer of $\widehat{\psi}_{\lambda}^{+}$, one can conclude that $K_{\widehat{\psi}_{\lambda}^{+}}^{a}=\left\{u_{\lambda}^{*}\right\}, \widehat{\psi}_{\lambda}^{+}\left(K_{\widehat{\psi}_{\lambda}^{+}}\right) \cap(a, 0)=\emptyset$.

Therefore we can apply the second deformation theorem in Papageorgiou et al. [16] (p. 386) and produce $h_{0}:[0,1] \times\left(\left(\widehat{\psi}_{\lambda}^{+}\right)^{0} \backslash K_{\widetilde{\psi}_{\lambda}^{+}}^{0}\right) \rightarrow\left(\widehat{\psi}_{\lambda}^{+}\right)^{a}$ such that

$$
\begin{align*}
& h_{0}(0, u)=u \quad \text { for all } u \in\left(\left(\widehat{\psi}_{\lambda}^{+}\right)^{0} \backslash\{0\}\right)\left(\text { note } K_{\psi_{\lambda}^{+}}^{0}=\{0\}\right),  \tag{4.11}\\
& \left.h_{0}(t, u)=u_{\lambda}^{*} \quad \text { for all } u \in\left(\left(\widehat{\psi}_{\lambda}^{+}\right)^{0} \backslash\{0\}\right) \text {, all } t \in[0,1] \text { (note } K_{\widehat{\psi}_{\lambda}^{+}}^{a}=\left\{u_{\lambda}^{*}\right\}\right),  \tag{4.12}\\
& \widehat{\psi}_{\lambda}^{+}\left(h_{0}(t, u)\right) \leq \widehat{\psi}_{\lambda}^{+}\left(h_{0}(s, u)\right) \quad \text { for all } 0 \leq s \leq t \leq 1, \text { all } u \in\left(\left(\widehat{\psi}_{\lambda}^{+}\right)^{0} \backslash\{0\}\right) . \tag{4.13}
\end{align*}
$$

These properties of the deformation $h_{0}$ imply that $K_{\widehat{\psi}_{\lambda}^{+}}^{a}$ is a strong deformation retract of $\left(\widehat{\psi}_{\lambda}^{+}\right)^{0} \backslash$ $\{0\}$ and the deformation is $\widehat{\psi}_{\lambda}^{+}$-decreasing. We set $\gamma_{+}(t)=h_{0}\left(t, \xi \widehat{u}_{1}\left(q, a_{2}\right)\right)^{+}$for all $t \in[0,1]$, i.e., a continuous path in $W_{0}^{1, p}(\Omega)$ and its trace is in the positive cone of $W_{0}^{1, p}(\Omega)$. Note $\xi \widehat{u}_{1}\left(q, a_{2}\right) \in\left(\widehat{\psi}_{\lambda}^{+}\right)^{0}$ (see (4.10)) and $\widehat{\psi}_{\lambda}^{+}\left(\xi \widehat{u}_{1}\left(q, a_{2}\right)\right)=\widehat{\psi}_{\lambda}\left(\xi \widehat{u}_{1}\left(q, a_{2}\right)\right)$. So, we have

$$
\begin{align*}
& \gamma_{+}(0)=\xi \widehat{u}_{1}\left(q, a_{2}\right) \quad(\operatorname{see}(4.11)), \\
& \gamma_{+}(1)=u_{\lambda}^{*} \quad(\text { see }(4.12)), \\
& \widehat{\psi}_{\lambda}^{+}\left(\gamma_{+}(t)\right) \leq \widehat{\psi}_{\lambda}^{+}\left(\gamma_{+}(0)\right) \quad \text { for all } t \in[0,1](\text { see }(4.13)), \\
\Rightarrow & \widehat{\psi}_{\lambda}\left(\gamma_{+}(t)\right) \leq \widehat{\psi}_{\lambda}\left(\xi \widehat{u}_{1}\left(q, a_{2}\right)\right)<0 \quad \text { for all } t \in[0,1] \text { (see (4.2), (4.10)) }, \\
\Rightarrow & \left.\widehat{\psi}_{\lambda}\right|_{\gamma_{+}}<0, \tag{4.14}
\end{align*}
$$

with $\gamma_{+}$being a continuous path in $W_{0}^{1, p}(\Omega)$, linking $\xi \widehat{u}_{1}\left(q, a_{2}\right)$ to $u_{\lambda}^{*}$. For $\widehat{\psi}_{\lambda}^{-}$, we can produce in a similar way a continuous path $\gamma_{-}$in $W_{0}^{1, p}(\Omega)$, connecting $-\xi \widehat{u}_{1}\left(q, a_{2}\right)$ and $v_{\lambda}^{*}$. and such that

$$
\begin{equation*}
\left.\widehat{\psi}_{\lambda}\right|_{\gamma_{-}}<0 . \tag{4.15}
\end{equation*}
$$

Merging $\gamma_{-}, \gamma_{0}, \gamma_{+}$, we get $\gamma_{*} \in \Gamma$ satisfying

$$
\left.\widehat{\psi}_{\lambda}\right|_{\gamma_{*}}<0 \quad(\operatorname{see}(4.10),(4.14),(4.15)),
$$

$$
\Rightarrow \quad \widehat{\psi}_{\lambda}\left(y_{\lambda}\right)<0=\widehat{\psi}_{\lambda}(0)
$$

which implies $y_{\lambda} \neq 0$, and so $y_{\lambda} \in\left[v_{\lambda}^{*}, u_{\lambda}^{*}\right] \cap C_{0}^{1}(\bar{\Omega})$ is nodal solution to $\left(P_{\lambda}\right)$.

So, we have the following multiplicity result of $\left(P_{\lambda}\right)$. We emphasize that in this theorem, one has sign information for all the solutions and the solutions are ordered.

Theorem 4.1. Let $H_{0}, H_{1}$ be satisfied. Thus:
(a) if $\lambda>\widehat{\lambda}_{1}\left(q, a_{2}\right)$, then $\left(P_{\lambda}\right)$ admits at least two constant sign solutions $u_{\lambda} \in \operatorname{int} C_{+}, v_{\lambda} \in-\operatorname{int} C_{+}$;
(b) if $\lambda>\bar{\lambda}_{2}\left(q, a_{2}\right)$, then there is also a nodal solution of $\left(P_{\lambda}\right)$, namely $y_{\lambda} \in\left[v_{\lambda}, u_{\lambda}\right] \cap C_{0}^{1}(\bar{\Omega})$.

If $q=2$ (weighted ( $p, 2$ )-equation), then we can improve a little Theorem 4.1(b).
Theorem 4.2. Let $H_{0}, H_{1}^{\prime}($ with $q=2)$ be satisfied, and $\lambda>\widehat{\lambda}_{2}\left(2, a_{2}\right)$. Then, $\left(P_{\lambda}\right)($ with $q=2)$ admits at least three nontrivial smooth solutions with sign information and ordered $u_{\lambda} \in \operatorname{int} C_{+}, v_{\lambda} \in-\operatorname{int} C_{+}$, $y_{\lambda} \in \operatorname{int}_{C_{0}^{1}(\bar{\Omega})}\left[v_{\lambda}, u_{\lambda}\right]$.
Proof. We start from the solutions provided by Theorem 4.1, namely $u_{\lambda} \in \operatorname{int} C_{+}, v_{\lambda} \in-\operatorname{int} C_{+}$and $y_{\lambda} \in\left[v_{\lambda}, u_{\lambda}\right] \cap C_{0}^{1}(\bar{\Omega})$ nodal.

Let $a(z, y)=a_{1}(z)|y|^{p-2}+a_{2}(z) y$ for all $z \in \Omega$, all $y \in \mathbb{R}^{N}$. Thus $\operatorname{div} a(z, \nabla u)=\Delta_{p}^{a_{1}} u+\Delta^{a_{2}} u$ for all $u \in W_{0}^{1, p}(\Omega)$. Observe $a(z, \cdot) \in C^{1}\left(\mathbb{R}^{N}, \mathbb{R}^{N}\right)$ (recall that $2<p$ here) and

$$
\begin{aligned}
& \nabla_{y} a(z, y)=a_{1}(z)|y|^{p-2}\left[\text { id }+(p-2) \frac{y \otimes y}{|y|^{2}}\right]+a_{2}(z) \mathrm{id} \\
\Rightarrow \quad & \left(\nabla_{y} a(z, y) \xi, \xi\right) \geq c_{1}|\xi|^{2} \quad \text { for all } y, \xi \in \mathbb{R}^{N} .
\end{aligned}
$$

Also, if $\rho=\max \left\{\left\|v_{\lambda}\right\|_{\infty},\left\|u_{\lambda}\right\|_{\infty}\right\}$ and $\widehat{\xi}_{\rho}>0$ is taken from $H_{1}^{\prime}(v i)$, then

$$
f(z, x)-f(z, u) \geq-\widehat{\xi}_{\rho}|x-u| \quad \text { for all } x, u \in[-\rho, \rho] .
$$

The tangency principle (Pucci and Serrin [25, Theorem 2.5.2]) leads to

$$
\begin{equation*}
v_{\lambda}(z)<y_{\lambda}(z)<u_{\lambda}(z) \quad \text { for all } z \in \Omega . \tag{4.16}
\end{equation*}
$$

Then we have

$$
\begin{align*}
-\Delta_{p}^{a_{1}} y_{\lambda}-\Delta^{a_{2}} y_{\lambda}+\widehat{\xi}_{\rho}\left|y_{\lambda}\right|^{p-2} y_{\lambda} & =\lambda y_{\lambda}+f\left(z, y_{\lambda}\right)+\widehat{\xi}_{\rho}\left|y_{\lambda}\right|^{p-2} y_{\lambda} \\
& \leq \lambda u_{\lambda}+f\left(z, u_{\lambda}\right)+\widehat{\xi}_{\rho} u_{\lambda}^{p-1} \quad\left(\text { see }(4.16) \text { and } H_{1}^{\prime}(v i)\right) \\
& =-\Delta_{p}^{a_{1}} u_{\lambda}-\Delta^{a_{2}} u_{\lambda}+\widehat{\xi}_{\rho} u_{\lambda}^{p-1} . \tag{4.17}
\end{align*}
$$

On account of (4.16) we have $0 \leq \lambda\left[u_{\lambda}-y_{\lambda}\right]$. Returning to (4.17), we obtain $u_{\lambda}-y_{\lambda} \in \operatorname{int} C_{+}$(by Gasiński et al. [6, Proposition 3.2]). On the other side, one can establish that $y_{\lambda}-v_{\lambda} \in \operatorname{int} C_{+}$. We deduce that $y_{\lambda} \in \operatorname{int}_{C_{0}^{1}(\bar{\Omega})}\left[v_{\lambda}, u_{\lambda}\right]$.

Finally under assumption $H_{1}^{\prime \prime}$ we can have a nonexistence result.

Theorem 4.3. Let $H_{0}$, $H_{1}^{\prime \prime}$ be satisfied, and $\lambda<\widehat{\lambda}_{1}\left(q, a_{2}\right)$. Then, $\left(P_{\lambda}\right)$ does not admit nontrivial solution.
Proof. At the beginning we postulate the existence of $u \in \mathcal{S}_{\lambda}^{+} \subseteq \operatorname{int} C_{+}$so that

$$
\left\langle A_{p}^{a_{1}}(u), h\right\rangle+\left\langle A_{q}^{a_{2}}(u), h\right\rangle=\int_{\Omega}\left[\lambda|u|^{q-2} u+f(z, u)\right] h d z \quad \text { for all } h \in W_{0}^{1, p}(\Omega) .
$$

For $h=u \in W_{0}^{1, p}(\Omega)$, by $H_{1}^{\prime \prime}(v i)$ we deduce that

$$
\int_{\Omega} a_{1}(z)|\nabla u|^{p} d z-\widehat{\lambda}_{1}\left(p, a_{1}\right)\|u\|_{p}^{p}+\int_{\Omega} a_{2}(z)|\nabla u|^{q} d z-\lambda\|u\|_{q}^{q} \leq 0,
$$

which implies $\left[\widehat{\lambda}_{1}\left(q, a_{2}\right)-\lambda\right]\|u\|_{q}^{q} \leq 0$, a contradiction since $\lambda<\widehat{\lambda}_{1}\left(q, a_{2}\right)$. Therefore $\mathcal{S}_{\lambda}^{+}=\emptyset$ for all $\lambda<\widehat{\lambda}_{1}\left(q, a_{2}\right)$.

Remark 4.1. For $(p, q)$-equations with no weights but with variable exponents we refer to the survey paper of Rădulescu [26].

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## Conflict of interest

The authors declare that there are no conflicts of interest regarding the publication of this paper.

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