



Research article

Time and frequency responses of non-integer order RLC circuits

Mehmet Emir Koksals*

Department of Mathematics, Ondokuz Mayıs University, 55139 Atakum Samsun, Turkey

* **Correspondence:** Email: mekoksals@omu.edu.tr.

Abstract: By introducing an auxiliary parameter, the dynamic of RLC electrical circuits of non-integer order is described by a fractional order differential equation. The order of derivative in the component models is assumed to be $0 < \gamma \leq 1$. The time and frequency domain characteristics of the circuit is investigated, and it is shown that three different filter characteristics of low-pass, high-pass and band-pass filters are obtained. The filter parameters are determined analytically, and the results are verified numerically.

Keywords: fractional calculus; electrical filters; RLC circuits; frequency response; step response; Mittag-Leffler function

Mathematics Subject Classification: 33E12, 34A08, 93E11, 94C05, 97Mxx

1. Introduction

Many of the physical systems show intrinsic fractional order behavior [1, 2–6]. Therefore, fractional calculus provides more accurate models than classical calculus for such systems [7, 8]. A significant benefit of fractional modelling appears for systems where hereditary and memory behaviors are involved since fractional derivative takes into account the past history as well [3, 9, 10]. Another beneficial use appears for the analysis of porous and/or self-similar structures where theory of fractals are involved [10, 11].

Like many of the physical systems, some electrical systems are also better modelled by fractional calculus either due to their structures (such as tree structures, domino ladders, etc.) or element behaviors (coils, capacitors, memristors) [9, 12–15].

A systematic way of constructing fractional differential equation for physical systems consists in to analyze the dimensionality and to bring the fractional derivative operator in consistency [16]. As the generalization of this work, F. Gomez, J. Rosales, and M. Guia have obtained an analytical time domain solution in terms of the Mittag-Leffler function for the fractional RLC circuit [17]. This circuit is not externally excited but it is forced by the initial charge on the capacitor. They have not studied

the forced solution of the circuit when it is excited by an external source but commented about the use of the circuit as fractional filters analogical to ordinary second order electrical RLC filters; the element behavior equations are also missing.

The purpose of this contribution is to investigate the time and frequency domain responses of the RLC circuit proposed by F. Gomez et al. under the excitement by an external voltage source; the frequency responses obtained show that the circuit can be used as fractional low-pass, high-pass and band-pass filters. Section II is devoted to the introduction of the fractional RLC circuit and its natural response. Section III covers its frequency response characteristics and its usage as three basic types of electrical filters. The sinusoidal responses of the obtained filters are studied in Section IV. Finally, Section V covers the conclusions.

2. Fractional RLC circuit; Natural response

It is well known that a series RLC circuit is described by the homogenous differential equation

$$L \frac{d^2 q(t)}{dt^2} + R \frac{dq(t)}{dt} + \frac{q(t)}{C} = 0 \quad (2.1)$$

where R, L, C are the resistance, inductance, and capacitance in ohms (Ω), henries (H), and farads (F), respectively; $q(t)$ is the charge circulating in the circuit and accumulated on the capacitor. Each term in the above formula represents the potential across the associated component so that the total potential in the circuit sums to zero [18].

A systematic way to construct the fractional differential equation from Eq. (2.1) is to replace the derivative operator by fractional operator as follows:

$$\frac{d}{dt} \rightarrow \frac{1}{\sigma^{1-\gamma}} \frac{d^\gamma}{dt^\gamma}, \quad (2.2)$$

where the parameter σ has the dimension of seconds and $0 < \gamma \leq 1$ represents the order of Caputo fractional time derivative [16, 17] defined by

$$\frac{d^\gamma}{dt^\gamma} f(t) = {}_0^C D_t^\gamma f(t) = \frac{1}{\Gamma(n-\gamma)} \int_0^t \frac{f^n(\tau)}{(t-\tau)^{\gamma-n+1}} d\tau$$

where $\tau \in R, n-1 < \gamma \leq n \in N = \{1, 2, \dots\}$, f^n represents the ordinary derivative of order n , and Γ is the Gamma function. For $\Gamma = 1$, Eq. (2.2) gives the usual derivative. Result of application of Eq. (2.2) to Eq. (2.1) yields the fractional differential equation

$$\frac{L}{\sigma^{2(1-\gamma)}} \frac{d^{2\gamma} q}{dt^{2\gamma}} + \frac{R}{\sigma^{1-\gamma}} \frac{d^\gamma q}{dt^\gamma} + \frac{q(t)}{C} = v_s(t), \quad (2.3)$$

where $v_s(t)$ is the added voltage source excitation in the circuit as shown in Figure 1. The voltage across each component and hence the element behavior equations in the circuit are represented by the following equations:

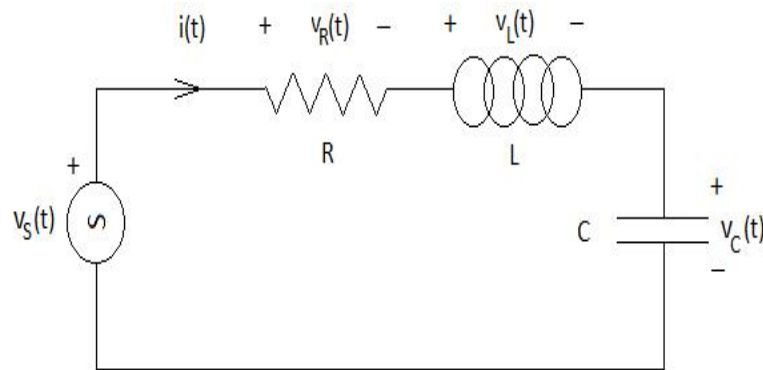


Figure 1. Fractional order RLC circuit excited by voltage source $v_s(t)$.

$$v_C(t) = \frac{q(t)}{C}, \quad (2.4)$$

$$v_R(t) = Ri(t), \quad i(t) = \frac{1}{\sigma^{1-\gamma}} \frac{d^\gamma q}{dt^\gamma}, \quad (2.5)$$

$$v_L(t) = L \frac{1}{\sigma^{1-\gamma}} \frac{d^\gamma i(t)}{dt^\gamma}. \quad (2.6)$$

The homogeneous solution of the fractional differential equation in (2.3) with the initial condition $q(t) = q_0$ is obtained by F. Gomez et al. [17]. But there are some conflicts in their solutions and plots. To be sure, the correct solution is derived for the underdamped case $R < 2\sqrt{L/C}$, the result is

$$q(t) = q_0 \operatorname{Im} \left[\frac{t^{-\gamma}}{w_0 \sigma^{1-\gamma} \sqrt{1-\xi^2}} E_{\gamma, 1-\gamma}(\lambda_1 t^\gamma) + \frac{2\xi}{\sqrt{1-\xi^2}} \right] E_{\gamma, 1}(\lambda_1 t^\gamma), \quad (2.7)$$

where

$$w_0 = 1/\sqrt{LC}, \quad (2.8)$$

$$\xi = R/2Lw_0, \quad (2.9)$$

$$\lambda_1 = w_0 \sigma^{1-\gamma} \left(-\xi + j\sqrt{1-\xi^2} \right), \quad (2.10)$$

$$\gamma = w_0 \sqrt{|1-\xi^2|} \sigma. \quad (2.11)$$

In these formulas, w_0 is the undamped natural frequency, ξ is the damping ratio, λ_1 is a zero of the characteristic polynomial in powers of s^γ , and $E_{\alpha, \beta}(\bullet)$ is the 2-parameter Mittag-Leffler function.

The solutions show that the oscillatory damping occurs for values of γ near to 1. For values of γ near to zero, the solutions monotonically damp to 0 as $t \rightarrow \infty$.

3. Fractional RLC circuit; Forced response

Although analytical and explicit solutions are given for the natural response in [17] and corrected in (2.7–2.11), these results do not give direct information about the behavior of the RLC circuit of non-integer order in respect to its time and frequency domain characteristics, that is they do not give information about the input-output relation of the circuit since it is unexcited by an external forcing source. In this section, the frequency domain characteristics (gain and phase) and the time domain characteristics for step responses are obtained, and the relevant filtering property of the circuit excited by a voltage source is investigated.

Taking the Laplace transform of the equations in (2.3) and (2.4–2.6), and solving them for $V_R(s), V_L(s), V_C(s)$, where the uppercase lettered signals represent the Laplace transform of the lowercase lettered signals, we obtain the following transfer functions in the complex frequency variable $s = \sigma + j\omega$:

$$H_{BP} = \frac{V_R}{V_S} = \frac{\frac{R\sigma^{1-\gamma}}{L} s^\gamma}{s^{2\gamma} + \frac{R\sigma^{1-\gamma}}{L} s^\gamma + \frac{\sigma^{2(1-\gamma)}}{LC}}, \quad (3.1)$$

$$H_{HP} = \frac{V_L}{V_S} = \frac{s^{2\gamma}}{s^{2\gamma} + \frac{R\sigma^{1-\gamma}}{L} s^\gamma + \frac{\sigma^{2(1-\gamma)}}{LC}}, \quad (3.2)$$

$$H_{LP} = \frac{V_C}{V_S} = \frac{\frac{\sigma^{2(1-\gamma)}}{LC}}{s^{2\gamma} + \frac{R\sigma^{1-\gamma}}{L} s^\gamma + \frac{\sigma^{2(1-\gamma)}}{LC}}. \quad (3.3)$$

Obviously, these are the fractional transfer functions resulting with the band pass, high pass and low pass filter characteristics, respectively. Gain (M) and phase (θ) characteristics are obtained by substituting $s = j\omega$ and taking the magnitude and phase angle of the transfer functions. The results are:

$$M_{BP} = \frac{R\sigma^{1-\gamma} w^\gamma}{L |\Delta|}, \theta_{BP} = \frac{\pi}{2} \gamma - \text{Arg} \tan \frac{B}{A}, \quad (3.4)$$

$$M_{HP} = \frac{w^{2\gamma}}{|\Delta|}, \theta_{HP} = \pi \gamma - \text{Arg} \tan \frac{B}{A}, \quad (3.5)$$

$$M_{LP} = \frac{\sigma^{2(1-\gamma)}}{LC |\Delta|}, \theta_{LP} = -\text{Arg} \tan \frac{B}{A}, \quad (3.6)$$

where

$$A = w^{2\gamma} \cos(\pi\gamma) + \frac{R\sigma^{1-\gamma}}{L} w^\gamma \cos\left(\frac{\pi}{2}\gamma\right) + \frac{\sigma^{2(1-\gamma)}}{LC},$$

$$B = w^{2\gamma} \sin(\pi\gamma) + w^\gamma \sin\left(\frac{\pi}{2}\gamma\right),$$

$$|\Delta| = |A + jB| = \sqrt{A^2 + B^2}.$$

As the numerical example for the underdamped case, consider $R = 0.3922327 \text{ } \Omega, L = 1 \text{ H}, C = 1 \text{ F}$, which corresponds to the case $R/2L \sqrt{w_0^2 - \xi^2} = 1/5$ as in [17], and $\omega_0 = 1 \text{ r/s}, \xi = 0.196116, \gamma = 0.980581 \sigma$. The frequency response curves (magnitude and phase) and the step responses are obtained and plotted in Figures 2–7, for the band pass, high pass and low pass cases. The plots are

carried for values of $\gamma = 1.00, 0.92, 0.75, 0.25$. The overdamped case is not considered as the numerical example since all the three types of basic filter characteristics valid for underdamped case do not occur for the overdamped case.

The band pass characteristics plotted in Figure 2 show that the peak value of the magnitude, the center frequency, and the quality factor decrease with decreasing γ . The data including the cutoff frequencies are shown in Table 1 where ω_p : Center frequency, M_p : Peak magnitude, ω_1, ω_2 : Cut off frequencies, BW : Band width, $Q = \omega_p/BW$: Quality factor.

For $\gamma = 0.25$, it is hard to say that the characteristic is of band-pass type since Q is very small. The phase responses in the same figure is seen to change from 90° to -90° but flatter with decreasing γ .

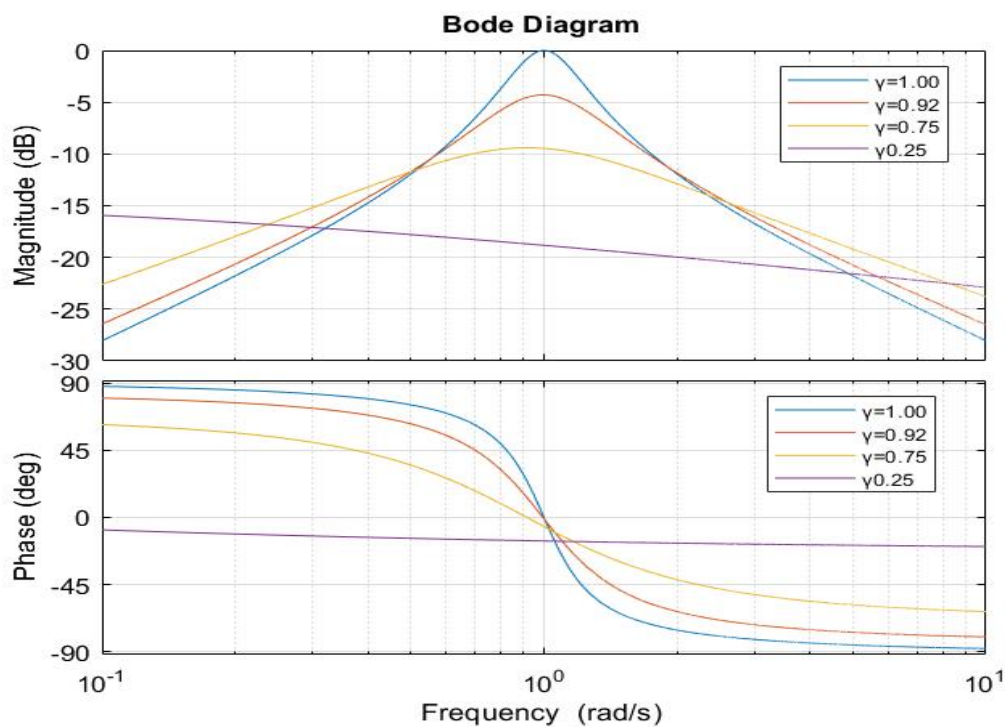


Figure 2. Frequency response of the band pass RLC filter of fractional order.

Table 1. The frequency response characteristics of band pass filter.

$\gamma \backslash Char$	ω_p (r/s)	M_p (dB)	ω_1 (r/s)	ω_2 (r/s)	BW (r/s)	Q
1.00	1	0	0.823	1.215	0.391	2.556
0.92	0.995	-4.292	0.709	1.396	0.687	1.447
0.75	0.915	-9.400	0.450	1.861	1.411	0.648
0.25	0.017	-15.13	0.001	0.635	0.634	0.026

The step responses of the filters are obtained by multiplying the transfer functions in Eqs. (3.7),

(3.8), (3.9) by $1/s$ and taking their inverse Laplace transforms. The results for *BP*, *HP*, *LP* filters are obtained, respectively, as

$$y_{BP_s}(t) = \frac{2\xi}{\sqrt{1-\xi^2}} \operatorname{Im} \left[E_{\gamma,1}(\lambda_1 t^\gamma) \right], \quad (3.7)$$

$$y_{HP_s}(t) = \frac{t^{1-\gamma}}{w_0 \sigma^{1-\gamma} \sqrt{1-\xi^2}} \operatorname{Im} \left[E_{\gamma,1-\gamma}(\lambda_1 t^\gamma) \right], \quad (3.8)$$

$$y_{LP_s}(t) = \frac{w_0 \sigma^{1-\gamma} t^\gamma}{\sqrt{1-\xi^2}} \operatorname{Im} \left[E_{\gamma,1}(\lambda_1 t^\gamma) \right]. \quad (3.9)$$

The step response of the band pass filter is plotted in Figure 3 for values of $\gamma = 1.00, 0.92, 0.75, 0.25$. It is observed that all step responses start from 0 and after several peaks damp to 0. Number of peaks and their magnitudes (M_{p1}, M_{p2}, \dots) decrease as γ decreases from 1 to 0.25. The data including characteristic time values are shown in Table 2. Note that the rise time and the settling time are not defined for reference value is not defined for *BP* characteristics. It is important to remark also that although the element values of R, L, C correspond to an underdamped second order dynamics, the fractional *RLC* circuit do not show the oscillatory time response characteristics of an ordinary underdamped *RLC* circuit for small values of γ (see the plot for $\gamma = 0.25$ in Figure 3, see also the natural responses for $\gamma = 0.25, 0.50$ in Figure 1 of [17]). In fact, number of peaks (and hence oscillations) decreases with decreasing γ .

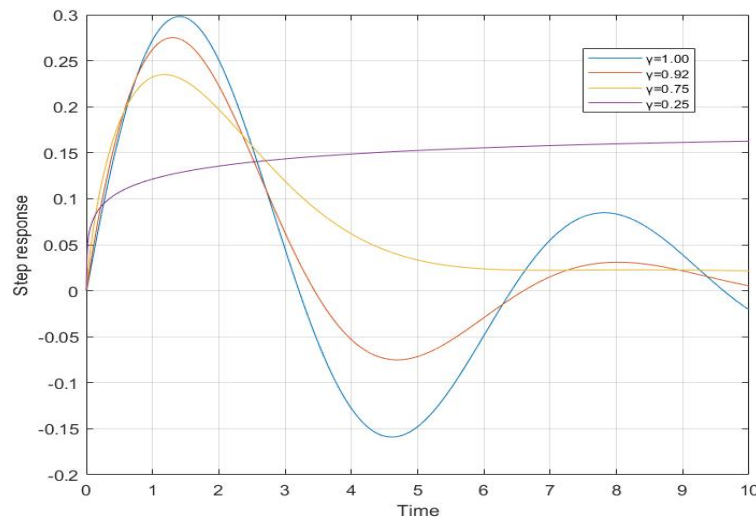
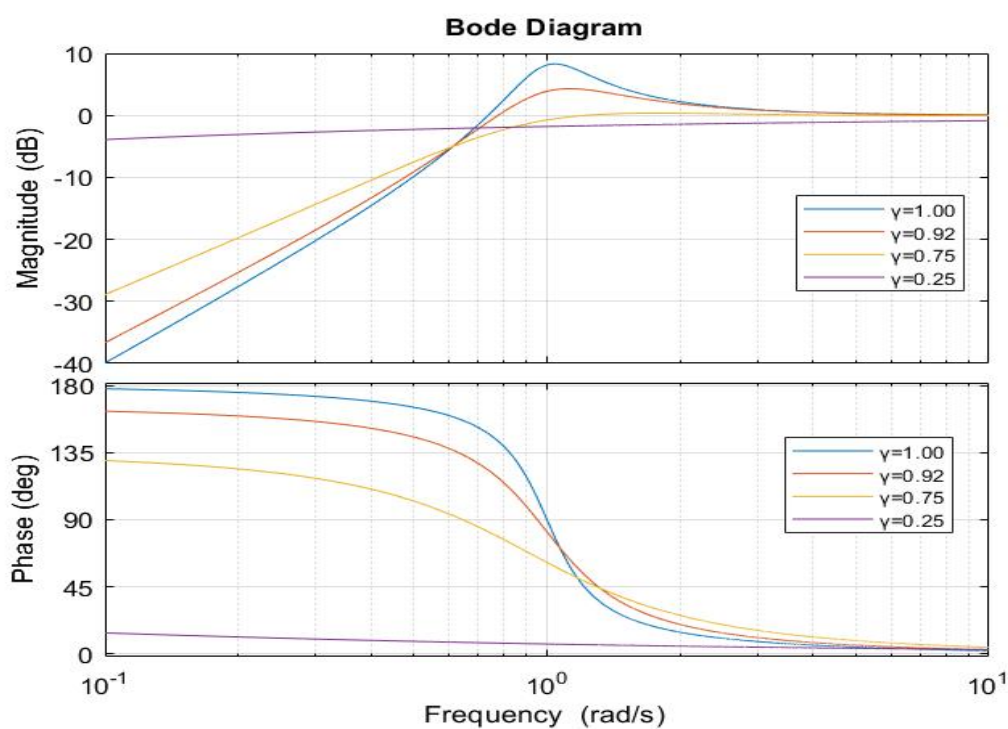


Figure 3. Step response of the band pass *RLC* filter of fractional order.

The high pass characteristics plotted in Figure 4 show that the 0 frequency and ∞ frequency gains are 0 and 1, respectively. But the peak value of the magnitude decreases with decreasing γ whilst the peak frequency is increasing. The relevant data are shown in Table 3 in detail. There is no peak value and stop bandwidth is very narrow for $\gamma = 0.25$. So, it is hard to say that the characteristic is of high pass type. The characteristic is like that of a band pass type (with higher peak gains > 0) around the center frequency for $\gamma = 1.00$ and 0.92 . For $\gamma = 0.25$, there is no peak in the gain. The phase responses in the same figure is seen to change from 180° to 0° , but flatter with decreasing γ as in the band pass case.

Table 2. The step response characteristics of band pass filter.

$\gamma \backslash Char.$	M_{p1} $t_{p1} (s)$	M_{p2} $t_{p2} (s)$	M_{p3} $t_{p3} (s)$	M_{p4} $t_{p4} (s)$	M_{p5} $t_{p5} (s)$
1.00	0.298	0.085	0.024	0.007	0.002
	1.401	7.809	14.216	20.624	27.032
0.92	0.275	0.031	0.006	0.002	—
	1.298	8.015	14.697	21.179	—
0.75	0.235	0.023	—	—	—
	1.170	8.366	—	—	—
0.25	0.171	—	—	—	—
	37.36	—	—	—	—

**Figure 4.** Frequency response of the high pass *RLC* filter of fractional order.**Table 3.** The frequency characteristics of high pass filter.

$\gamma \backslash Char$	$\omega_p (r/s)$	$M_p (dB)$	$\omega_1 (r/s)$	$\omega_2 (r/s)$	$BW (r/s)$	Q
1.00	1.041	8.300	1.362	1.215	0.875	2.138
0.92	1.124	4.299	2.414	1.396	0.829	0.709
0.75	1.778	0.359	—	1.861	0.768	—
0.25	—	—	—	0.635	0.103	—

The step response of the high pass filter is plotted in Figure 5 for values of $\gamma = 1.00, 0.92, 0.75, 0.25$; all the responses start from 1 at $t = 0$ and approach to zero with increasing time as expected due to high pass characteristics. It is observed that the magnitude and number of the oscillations reduces with decreasing γ . And the circuit step response gets away from the typical underdamped oscillatory characteristics of an underdamped second order RLC circuit as γ approaches to 0. In fact, the step response monotonically decays to zero for $\gamma = 0.25$. The relevant data are shown in Table 4.

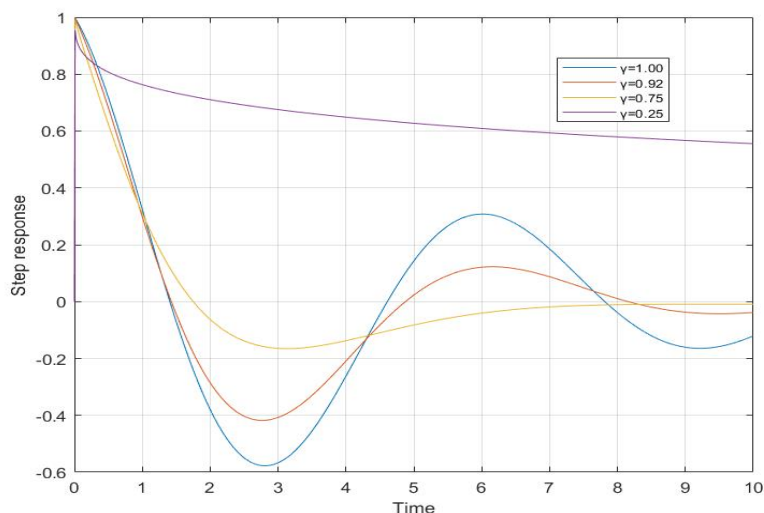


Figure 5. Step response of the high pass RLC filter of fractional order.

Table 4. The step response characteristics of high pass filter.

$\gamma \backslash Char.$	M_{p1} $t_{p1} (s)$	M_{p2} $t_{p2} (s)$	M_{p3} $t_{p3} (s)$	M_{p4} $t_{p4} (s)$	M_{p5} $t_{p5} (s)$
1.00	1.000 0	0.308 6.006	0.088 12.413	0.025 18.821	0.007 25.228
0.92	1.000 0	0.123 6.152	0.011 12.885	0.006 19.625	—
0.75	1.000 0	—	—	—	—
0.25	1.000 0	—	—	—	—

The low pass characteristics plotted in Figure 6 show that the 0 frequency and ∞ frequency gains are 1 and 0, respectively. The peak value of the magnitude decreases and disappears with decreasing γ . For $\gamma = 1.00$ and $\gamma = 0.92$, the characteristics is like that of band pass due to high peak values of the magnitudes. The relevant data are shown in Table 5. For $\gamma = 0.25$, the characteristic is of low-pass type with a very narrow band width. The phase responses in Figure 6 are seen to change from 0° to -180° , but flatter with decreasing γ as in the band pass case.

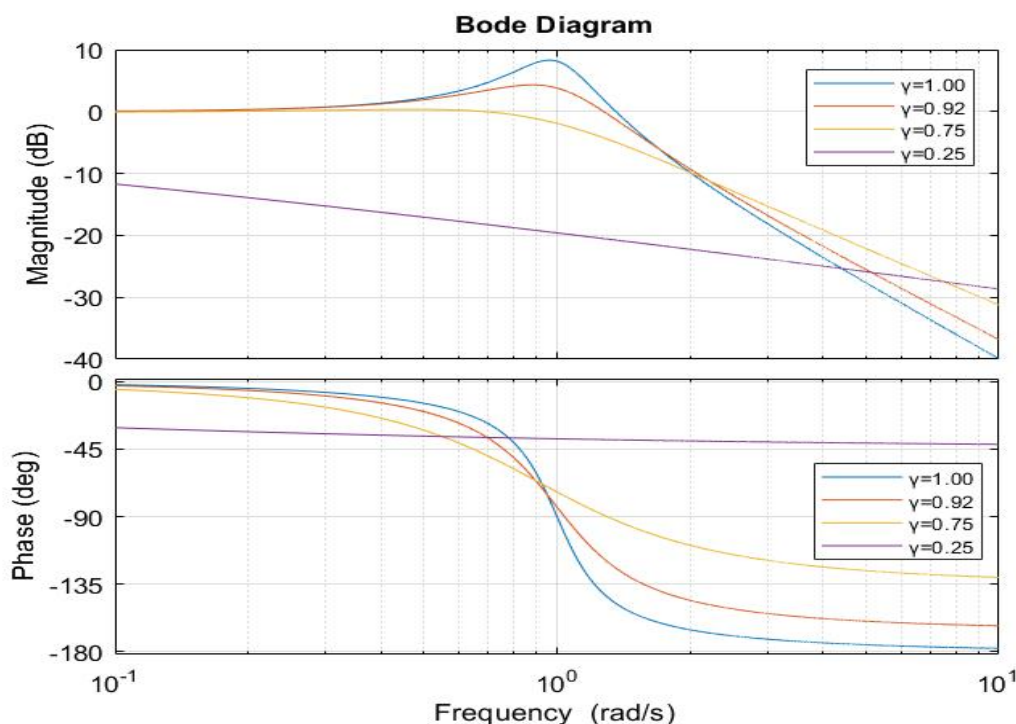


Figure 6. Frequency response of the low pass *RLC* filter of fractional order.

Table 5. Frequency response characteristics of low pass filter.

$\gamma \backslash Char$	ω_p (r/s)	M_p (dB)	ω_1 (r/s)	ω_2 (r/s)	BW (r/s)	Q
1.00	0.961	8.300	0.734	1.143	0.409	2.351
0.92	0.880	4.299	0.410	1.194	0.784	1.122
0.75	0.359	0.471	—	1.089	1.089	—
0.25	—	—	—	0.009	0.009	—

The step response of the low pass filter is shown in Figure 7 for values of $\gamma = 1.00, 0.92, 0.75, 0.25$. All the responses start from 0 at $t = 0$ and increase to unity as $t \rightarrow \infty$. It is observed that the magnitude of the oscillations and the number of oscillations reduce with decreasing γ . Again, the plots get away from the typical underdamped oscillatory characteristics of an underdamped second order *RLC* circuit as γ approaches to 0. In fact, the step response monotonically increases to 1 for $\gamma = 0.25$. The relevant data is shown in shown in Table 6. The rise time and the settling time could not be detected be for the case $\gamma = 0.25$ since the data is taken up to $t = 50s$.

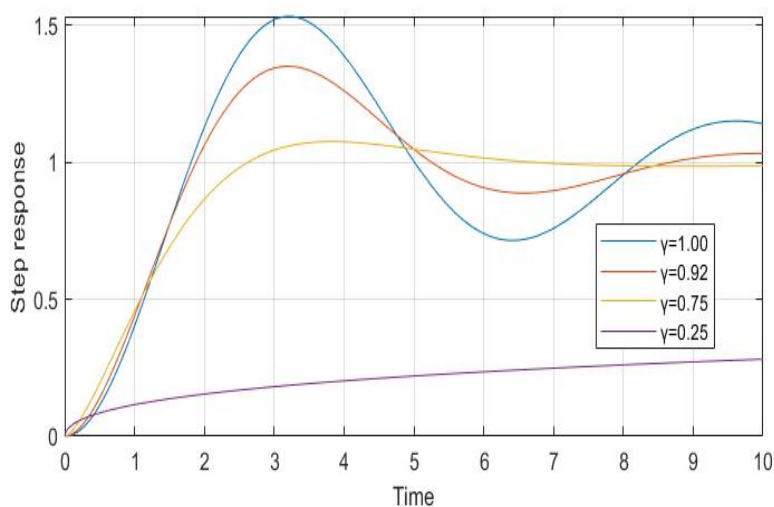


Figure 7. Step response of the low pass *RLC* filter of fractional order.

Table 6. The step response characteristics of low pass filter.

$\gamma \backslash Char.$	M_{p1} $t_{p1} (s)$	M_{p2} $t_{p2} (s)$	M_{p3} $t_{p3} (s)$	M_{p4} $t_{p4} (s)$	M_{p5} $t_{p5} (s)$	$T_r (s)$ $T_s (s)$
1.00	1.534	1.152	1.043	1.012	1.004	7.273
	3.204	9.612	16.020	22.427	28.835	13.738
0.92	1.351	1.033	1.002	0.9999	—	6.654
	3.190	9.928	11.971	23.500	—	7.875
0.75	1.077	—	—	—	—	1.815
	3.819	—	—	—	—	4.959
0.25	1	—	—	—	—	> 50
	∞	—	—	—	—	> 50

4. Sinusoidal responses

When a sine wave voltage is applied at the input, the resulting sinusoidal responses of the filters are discussed in this section. A sine wave input voltage of magnitude 1 and phase 0 is used as the voltage excitation in all the simulations.

The band pass filter sinusoidal response is checked by applying a sine wave of frequency $f = 0.125 \text{ Hz}$, $\omega = 0.25\pi \text{ r/s}$ which is smaller than the center frequencies in Table 1 (except $\gamma = 0.25$). The sinusoidal responses are shown in Figure 8, where the peaks and phases of the sine waves as approaching to steady-state are consistent with the frequency response curves in Figure 2. All the responses are leading the input by a phase decreasing with γ . All the magnitudes decrease with decreasing γ as well. Note that the consistency holds after a few periods of time when the steady state sinusoidal response has been reached. When the input frequency is increased to $\omega = 1$ and 0.4π

r/s , the sinusoidal responses are shown in Figures 11, 12 and 13 respectively. The consistency with the frequency response curves in Figure 2 still holds. For example, for $\omega = 0.4\pi$, all the responses are lagging due to the negative phase characteristics at this frequency. Note that all the responses have almost in phase with the input for $\omega = 1$ as seen in Figures 11 and 12, which is expected due to the small phases in Figure 2 at this frequency for all γ .

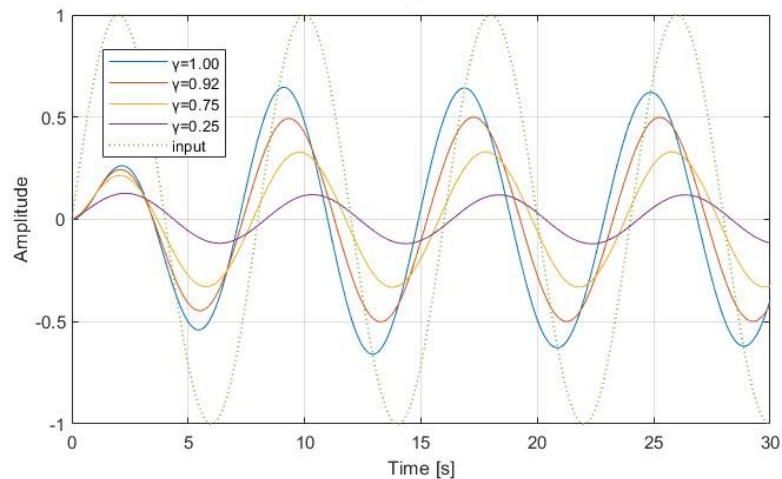


Figure 8. The sinusoidal response of BP filter with $w < 1r/s$.

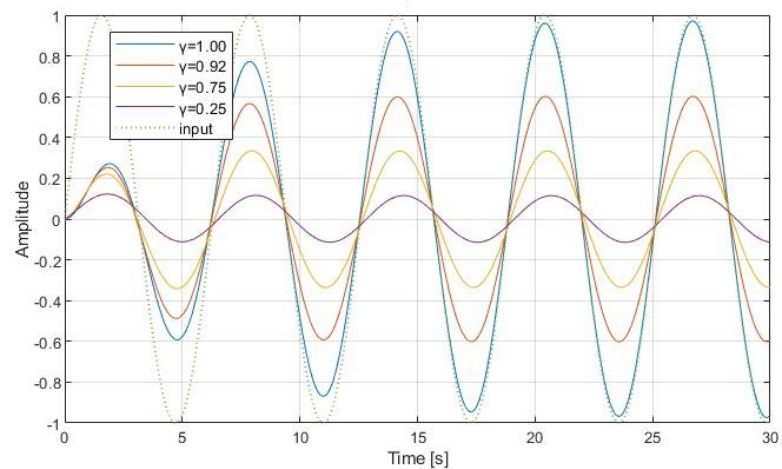


Figure 9. The sinusoidal response of BP filter with $w = 1r/s$.

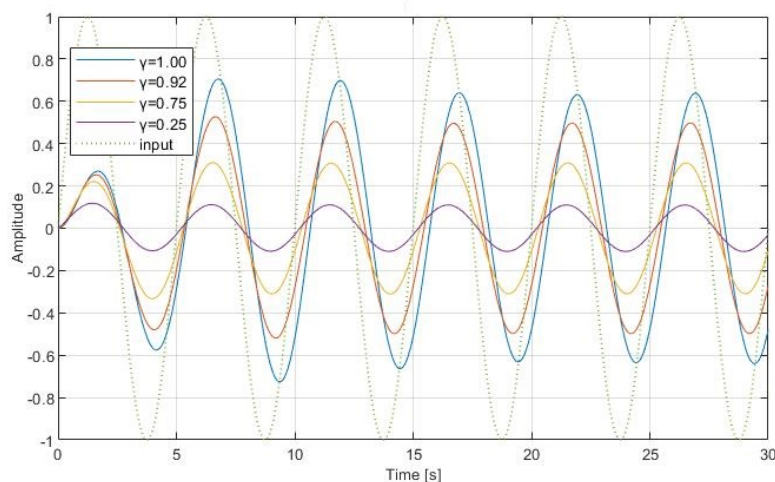


Figure 10. The sinusoidal response of BP filter with $w > 1r/s$.

The sinusoidal response of the high pass filter is obtained at two frequencies $\omega = 0.5$ and 1.5 . The results are shown in Figures 11 and 12. It is observed that the sinusoidal responses at the steady state conditions show the gain and phase characteristics implied by Figure 4.

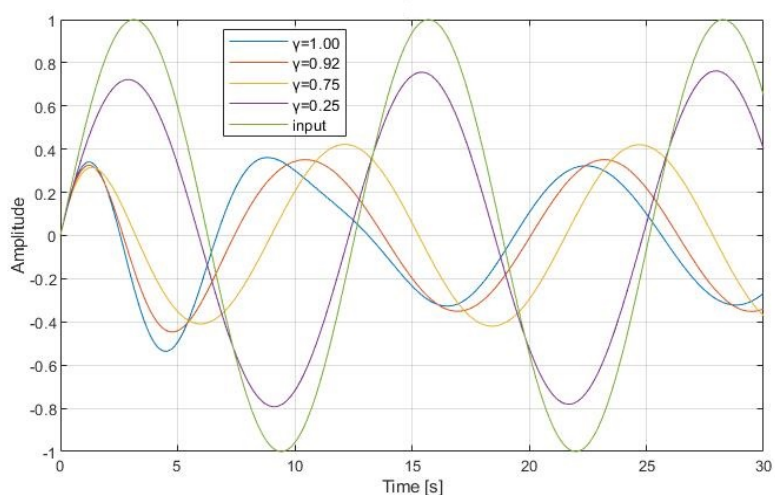


Figure 11. The sinusoidal response of HP filter with $w = 0.5r/s$.

Finally, the sinusoidal response of the low pass filter is obtained for a single frequency $\omega = 0.5 r/s$ in the pass band. It is seen in Figure 13 that all the responses are lagging the input in accordance to the phase response seen in Figure 6. Further, the gain is decreasing with decreasing γ as implied by the magnitude characteristic of the same figure.

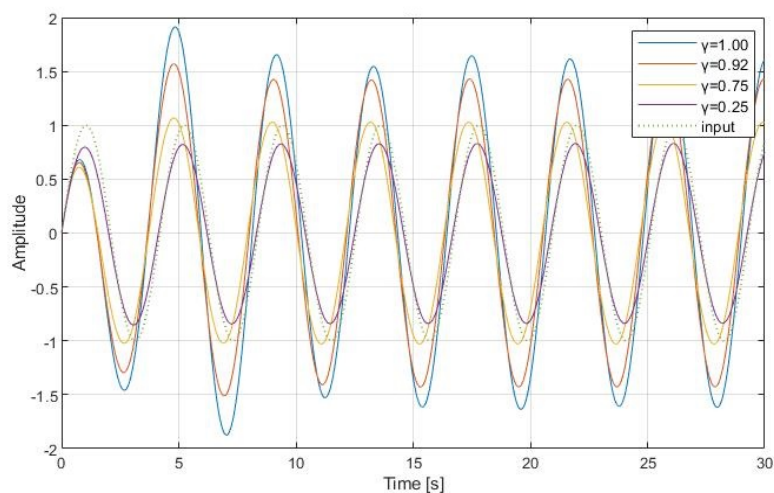


Figure 12. The sinusoidal response of HP filter with $w = 1.5r/s$.

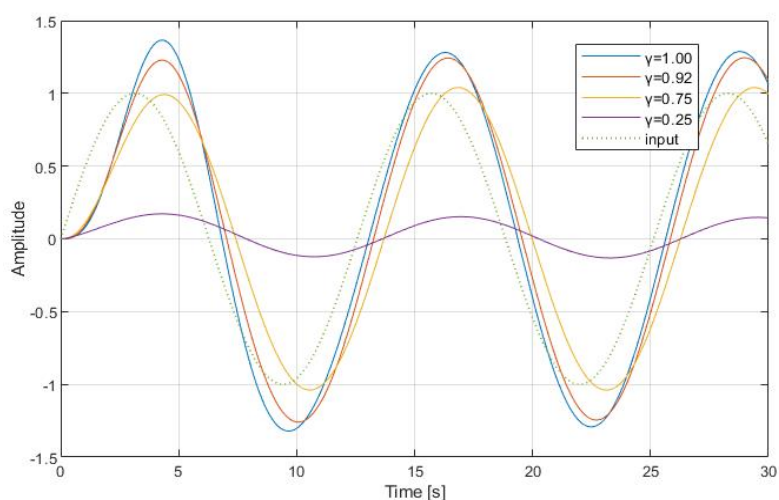


Figure 13. The sinusoidal response of LP filter with $w = 0.5r/s$.

The steady state variations of the sinusoidal responses in Figures 8–13, that is their variations for large time, are seen to be in coherence with the following formulas giving the steady state sinusoidal responses

$$y_{BP_{ss}}(t) = M_{BP}(w) \sin(\omega t + \theta_{BP}(w)),$$

$$y_{HP_{ss}}(t) = M_{HP}(w) \sin(\omega t + \theta_{HP}(w)),$$

$$y_{LP_{ss}}(t) = M_{LP}(w) \sin(\omega t + \theta_{LP}(w)),$$

where the gains M_i 's and the phases θ_i ' are as defined in Eqs. (3.4), (3.5), (3.6), respectively.

5. Conclusion

The frequency and time response characteristics of an RLC electrical circuit of non-integer order are examined in this paper. The circuit is shown to be used as band pass, high pass and low pass filters. The original RLC values are chosen for a second order ordinary RLC circuit of an underdamped type. The overdamped case is not considered since not all the considered three types of filter characteristics occur for this case. FOMCON Toolbox integrated with MATLAB R2017a is used to evaluate the fractional operations and Bode plots [19]. Fractional order RLC circuit is derived by applying the fractional derivative operator in consistency with the dimensionality [16] to the Kirchhoff's voltage law equation of the ordinary series RLC circuit excited by a voltage source [18]. The results stand as the answer to the future work indicated by F. Gomez et al. who obtained analytical solutions for time domain natural responses in terms of the Mittag-Leffler function for the treated fractional unexcited RLC circuit [17].

As the future work, the underdamped forced responses can be studied and the associated analytical results for the step responses can be searched. Similar research could also be conducted for higher order fractional RLC circuits.

Conflict of Interest

Authors declare that there is no conflict of interest.

References

1. R. Metzler, J. Klafter, *The restaurant at the end of the random walk: recent developments in the description of anomalous transport by fractional dynamics*, J. Phys. A: Math. Gen., **37** (2004), 161–208.
2. O. Tokatlı, V. Patoglu, *Using fractional order elements for haptic rendering*, Springer Proceedings in Advanced Robotics, **2** (2018), 373–388.
3. A. Karthikeyan, K. Rajagopal, *FPGA implementation of fractional-order discrete memristor chaotic system and its commensurate and incommensurate synchronizations*, Pramana-J. Phys., **90** (2018), 1–13.
4. A. Pratap, R. Raja, C. Sowmiya, et al. *Robust generalized Mittag-Leffler synchronization of fractional order neural networks with discontinuous activation and impulses*, Neural Networks, **103** (2018), 128–141.
5. Q. Sun, M. Xiao, B-B. Tao, et al. *Hopf bifurcation analysis in a fractional-order survival red blood cells model and PD^α control*, Adv Differ Equations, **10** (2018), 1–10.
6. J. Sabatier, F. Guillemard, L. Lavigne et al. *Fractional models of lithium-ion batteries with application to state of charge and ageing estimation*, Lecture Notes in Electrical Engineering, **430** (2018), 55–72.
7. K. Diethelm, D. Baleanu, E. Scalas, *Fractional calculus: Models and numerical methods*, Singapore: World Scientific, 2012.
8. M. A. E. Herzallah, *Notes on some fractional calculus operators and their properties*, J. Fract. Calc. Appl., **5** (2014), 1–10.

9. J. Ma, P. Zhou, B. Ahmad et al. *Chaos and multi-scroll attractors in RCL-shunted junction coupled Jerk circuit connected by memristor*, PloS One, **13** (2018), 1–21.
10. S. Bhalekar, V. D. Gejji, D. Baleanu, et al. *Transient chaos in fractional Bloch equations*, Comput. Math. Appl., **64** (2012), 3367–3376.
11. A. Razminia, D. Baleanu, *Fractional synchronization of chaotic systems with different orders*, P. Romanian Acad. A, **13** (2012), 314–321.
12. A. Jakubowska-Ciszek, J. Walczak, *Analysis of the transient state in a parallel circuit of the class $RL^\beta C^\alpha$* , Appl. Math. Comput., **319** (2018), 287–300.
13. A. Buscarino, R. Caponetto, G. D. Pasquale, et al. *Carbon black based capacitive fractional order element towards a new electronic device*, AEU-Int. J. Electron C., **84** (2018), 307–312.
14. P. Bertias, C. Psychalinos, *Differentiator based fractional-order high-pass filter designs*, 7th International Conference on Modern Circuits and Systems Technologies, 7–9 May 2018, (2018), 1–4.
15. A. Obeidat, M. Gharibeh, et al. *Fractional calculus and applied analysis*, **14**, Springer, 2011.
16. J. F. Gomez-Aguilar, J. J. Rosales-Garcia, J. J. Bernal-Alvarado, et al. *Fractional mechanical oscillators*, Rev. Mex. Fis., **58** (2012), 348–352.
17. F. Gomez, J. Rosales, M. Guia, *RLC electrical circuits of non-integer orders*, Cent. Eur. J. Phys., **11** (2013), 1361–1365.
18. C. Alexander, M. Sadiku, *Fundamentals of electric circuits*, 5 Eds., New York: Mc Graw Hill, 2013.
19. A. Tepljakov, E. Petlenkov, J. Belikov, et al. *Fractional-order controller design and digital implementation using FOMCON toolbox for MATLAB*, 2013 IEEE Conference on Computer Aided Control System Design, 28–30 Aug 2013, Hyderabad, India, (2013), 340–345.



AIMS Press

©2019 the Author(s), licensee AIMS Press. This is an open access article distributed under the terms of the Creative Commons Attribution License (<http://creativecommons.org/licenses/by/4.0>)