



Research article

Finite-time stability of elastoplastic lattice spring systems under simultaneous displacement-controlled and stress-controlled loadings[†]

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Abstract: According to Moreau (C.I.M.E. 1973), plastically deforming springs of lattice spring systems with a time-dependent displacement-controlled loading correspond to the attractor of a differential inclusion with a moving constraint of the form $C(t) = C + c(t)$, where C is a polytope and $c(t)$ is a time-dependent vector. Finite-time stability of differential inclusions of this type is established in Gudoshnikov et al. [SIAM J. Control Optim. 60 (2022)]. The work by Moreau also implies that accounting for a stress-controlled loading no longer allows to split $C(t)$ as $C + c(t)$. In the present paper we show that if we are interested in attractivity of a particular vertex of $C(t)$, then $C(t)$ can again be viewed as $C + c(t)$ for a specially constructed $c(t)$ (which depends on the vertex of interest), so that the technique of Gudoshnikov et al. can be used to obtain a criterion for finite-time stability of the vertex. The criterion obtained is illustrated with a benchmark example where we discover a drastic increase of diversity of possible combinations of plastically deforming springs when stress-controlled loading is introduced on top of displacement-controlled loading compared to the case where displacement-controlled loading is the only forcing.

Keywords: sweeping process; finite-time stability; time-dependent constraint; Lyapunov function; elastoplastic lattice spring model

1. Introduction

Finite-time stability in differential equations with nonsmooth right-hand-sides is addressed in Bernuau et al. [4], Bhat-Bernstein [5], Oza et al. [16], Sanchez et al. [19] over the Lyapunov function

V that solves

$$\frac{d}{dt}[V(x(t))] + 2\varepsilon \sqrt{V(x(t))} \leq 0, \quad \text{a.e. on } [0, \infty), \quad (1.1)$$

where $\varepsilon > 0$ and x is a solution. Adly et al. [1] extended the Lyapunov function approach to differential inclusions with a subdifferential of a convex function $\Phi : \mathbb{R} \rightarrow \mathbb{R}^n$ (given by a cone, see Rockafellar [18])

$$-\ddot{x}(t) - \nabla f(x(t)) \in \partial\Phi(\dot{x}(t)), \quad (1.2)$$

and derived an inequality of form (1.1) (and corresponding finite-time stability) from a cone-type condition

$$-\nabla f(x(t)) + B_\varepsilon(0) \subset \partial\Phi(0), \quad \text{a.e. on } [0, \infty), \quad (1.3)$$

where $B_\varepsilon(0)$ is the ball of \mathbb{R}^n of radius ε centered at 0.

More recently, a significant interest in the study of finite-time stability of differential inclusions has been due to new applications in elastoplasticity (see e.g., Gudoshnikov et al. [8]). We remind the reader that according to the pioneering work by Moreau [15] (see also Gudoshnikov and Makarenkov [9]), the stresses in a network of m elastoplastic springs with time-varying displacement-controlled loadings are governed by

$$-\dot{y} \in N_{C(t)}^A(y), \quad y \in V, \quad \begin{array}{l} V \text{ is a } d - \text{dimensional subspace of } \mathbb{R}^m \\ \text{with the scalar product } (x, y)_A = \langle x, Ay \rangle, \end{array} \quad (1.4)$$

where A is a positive diagonal $m \times m$ -matrix, and $N_{C(t)}^A(y)$ is a normal cone to the set

$$C(t) = \bigcap_{j=1}^m V_j(g(t), h(t)), \quad \begin{array}{l} V_j(g, h) = L(-1, j, g, h) \cap L(+1, j, g, h), \\ L(\alpha, j, g, h) = \{y \in V : \langle \alpha e_j, Ay + Ag \rangle \leq \alpha c_j^- + \alpha \alpha_j h\}, \end{array} \quad (1.5)$$

at a point y , with appropriate $d, c_j^-, c_j^+, g(t), h(t)$ that define mechanical parameters of the network of elastoplastic springs, displacement-controlled loadings and stress-controlled loadings (to be discussed in Section 4 in details), and where $e_j \in \mathbb{R}^m$ is the vector with 1 in the j -th component and zeros elsewhere. The solutions $y(t)$ of differential inclusion (1.4) never escape from $C(t)$ (i.e., $y(t)$ is swept by $C(t)$) for which reason (1.4) is called *sweeping process*. Spring j undergoes *plastic deformation* when the inequality $c_j^- < \langle e_j, Ay(t) - Ah(t) + Ag(t) \rangle < c_j^+$ is violated. Therefore, knowledge of the evolution of $y(t)$ allows to make conclusions about the regions of plastic deformation (that lead to *low-cycle fatigue* or *incremental failure*, see Yu [20, §4.6]).

Krejci [13] proved that if $C(t)$ is T -periodic then any solution $y(t)$ of (1.4) always converges to a T -periodic solution. Colombo et al. [7] proved the existence of the attractor in the case when $C(t)$ is a parallelepiped (of potentially changing dimensions). Extending their earlier two-dimensional version [8], Gudoshnikov et al. [11] offered a rule to compute the T -periodic attractor of (1.4) in the case where $h(t) \equiv 0$, i.e., when just displacement-controlled loading is present (meaning that the shape of the moving polytop doesn't change). Details of geometry of the attractor of (1.4) are addressed in Gudoshnikov et al. [12]. The goal of the present work is to investigate finite-time convergence of (1.4) in the case of displacement-controlled loading and stress-controlled loading present simultaneously.

As also mentioned in Gudoshnikov et al. [11], predicting the behavior of solutions of sweeping process (1.4) within a guaranteed time is of crucial importance for materials science. Current

methods of computing the asymptotic response of networks of elastoplastic springs (see e.g., Boudy et al. [6], Zouain-SantAnna [21]) run the numeric routine until the difference between the responses corresponding to two successive cycles of loading get smaller than a prescribed tolerance (without any estimate as for how soon such a desired accuracy will be reached).

The approach of the present paper is a suitable generalization of Gudoshnikov et al. [11] (pioneered by Adly et al. [1]). Specifically, let $y_*(t)$ be a vertex of $C(t)$. We prove that if $y_*(t)$ can be expressed as

$$y_*(t) = y_* + c(t), \quad (1.6)$$

$$\{y_*\} = \bigcap_{(\alpha,j) \in I_0} \bar{L}(\alpha, j), \quad (1.7)$$

$$\bar{L}(\alpha, j) = \{y \in V: \langle e_j, Ay \rangle = c_j^\alpha\},$$

where

$$\bar{L}(\alpha, j), (\alpha, j) \in I_0 \text{ are independent : } |I_0| = \dim V, \quad (1.8)$$

and if

$$-c'(t) + B_\varepsilon^{A(0)} \subset N_{C(t)-c(t)}^A(y_*), \quad \text{a.a. } t \in [0, \tau_d], \quad (1.9)$$

where $B_\varepsilon^A(0)$ is a ball in the norm induced by the scalar product (1.4), then, for any solution $y(t)$ of (1.4), the function

$$x(t) = y(t) - c(t) \quad (1.10)$$

satisfies the estimate (1.1) on $[0, \tau_d]$ for a suitable Lyapunov function V that measures the distance from $x(t)$ to y_* . In the earlier work [11], the difference $C(t) - c(t)$ is independent of t due to the lack of the stress-controlled loading. Accordingly, [11] assumes

$$-c'(t) + B_\varepsilon^A(0) \cap N_F^A(y) \subset N_C^A(F) \quad (1.11)$$

instead of (1.9), where F is however allowed to be a face, not just a vertex of C . The complication we encounter when adding a stress-controlled loading is that $C(t) - c(t)$ is no longer constant, but we discover that, for applications to lattice spring models, $y_*(t) - c(t)$ can still be assumed to be constant thanks to an appropriate construction of $c(t)$ (Lemma 4.2), which is the main result of the paper.

The paper is organized as follows: In Section 2, we adjust the proof of [11, Theorem 3.1] for the case where assumption (1.11) is replaced by (1.9). The corresponding theorem (Theorem 2.1) provides an estimate for the time it takes for any solution of (1.4) to reach $y_*(t)$. A corollary of Theorem 2.1 for the case where $c(t)$ is T -periodic is given in Section 3. Section 4 links the entries of Theorem 2.1 to parameters of a lattice spring model, which allows to specify the structure of $C(t)$ and $c(t)$ and to prove the main result of the paper (Lemma 4.2) about the existence of representation (1.6)-(1.7) for the case where sweeping process (1.4) comes as a model of an elastoplastic lattice spring model (as introduced in Moreau [15] and adapted to lattice spring models in Gudoshnikov et al. [9]). Section 5 combines Sections 2 and 4 in order to provide a step-by-step guide for computation of the entries of Theorem 2.1 in terms of mechanical parameters of elastoplastic lattice spring models. Section 6 follows the guide of Section 5 in order to investigate several instructional cases of a benchmark model of 5 springs on 4 nodes (Rachinskiy [17], Gudoshnikov et al. [11]) that allows us to clarify (Section 6.4) the role of the stress-controlled loading in the diversity of possible attractors of sweeping process (1.4) and, accordingly, in the diversity of different distributions of plastic deformations in elastoplastic lattice spring models. Conclusions section concludes the paper.

2. A sufficient condition for finite-time stability of sweeping process with a moving constraint of changing shape

We remind the reader that the normal cone $N_C^A(y)$ to the set C at a point $y \in C$ in a scalar product space V with the scalar product

$$(x, y)_A = \langle x, Ay \rangle, \quad \text{where } A \text{ is a diagonal positive } m \times m\text{-matrix}, \quad (2.1)$$

is defined as (see Bauschke and Combettes [3, §6.4])

$$N_C^A(y) = \begin{cases} \{x \in V : \langle x, A(\xi - y) \rangle \leq 0, \text{ for any } \xi \in C\}, & \text{if } y \in C, \\ \emptyset, & \text{if } y \notin C. \end{cases}$$

In what follows (see Bauschke and Combettes [3, §3.2])

$$\|x\|^A = \sqrt{\langle x, Ax \rangle}. \quad (2.2)$$

Definition 2.1. (see e.g., [14]) A set-valued function $C(t)$ (acting from \mathbb{R} to a vector space $V \subset \mathbb{R}^m$) is called Lipschitz continuous, if, for any $T > 0$, there exists $L > 0$ such that

$$d_H(C(t), C(s)) \leq L|t - s|, \quad t, s \in [0, T],$$

where $d_H(C_1, C_2)$ is Hausdorff distance between closed bounded sets $C_1, C_2 \in \mathbb{R}^m$.

We remind the reader that solution of an initial-value problem for sweeping processes (1.4) with Lipschitz continuous moving constraint $C(t)$ exists, unique and features continuous dependence on initial conditions (see e.g., Kunze and Monteiro Marques [14, Theorems 1–3]).

The statement of the following theorem and its proof follow the corresponding statement and proof of [11, Theorem 3.1], but we still rewrite the proof for completeness because [11, Theorem 3.1] uses condition (1.9) with $C(t) - c(t)$ replaced by C , i.e., $C(t) - c(t)$ is assumed totally constant in [11]. It turns out that replacing C by $C(t) - c(t)$ requires almost no changes in the proof.

Theorem 2.1. Let V be a d -dimensional linear subspace of \mathbb{R}^m with scalar product (2.1), $t \mapsto C(t)$ be a Lipschitz continuous multi-valued function with closed convex values, and $c : [0, \infty) \rightarrow V$ be Lipschitz continuous, and

$$y_* \in C(t) - c(t), \quad t \in [0, \infty). \quad (2.3)$$

Assume that there exists an $\varepsilon > 0$ such that condition (1.9) holds on an interval $[0, \tau_d]$ with

$$\tau_d \geq \frac{1}{\varepsilon} \cdot \sup_{v \in C(t) - c(t), t \geq 0} \|y_* - v\|^A. \quad (2.4)$$

Then, every solution y of (1.4) with the initial condition $y(0) \in C(0)$ satisfies $y(\tau_d) = y_* + c(\tau_d)$.

The proof of Theorem 2.1 follows the lines of [11]. The idea is based on observing that

$$V(v) = (\|v - y_*\|^A)^2 = \langle v - y_*, A(v - y_*) \rangle \quad (2.5)$$

is a Lyapunov function for the sweeping process

$$-x'(t) - c'(t) \in N_{C(t)-c(t)}^A(x(t)), \quad (2.6)$$

which is related to (1.4) through the change of the variables (1.10). Accordingly, $V(x(t_1)) = 0$ will imply $y(t_1) = y_* + c(t_1)$. For completeness, the proof of Theorem 2.1 is included in Appendix.

The supremum in (2.4) of Theorem 2.1 can be estimated using the following proposition (which extends [9, Proposition 3.14]).

Proposition 2.1. *For any $y \in C(t) - c(t)$, it holds that*

$$\sup_{v \in C(t) - c(t), t \geq 0} \|y - v\|^A \leq \|A^{-1}c^+ - A^{-1}c^-\|^A. \quad (2.7)$$

Proof. We have

$$\begin{aligned} \sup_{v \in C(t) - c(t), t \geq 0} \|y - v\|^A &\leq \sup_{v_1, v_2 \in C(t) - c(t), t \geq 0} \|v_1 - v_2\| \\ &= \sup_{u_1, u_2 \in C(t), t \geq 0} \|u_1 - u_2\| \\ &= \sup_{u_1, u_2 \in \Pi(t) \cap V, t \geq 0} \|u_1 - u_2\| \leq \sup_{u_1, u_2 \in \Pi(t), t \geq 0} \|u_1 - u_2\| \\ &= \sup_{u_1, u_2 \in A^{-1}C + h(t) - g(t), t \geq 0} \|u_1 - u_2\| = \max_{u_1, u_2 \in A^{-1}C} \|u_1 - u_2\|, \end{aligned}$$

which can be estimated from above by $\|A^{-1}c^+ - A^{-1}c^-\|^A$ according to [9, Proposition 3.14]. The proof of the proposition is complete. \square

3. The existence of a globally one-period stable periodic attractor

Corollary 3.1. *If, in the settings of Theorem 2.1, we additionally have that $c(t)$ is T -periodic with $T \geq \tau_d$, then y_* is a globally one-period stable T -periodic solution of (1.4).*

Corollary 3.1 follows by observing that

$$y_*(\tau_d + T) = y_* + c(\tau_d + T) = y_* + c(\tau_d) = y_*(\tau_d).$$

4. Application to a general elastoplastic system with displacement-controlled loading and stress-controlled loading present simultaneously

We remind the reader that according to Moreau [15] a network of m elastoplastic springs on n nodes with 1 displacement-controlled loading and subjected to a stress-controlled loading at all nodes is fully defined by an $m \times n$ kinematic matrix D of the topology of the network, $m \times m$ matrix of stiffnesses (Hooke's coefficients) $A = \text{diag}(a_1, \dots, a_m)$, an m -dimensional hyperrectangle $C = \prod_{j=1}^m [c_j^-, c_j^+]$ of the achievable stresses of springs (beyond which plastic deformation begins), a vector $R \in \mathbb{R}^m$ of the location of the displacement-controlled loading, a scalar function $l(t)$ that defines the magnitude of the displacement-controlled loading, and a function $\bar{h}(t) \in \mathbb{R}^m$ such that

$$f(t) = -D^T \bar{h}(t) \in \mathbb{R}^n \quad (4.1)$$

defines the forces applied at the n nodes of the network. When all springs are connected (form a connected graph), we have (see Bapat [2, Lemma 2.2])

$$\text{rank } D = n - 1. \quad (4.2)$$

We furthermore assume that

$$m \geq n \quad \text{and} \quad \text{rank}(D^T R) = 1. \quad (4.3)$$

To formulate the Moreau sweeping process corresponding to the elastoplastic system $(D, A, C, R, l(t), f(t))$, we first follow the 3 steps described in Gudoshnikov and Makarenkov [10, §5]:

- (1) Find an $n \times (n-2)$ -matrix M of $\text{rank}(DM) = n-2$ that solves $R^T DM = 0$ and use M to introduce $U_{basis} = DM$.
- (2) Find a matrix V_{basis} of $m-n+2$ linearly independent column vectors of \mathbb{R}^m that solves

$$(U_{basis})^T A V_{basis} = 0. \quad (4.4)$$

- (3) Find an $m \times (m-n+1)$ -matrix D^\perp that solves $(D^\perp)^T D = 0_{(m-n+1) \times n}$ and such that

$$\text{rank}(D^\perp) = m - n + 1. \quad (4.5)$$

With the new matrices introduced, the moving constraint $C(t)$ of sweeping process (1.4) corresponding to the elastoplastic system $(D, A, C, R, l(t), f(t))$ is given by

$$C(t) = \bigcap_{j=1}^m \left\{ y \in V : c_j^- + a_i h_i(t) \leq \langle e_j, Ay + A V_{basis} \bar{L} l(t) \rangle \leq c_j^+ + a_i h_i(t) \right\}, \quad (4.6)$$

where, for each $j \in \overline{1, m}$,

$$\bar{L} = W^{-1} \begin{pmatrix} 1 \\ 0_{m-n+1} \end{pmatrix}, \quad W = \begin{pmatrix} R^T \\ (D^\perp)^T \end{pmatrix} V_{basis} \quad (4.7)$$

with e_j being the basis vectors of \mathbb{R}^m , i.e., $e_j = (\underbrace{0, \dots, 0}_{j-1}, 1, 0, \dots, 0)^T$, and where $h(t)$ is found from $\bar{h}(t)$

according to a formula given by the following proposition ([11] focused on $h(t) = 0$ and didn't derive explicit formula for h). The existence of W^{-1} is proved in Gudoshnikov et al. [11].

Lemma 4.1. *If the stress-controlled loading $f(t)$ verifies the representation (4.1), then the corresponding function $h(t)$ in (4.6) is given by*

$$h(t) = U_{basis} \left((U_{basis})^T A U_{basis} \right)^{-1} (U_{basis})^T \bar{h}(t),$$

or, equivalently,

$$h(t) = -U_{basis} \left((U_{basis})^T A U_{basis} \right)^{-1} M^T f(t). \quad (4.8)$$

In particular, the $(n-2) \times (n-2)$ -matrix $(U_{basis})^T A U_{basis}$ is invertible.

Proof. First of all, we observe that the $(n-2) \times (n-2)$ matrix $(U_{basis})^T A U_{basis}$ is invertible. Indeed, if $(U_{basis})^T A U_{basis} x = 0$, then $(U_{basis} x)^T A (U_{basis} x) = 0$. But since the stiffness matrix A is positive definite, then $U_{basis} x = 0$. This means that $\text{Ker}(U_{basis}) = \text{Ker}((U_{basis})^T A U_{basis})$ and hence, $\text{Rank}(U_{basis}) = \text{Rank}((U_{basis})^T A U_{basis}) = n-2$.

According to [9, formula (16)],

$$h(t) = P_U A^{-1} \bar{h}(t), \quad (4.9)$$

where P_U is a linear (orthogonal in sense of scalar product (2.1)) projection map on U along V . To compute P_U we decompose x as

$$x = U_{basis}u + V_{basis}v, \quad (4.10)$$

and compute $u \in \mathbb{R}^{\dim U}$ from (4.10). Applying $(U_{basis})^T A$ to both sides of (4.10), using (4.4), and solving for u , we get

$$u = \left((U_{basis})^T K U_{basis} \right)^{-1} (U_{basis})^T A x,$$

which implies

$$P_U = U_{basis} \left((U_{basis})^T K U_{basis} \right)^{-1} (U_{basis})^T A x. \quad (4.11)$$

The statement of the proposition is obtained by substituting (4.11) to (4.9). The proof of the proposition is complete. \square

Using Lemma 4.1 and [9, Theorem 3.1], the solution $y(t)$ of sweeping process (1.4) is related to the vector $s(t) = (s_1(t), \dots, s_m(t))^T$ of the stresses of springs via

$$y(t) = A^{-1} s(t) - V_{basis} \bar{L}(t) - U_{basis} \left((U_{basis})^T A U_{basis} \right)^{-1} M^T f(t), \quad (4.12)$$

provided that $f(t)$ admits representation (4.1) or, equivalently (see [9, Remark 3.5]),

$$f_1(t) + \dots + f_n(t) = 0, \quad t \geq 0. \quad (4.13)$$

In contrast with Gudoshnikov et al. [11], the function $c(t)$, that we need to verify condition (1.9) for, depends on the choice of the vertex y_* , i.e., on the choice of I_0 in (1.7).

Definition 4.1. A set of indices $I_0 \subset \{-1, 1\} \times \overline{1, m}$ with $|I_0| = d$ will be called non-singular, if the matrix

$$\left(\{e_j, (\alpha, j) \in I_0\} \right)^T A V_{basis} \quad (4.14)$$

is invertible.

Remark 4.1. [11, formula (7.9)] If $I_0 \subset \{-1, 1\} \times \overline{1, m}$ with $|I_0| = d$ is non-singular then $\bigcap_{(\alpha, j) \in I_0} \bar{L}(\alpha, j)$ is a singleton and y_* in (1.7) is well defined.

Lemma 4.2. Let $I_0 \subset \{-1, 1\} \times \overline{1, m}$ be non-singular and let y_* be given by (1.7). Then, $y_*(t)$ given by (1.6) belongs to $C(t)$, if $c(t)$ is defined by

$$c(t) = -V_{basis} \bar{L}(t) + V_{basis} \Delta_{I_0} h(t), \quad (4.15)$$

$$\Delta_{I_0} = Z^{-1} L, \quad (4.16)$$

$$Z = \left(\{e_j : (\alpha, j) \in I_0\} \right)^T A V_{basis}, \quad (4.17)$$

$$L = \left(\{e_j : (\alpha, j) \in I_0\} \right)^T A, \quad (4.18)$$

and if the feasibility condition

$$c_j^- + a_j h_j(t) < \left\langle e_j, A y_* + A V_{basis} \Delta_{I_0} h(t) \right\rangle < c_j^+ + a_j h_j(t), \quad (\alpha, j) \notin I_0, \quad t \geq 0, \quad (4.19)$$

holds.

Remark 4.2. Equation (4.16) is equivalent to

$$a_j h_j(t) = \langle e_j, AV_{basis} \Delta_{I_0} h(t) \rangle, \quad (\alpha, j) \in I_0. \quad (4.20)$$

Proof of Lemma 4.2. Introduce

$$\begin{aligned} C_*(t) &= C_0(t) + C_1(t), \\ C_0(t) &= \bigcap_{(\alpha, j) \in I_0} \{y \in V : c_j^\alpha + a_j h_j(t) = \langle e_j, Ay + AV_{basis} \bar{L}l(t) \rangle\}, \\ C_1(t) &= \bigcap_{(\alpha, j) \notin I_0} \{y \in V : c_j^- + a_j h_j(t) \leq \langle e_j, Ay + AV_{basis} \bar{L}l(t) \rangle \leq c_j^+ + a_j h_j(t)\}. \end{aligned}$$

Since $C_*(t) \subset C(t)$, the problem of finding $y_*(t)$ with (1.6)-(1.7) reduces to finding $y_*(t)$ that satisfies

$$\{y_*(t)\} = C_0(t), \quad t \geq 0, \quad (4.21)$$

$$y_*(t) \subset C_1(t), \quad t \geq 0. \quad (4.22)$$

By (4.20), $C_0(t)$ can be rewritten in the form

$$C_0(t) = \bigcap_{(\alpha, j) \in I_0} \{y \in V : c_j^\alpha = \langle e_j, Ay + AV_{basis} \bar{L}l(t) - AV_{basis} \Delta_{I_0} h(t) \rangle\}. \quad (4.23)$$

To solve (4.20) we rewrite (4.20) as

$$L(t) = Z \Delta_{I_0} h(t).$$

Observe that if $y \in \{y \in V : F(y + k) = 0\}$ for some $k \in V$, then letting $x = y + k$, we have $x \in V$ and $F(x) = 0$, meaning that y can be represented as $y = x - k \in \{x \in V : F(x) = 0\} - k$.

Therefore, with formula (4.23), the expression for $C_0(t)$ can be written as

$$C_0(t) = C_0 + c(t),$$

where

$$C_0(t) = \bigcap_{(\alpha, j) \in I_0} \{y \in V : c_j^\alpha = \langle e_j, Ay \rangle\}$$

and where $c(t)$ is given by (4.15). Therefore, $C_0(t)$ is a singleton given by

$$C_0(t) = \{y_*(t)\},$$

where $y_*(t)$ is defined by (1.6)-(1.7), i.e., (4.21) is established.

The inclusion (4.22) follows from (4.19). The proof of the lemma is complete. \square

To understand what the conclusion

$$y(t) = y_* + c(t), \quad t \geq \tau_d, \quad (4.24)$$

of Theorem 2.1 says about the dynamics of the elastoplastic system $(D, A, C, R, l(t), f(t))$, recall that by [9, Theorem 3.1 and §3.2],

$$y(t) = A^{-1} s(t) + h(t) - V_{basis} \bar{L}l(t),$$

or, by combining with (4.24) and (4.15),

$$y_* + V_{basis} \Delta_{I_0} h(t) = A^{-1} s_*(t) + h(t),$$

i.e.,

$$s_*(t) = Ay_* + (V_{basis} \Delta_{I_0} - I)h(t), \quad t \geq \tau_d. \quad (4.25)$$

5. A step-by-step guide for analytic computations

The following lemma from Rockafellar-Wets [18, Theorem 6.46] (see also [11, Lemma 5.1]) gives a computational recipe to check condition (1.9). Recall that $\text{cone}\{\xi_1, \dots, \xi_K\}$ stays for the cone formed by vectors ξ_1, \dots, ξ_K .

Lemma 5.1. *Let V be a d -dimensional linear subspace of \mathbb{R}^m with scalar product (2.1). Consider*

$$\widetilde{C} = \bigcap_{k=1}^K \{y \in V : \langle \widetilde{n}_k, Ay \rangle \leq c_k\}, \quad (5.1)$$

where $\widetilde{n}_k \in V$, $c_k \in \mathbb{R}$, $K \in \mathbb{N}$. If $\widetilde{I}(y) = \{k \in \overline{1, K} : \langle \widetilde{n}_k, Ay \rangle = c_k\}$, then

$$N_{\widetilde{C}}^A(y) = \text{cone}\{\widetilde{n}_k : k \in \widetilde{I}(y)\}.$$

Corollary 5.1. *If $y_*(t) \in C(t)$ satisfies (1.6)-(1.7) for suitable $c(t)$ and y_* , then*

$$N_{C(t)-c(t)}^A(y_*) = \text{cone}\{\alpha n_j : (\alpha, j) \in I_0\}, \quad (5.2)$$

where $n_j \in V$, $j \in \overline{1, \dim V}$ are the vectors that solve the equality

$$\langle e_j, y \rangle = \langle n_j, y \rangle, \quad y \in V.$$

Proof. First note that

$$N_{C(t)-c(t)}^A(y_*) = N_{C(t)}^A(y_* + c(t)).$$

To apply Lemma 5.1, we rewrite (4.6) as follows:

$$\begin{aligned} C(t) = & \bigcap_{j=1}^m \left\{ y \in V : \langle -n_j, Ay + AV_{\text{basis}} \bar{L}(t) \rangle \leq -c_j^- - a_i h_i(t) \right\} \\ & \cup \left\{ y \in V : \langle n_j, Ay + AV_{\text{basis}} \bar{L}(t) \rangle \leq c_j^+ + a_i h_i(t) \right\}. \end{aligned}$$

Therefore, we want to compute $\widetilde{I}(y_* + c(t))$ given by

$$\widetilde{I}(y_* + c(t)) = \{(\alpha, j) \in \{-1, 1\} \times \overline{1, m} : \langle \alpha n_j, Ay_* + AV_{\text{basis}} \Delta_{I_0} h(t) \rangle = \alpha c_j^\alpha\}.$$

By successively using (4.19), (4.20), and (1.7) the above equality rewrites as

$$\begin{aligned} \widetilde{I}(y_* + c(t)) &= \{(\alpha, j) \in I_0 : \langle \alpha n_j, Ay_* + AV_{\text{basis}} \Delta_{I_0} h(t) \rangle = \alpha c_j^\alpha\} \\ &= \{(\alpha, j) \in I_0 : \langle \alpha n_j, Ay_* \rangle = \alpha c_j^\alpha\} = \{(\alpha, j) \in I_0 : \langle n_j, Ay_* \rangle = c_j^\alpha\} = I_0, \end{aligned}$$

which provides the required statement. The proof of the Corollary is complete. \square

Using (5.2), condition (1.9) can be rewritten as

$$-c'(t) + B_\varepsilon^A(0) \subset \text{cone}\{\alpha n_j : (\alpha, j) \in I_0\}. \quad (5.3)$$

Therefore, when $c'(t)$ is T -periodic, the required $\varepsilon > 0$ exists, if

$$\begin{aligned} I_0 \text{ admissibility condition: } & -c'(t) \in \text{cone} \left\{ \alpha n_j : (\alpha, j) \in I_0 \right\}, \quad t \in [0, T], \\ I_0 \text{ irreducibility condition: } & -c'(t) \notin \text{rb} \left(\text{cone} \left\{ \alpha n_j : (\alpha, j) \in I_0 \right\} \right), \quad t \in [0, T], \end{aligned} \quad (5.4)$$

where $\text{rb}(B)$ denotes the relative boundary of a set B of a linear subspace $V \subset \mathbb{R}^m$. The irreducibility condition can be further rewritten as

$$I_0 \text{ irreducibility condition: } \quad \tilde{I}_0 \text{ is not admissible for any } \tilde{I}_0 \subset I_0, \tilde{I}_0 \neq I_0. \quad (5.5)$$

In other words, (5.5) says that I_0 is irreducible, if any $\tilde{I}_0 \subset I_0$ with $\tilde{I}_0 \neq I_0$ makes

$$-c'(t) \notin \text{cone} \left\{ \alpha n_j : (\alpha, j) \in \tilde{I}_0 \right\}, \quad \text{for at least one } t \in [0, T].$$

According to [10, formula (27)], n_j can be taken as

$$n_j = V_{\text{basis}} W^{-1} \begin{pmatrix} R^T \\ (D^\perp)^T \end{pmatrix} e_j \quad (5.6)$$

and substitution of (5.2), (4.7), and (4.15) to (5.4) gives

$$\overbrace{V_{\text{basis}} W^{-1} \begin{pmatrix} l'(t) \\ 0_{m-n+1} \end{pmatrix}}^{-c'(t)} - V_{\text{basis}} \Delta_{I_0} h'(t) \in \text{cone} \left(V_{\text{basis}} W^{-1} \begin{pmatrix} R^T \\ (D^\perp)^T \end{pmatrix} \left\{ \alpha e_j : (\alpha, j) \in I_0 \right\} \right), \quad (5.7)$$

or

$$\begin{pmatrix} l'(t) \\ 0_{m-n+1} \end{pmatrix} - W \Delta_{I_0} h'(t) \in \text{cone} \left(\begin{pmatrix} R^T \\ (D^\perp)^T \end{pmatrix} \left\{ \alpha e_j : (\alpha, j) \in I_0 \right\} \right), \quad (5.8)$$

Based on Lemma 4.2 and formula (5.8) we can now split verification of conditions of Theorem 2.1 into the following steps.

Step 1. Fix appropriate indexes I_0 (springs that will reach plastic deformation). Spot a non-singular I_0 that is admissible and irreducible (i.e., solves (5.8) and such that \tilde{I}_0 doesn't satisfy (5.8) for any $\tilde{I}_0 \subset I_0, \tilde{I}_0 \neq I_0$).

Combination of Theorem 2.1 and Lemma 4.2 lead to the following qualitative description of the asymptotic behavior of elastoplastic system $(D, A, C, R, l(t), f(t))$ and of the associated sweeping process (1.4). Here the quantities y_* , $c(\cdot)$, I_0 , Δ_{I_0} , $h(\cdot)$ are those referred to in the statement of Lemma 4.2.

Proposition 5.1. (Conclusion of Step 1). *If*

- 1) both $l'(t)$ and $f'(t)$ are constant,
- 2) $f(t)$ satisfies the static balance condition (4.13),
- 3) there exists a non-singular and irreducible I_0 satisfying (5.8),

then there exist c_j^α , $(\alpha, j) \notin I_0$, such that the solution y of sweeping process (1.4) with initial condition $y(0) \in C(0)$ satisfies

$$y(\tau_d) = y_* + c(\tau_d), \quad (5.9)$$

for an appropriate $\tau_d > 0$, and the stress vector $s(t)$ of the elastoplastic system $(D, A, C, R, l(t), f(t))$ satisfies

$$s(\tau_d) = Ay_* + (V_{basis}\Delta_{I_0} - I)h(\tau_d), \quad (5.10)$$

$$s_j(\tau_d) = c_j^\alpha, \quad (\alpha, j) \in I_0. \quad (5.11)$$

Indeed, conditions of Proposition 5.1 imply that condition (1.9) of Theorem 2.1 is satisfied on $[0, \tau_d]$, where

$$\tau_d = \frac{1}{\varepsilon} \cdot \sup_{v \in C(t)-c(t), t \geq 0} \|y_* - v\|^A \quad (5.12)$$

with suitable $\varepsilon > 0$, which gives τ_d used in the statement. Relation (5.11) follows from (5.10) by Remark 4.2.

Remark 5.1. Relation (5.11) means that springs with indices I_0 get to plastic mode (i.e., capable to deform plastically) by the time τ_d .

One has to proceed to Steps 2 and 3 to come up with an explicit version of Proposition 5.1 where conditions for c_i^α and formulas for ε and τ_d are given in closed form.

Step 2. Compute y_* and impose the feasibility condition. By [11, formula (7.9)] the formula for y_* of (1.7) reads as

$$y_* = V_{basis} \left(\left(\left\{ e_j, (\alpha, j) \in I_0 \right\} \right)^T A V_{basis} \right)^{-1} \left(\left\{ c_j^\alpha, (\alpha, j) \in I_0 \right\} \right)^T. \quad (5.13)$$

Substitution of y_* , Δ_{I_0} , $h(t)$ as defined by (5.13), (4.16), (4.8) to (4.19) yields feasibility condition in terms of mechanical parameters of the lattice-spring model.

Note, function $f(t)$ must additionally satisfy (4.13) for formula (4.8) to be valid.

Step 3. Compute ε_0 . This step is devoted to finding ε for which assumption (5.3) holds. Assumption (5.3) requires computation of the distance from $-c'(t)$ to the boundary of cone $\text{cone}\{an_j : (\alpha, j) \in I_0\}$. The required boundary is $\partial \text{cone}\{an_j : (\alpha, j) \in I_0\}$.

Using formula (5.2), we compute

$$\begin{aligned} \varepsilon_0(t) &= \text{dist}^A \left(-c'(t), \partial \text{cone}\{an_j : (\alpha, j) \in I_0\} \right) \\ &= \min_{(\alpha_*, j_*) \in I_0} \text{dist}^A \left(-c'(t), \text{cone}\{an_j : (\alpha, j) \in I_0 \setminus \{(\alpha_*, j_*)\}\} \right). \end{aligned} \quad (5.14)$$

According to [11, formula (7.18)] and [11, Lemma 7.7], the quantity

$$\varepsilon_0(t) = \text{dist}^A \left(-c'(t), \partial \text{cone}\{an_j : (\alpha, j) \in I_0\} \right)$$

computes as (see [11, formula (7.20)])

$$\bar{\varepsilon}_0(t) = \min_{(\alpha_*, j_*) \in I_0} \left\| -c'(t) - \text{proj}^A \left(-c'(t), \text{span}\{an_j : (\alpha, j) \in I_0 \setminus \{(\alpha_*, j_*)\}\} \right) \right\|^A, \quad (5.15)$$

where

$$\begin{aligned}
 & \text{proj}^A \left(-c'(t), \text{span} \{ \alpha n_j : (\alpha, j) \in I_0 \setminus \{(\alpha_*, j_*)\} \} \right) \\
 &= -(\{n_j, (\alpha, j) \in I_0 \setminus \{(\alpha_*, j_*)\}\}) \circ \\
 &\circ \left[(\{n_j, (\alpha, j) \in I_0 \setminus \{(\alpha_*, j_*)\}\})^T A (\{n_j, (\alpha, j) \in I_0 \setminus \{(\alpha_*, j_*)\}\}) \right]^{-1} \\
 &\circ (\{n_j, (\alpha, j) \in I_0 \setminus \{(\alpha_*, j_*)\}\})^T A c'(t).
 \end{aligned} \tag{5.16}$$

Choose $\varepsilon_0 > 0$ such that $\varepsilon_0 \leq \bar{\varepsilon}_0(t)$ for all $t \in [0, \tau_d]$. Theorem 2.1, Proposition 2.1, and Lemma 4.2 then lead to the following conclusion.

Proposition 5.2. (Conclusion of Steps 1–3). *If assumptions 1)–3) of Proposition 5.1 hold, then condition (1.9) holds on $[0, \tau_d]$ for any $\tau_d \geq \tau$, where*

$$\tau = \frac{1}{\varepsilon_0} \cdot \|A^{-1}c^+ - A^{-1}c^-\|^A.$$

If, additionally,

4) condition (4.19) holds on some $[0, \tau_d]$ with $\tau_d \geq \tau$,

then conclusions (5.9)–(5.11) of Proposition 5.1 hold with τ_d replaced by t , $t \in [\tau, \tau_d]$.

6. A benchmark example

The focus of the present section is on the elastoplastic model shown in Figure 1 (earlier introduced in Rachinskiy [17]), which allows to fully illustrate the practical implementation of Theorem 2.1.

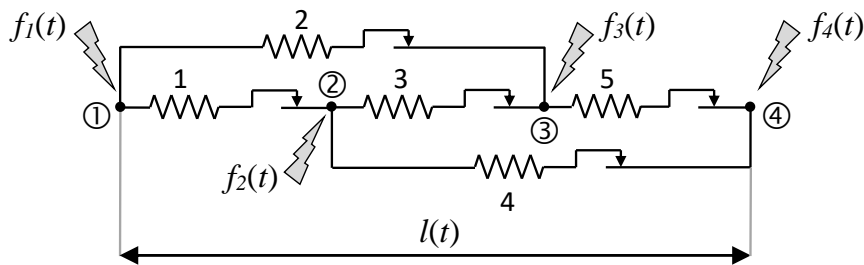


Figure 1. A system of 5 elastoplastic springs on 4 nodes with displacement-controlled loading $l(t)$ and stress-controlled loading $f(t) = (f_1(t), f_2(t), f_3(t), f_4(t))^T$. Circled numbers are indices of nodes and regular numbers are indices or springs.

According to Gudoshnikov et al. [11], the elastoplastic system of Figure 1 leads to the following expressions for D and R

$$D = \begin{pmatrix} -1 & 1 & 0 & 0 \\ -1 & 0 & 1 & 0 \\ 0 & -1 & 1 & 0 \\ 0 & -1 & 0 & 1 \\ 0 & 0 & -1 & 1 \end{pmatrix}, \quad R = \begin{pmatrix} 1 \\ 0 \\ 1 \\ 0 \\ 1 \end{pmatrix}.$$

We now follow Gudoshnikov et al. [11] to formulate the quantities used in Section 5.

First of all, based on [10, formula (17)], we compute the dimension of sweeping process (1.4) as

$$\dim V = m - n + q + 1 = 5 - 4 + 1 + 1 = 3.$$

According to [10, §5, Step 1], we then look for an 4×2 matrix M such that $R^T DM = 0$ and such that the matrix DM is full rank. Such a matrix M can be taken as

$$M = \begin{pmatrix} 0 & 0 \\ 1 & 1 \\ 1 & -1 \\ 0 & 0 \end{pmatrix} \quad \text{with} \quad DM = \begin{pmatrix} 1 & 1 \\ 1 & -1 \\ 0 & -2 \\ -1 & -1 \\ -1 & 1 \end{pmatrix} = U_{basis}.$$

The next step is determining V_{basis} which consists of $d = 3$ linearly independent columns of $\mathbb{R}^m = \mathbb{R}^5$ and solves $(DM)^T AV_{basis} = 0$. Such a V_{basis} can be taken as

$$V_{basis} = \begin{pmatrix} 0 & 1/a_1 & 1/a_1 \\ 0 & 1/a_2 & -1/a_2 \\ 1/a_3 & 0 & 1/a_3 \\ -1/a_4 & 1/a_4 & 0 \\ 1/a_5 & 1/a_5 & 0 \end{pmatrix} \quad \text{with} \quad AV_{basis} = \begin{pmatrix} 0 & 1 & 1 \\ 0 & 1 & -1 \\ 1 & 0 & 1 \\ -1 & 1 & 0 \\ 1 & 1 & 0 \end{pmatrix}.$$

Finally, a 5×2 full rank matrix D^\perp satisfying $(D^\perp)^T D = 0$ can be taken as

$$D^\perp = \begin{pmatrix} 0 & 1 \\ 0 & -1 \\ 1 & 1 \\ -1 & 0 \\ 1 & 0 \end{pmatrix} \quad \text{leading to} \quad \begin{pmatrix} R^T \\ (D^\perp)^T \end{pmatrix} = \begin{pmatrix} 1 & 0 & 1 & 0 & 1 \\ 0 & 0 & 1 & -1 & 1 \\ 1 & -1 & 1 & 0 & 0 \end{pmatrix}. \quad (6.1)$$

Formula (5.6) then yields

$$\begin{aligned} n_1 &= (a_2(a_3 + a_4) + a_4a_5 + a_3(a_4 + a_5), -a_1a_3, a_1(a_2 + a_5), a_1(a_2 + a_3 + a_5), a_1a_3)^T / k, \\ n_2 &= (-a_2a_3, a_4a_5 + a_1(a_3 + a_5) + a_3(a_4 + a_5), -a_2(a_1 + a_4), a_2a_3, a_2(a_1 + a_3 + a_4))^T / k, \\ n_3 &= (a_3(a_2 + a_5), -a_3(a_1 + a_4), (a_1 + a_4)(a_2 + a_5), -a_3(a_2 + a_5), a_3(a_1 + a_4))^T / k, \\ n_4 &= (a_4(a_2 + a_3 + a_4), a_3a_4, -a_4(a_2 + a_5), a_3(a_2 + a_5) + a_1(a_2 + a_3 + a_5), -a_3a_4)^T / k, \\ n_5 &= (a_3a_5, (a_1 + a_3 + a_4)a_5, (a_1 + a_4)a_5, -a_3a_5, a_1(a_2 + a_3) + a_3a_4 + a_2(a_3 + a_4))^T / k \end{aligned} \quad (6.2)$$

with

$$k = a_3a_4 + a_2(a_3 + a_4) + a_3a_5 + a_4a_5 + a_1(a_2 + a_3 + a_5). \quad (6.3)$$

In what follows, we consider

$$f(t) = \begin{pmatrix} b_1t + d_1 \\ b_2t + d_2 \\ b_3t + d_3 \\ b_4t + d_4 \end{pmatrix}, \quad \sum_{i=1}^4 b_i = \sum_{i=1}^4 d_i = 0, \quad l(t) = l_0 + l_1t, \quad (6.4)$$

where $l_0, l_1 > 0, b_1, b_2, b_3, b_4, d_1, d_2, d_3, d_4 \in \mathbb{R}$, which leads to

$$h(t) = \frac{1}{k} \begin{pmatrix} -(a_2(tb_2 + d_2) + a_5(tb_2 + d_2) + a_3(tb_2 + tb_3 + d_2 + d_3)) \\ -((a_1 + a_4)(tb_3 + d_3) + a_3(tb_2 + tb_3 + d_2 + d_3)) \\ a_2(tb_2 + d_2) + a_5(tb_2 + d_2) - (a_1 + a_4)(tb_3 + d_3) \\ a_2(tb_2 + d_2) + a_5(tb_2 + d_2) + a_3(tb_2 + tb_3 + d_2 + d_3) \\ (a_1 + a_4)(tb_3 + d_3) + a_3(tb_2 + tb_3 + d_2 + d_3) \end{pmatrix}$$

with k given by (6.3).

We now illustrate computations of Steps 1–3 of Section 5 and statements of Propositions 5.1 and 5.2 through three instructive cases of I_0 .

6.1. Case $I_0 = \{(+, 2), (+, 3), (+, 4)\}$

This case was addressed in [11] without stress-controlled loading.

Step 1. Computation of matrix Δ_{I_0} gives

$$\Delta_{I_0} = \begin{pmatrix} 0 & 1 & -1 \\ 1 & 0 & 1 \\ -1 & 1 & 0 \end{pmatrix}^{-1} \circ \begin{pmatrix} 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 \end{pmatrix} A = \frac{1}{2} \begin{pmatrix} 0 & a_2 & a_3 & -a_4 & 0 \\ 0 & a_2 & a_3 & a_4 & 0 \\ 0 & -a_2 & a_3 & a_4 & 0 \end{pmatrix}. \quad (6.5)$$

Substituting to (5.8) gives

$$\begin{pmatrix} l_1 - \frac{b_2}{a_1} + \frac{b_3}{a_5} \\ \frac{b_3}{a_5} \\ -\frac{b_2}{a_1} \end{pmatrix} \in \text{cone} \left\{ \begin{pmatrix} 0 \\ 0 \\ -1 \end{pmatrix}, \begin{pmatrix} 1 \\ 1 \\ 1 \end{pmatrix}, \begin{pmatrix} 0 \\ -1 \\ 0 \end{pmatrix} \right\}. \quad (6.6)$$

The vector in the left-hand-side won't belong to the boundary of the right-hand-side (i.e., I_0 will be admissible and irreducible), if

$$l_1 + \frac{b_3}{a_5} > 0, \quad l_1 - \frac{b_2}{a_1} + \frac{b_3}{a_5} > 0, \quad l_1 - \frac{b_2}{a_1} > 0. \quad (6.7)$$

Remark 6.1. Conditions (6.7) can be viewed as a condition for l_1 to be sufficiently large or as a condition for $|b_2| + |b_3|$ to be sufficiently small.

Proposition 6.1. (Conclusion of Step 1). Assume that stress-controlled loading $f(t)$ and displacement-controlled loading $l(t)$ of elastoplastic system of Figure 1 are given by (6.4) and satisfy (6.7). Then c_j^α , $(\alpha, j) \notin I_0$, can be amended in such a way that the stress vector $s(t)$ of the elastoplastic system $(D, A, C, R, l(t), f(t))$ satisfies

$$s(\tau_d) = Ay_* + (V_{\text{basis}}\Delta_{I_0} - I)h(\tau_d), \quad (6.8)$$

$$s_j(\tau_d) = c_j^\alpha, \quad (\alpha, j) \in I_0. \quad (6.9)$$

Step 2. Formula (5.13) yields

$$Ay_{I_0} = \begin{pmatrix} c_3^+ + c_4^+ \\ c_2^+ \\ c_3^+ \\ c_4^+ \\ c_2^+ + c_3^+ \end{pmatrix} \quad (6.10)$$

and substitution to (4.19) returns

$$\begin{aligned} c_1^- &< tb_2 + d_2 + c_3^+ + c_4^+ < c_1^+, \\ c_5^- &< -tb_3 - d_3 + c_2^+ + c_3^+ < c_5^+. \end{aligned} \quad (6.11)$$

Step 3. Since $|I_0| = 3$, for any $(\alpha_*, j_*) \in I_0$, the set $I_0 \setminus \{(\alpha_*, j_*)\}$ consists of two elements $\{(\alpha_1, j_1), (\alpha_2, j_2)\}$, and formulas (5.15) and (5.16) can be rewritten as

$$\varepsilon_0(t) = \min_{(\alpha_1, j_1), (\alpha_2, j_2) \in I_0} S_{j_1 j_2}, \quad (6.12)$$

where $S_{j_1 j_2} = \|-c'(t) - \text{proj}(-c'(t), \text{span}\{n_{j_1}, n_{j_2}\})\|^A$,

and $\text{proj}^A(-c'(t), \text{span}\{n_{j_1}, n_{j_2}\}) = -(n_{j_1} \ n_{j_2}) \begin{pmatrix} n_{j_1}^T A n_{j_1} & n_{j_1}^T A n_{j_2} \\ n_{j_2}^T A n_{j_1} & n_{j_2}^T A n_{j_2} \end{pmatrix}^{-1} \begin{pmatrix} n_{j_1}^T \\ n_{j_2}^T \end{pmatrix} A c'(t).$

Substituting n_j and $c'(t)$ given by (6.2) and (5.7) to (6.12), we get

$$\begin{aligned} \varepsilon_0 &= \min\{S_{23}, S_{34}, S_{24}\}, \\ S_{23} &= \sqrt{\frac{a_4(-l_1 a_1 + b_2)^2}{a_1(a_1 + a_4)}}, \quad S_{34} = \sqrt{\frac{a_2(l_1 a_5 + b_3)^2}{a_5(a_2 + a_5)}}, \\ S_{24} &= \sqrt{\frac{a_3(a_1 b_3 + a_5(-b_2 + a_1 l_1))^2}{a_1 a_5(a_3 a_5 + a_1(a_3 + a_5))}}. \end{aligned} \quad (6.13)$$

Proposition 6.2. (Conclusion of Steps 1–3). Assume that assumptions of Proposition 6.1 hold. Let ε_0 be given by (6.13) and assume that

$$\text{time } \tau = \frac{1}{\varepsilon_0} \cdot \|A^{-1}c^+ - A^{-1}c^-\|^A \text{ satisfies (6.11)}. \quad (6.14)$$

Then the stress vector $s(t)$ of the elastoplastic system $(D, A, C, R, l(t), f(t))$ satisfies

$$s(t) = Ay_{I_0} + (V_{\text{basis}} \Delta_{I_0} - I)h(t), \quad (6.15)$$

$$s_j(t) = c_j^\alpha, \quad (\alpha, j) \in I_0, \quad (6.16)$$

for all $t \geq \tau$ that satisfy (6.11).

Remark 6.2. When $b_2 = d_2 = b_3 = d_3 = 0$ (e.g., in the absence of the stress-controlled loading), condition (6.11) is the standard feasibility condition of vertex y_* ([11, formula (8.5)]). Therefore, condition (6.11) can be viewed as a condition for the stress-controlled loading to not be too big.

6.2. Case $I_0 = \{(+, 2), (-, 3), (+, 4)\}$

Formula (6.5) stays because computation of Δ_{I_0} is independent of the signs of α in $(\alpha, j) \in I_0$. The relations (6.6) and (6.10) take the form

$$\begin{pmatrix} l_1 - \frac{b_2}{a_1} + \frac{b_3}{a_5} \\ \frac{b_3}{a_5} \\ -\frac{b_2}{a_1} \end{pmatrix} \in \text{cone} \left(\left\{ \begin{pmatrix} 0 \\ 0 \\ -1 \end{pmatrix}, \begin{pmatrix} -1 \\ -1 \\ -1 \end{pmatrix}, \begin{pmatrix} 0 \\ -1 \\ 0 \end{pmatrix} \right\} \right), \quad Ay_{I_0} = \begin{pmatrix} c_3^- + c_4^+ \\ c_2^+ \\ c_3^- \\ c_4^+ \\ c_3^- + c_2^+ \end{pmatrix}.$$

Admissibility and irreducibility condition (6.7), feasibility condition (6.11), and convergence time (6.11) now compute as follows.

Admissibility and irreducibility condition:

$$l_1 + \frac{b_3}{a_5} > 0, \quad \frac{b_2}{a_1} - \frac{b_3}{a_5} - l_1 > 0, \quad l_1 - \frac{b_2}{a_1} > 0. \quad (6.17)$$

Feasibility condition:

$$\begin{aligned} c_1^- &< tb_2 + d_2 + c_3^- + c_4^+ < c_1^+, \\ c_5^- &< -tb_3 - d_3 + c_3^- + c_2^+ < c_5^+. \end{aligned} \quad (6.18)$$

Convergence time:

$$t \geq \max \left\{ \sqrt{\frac{a_1(a_1 + a_4)}{a_4(-l_1 a_1 + b_2)^2}}, \sqrt{\frac{a_5(a_2 + a_5)}{a_2(l_1 a_5 + b_3)^2}}, \sqrt{\frac{a_1 a_5 (a_3 a_5 + a_1 (a_3 + a_5))}{a_3 (a_1 b_3 + a_5 (-b_2 + a_1 l_1))^2}} \right\} \|A^{-1}c^+ - A^{-1}c^-\|^A. \quad (6.19)$$

Proposition 6.2 can now be concisely formulated as follows.

Proposition 6.3. Assume that condition (6.17) holds. Then, for any $t \geq 0$ that satisfies simultaneously (6.18) and (6.19) the stress vector $s(t)$ of elastoplastic system $(D, A, C, R, l(t), f(t))$ obeys properties (6.15) and (6.16). In particular, for these values of t , spring j plastically expands if $(+, j) \in I_0$ and spring j plastically contracts if $(-, j) \in I_0$.

6.3. Case $I_0 = \{(+, 1), (+, 2), (-, 3)\}$

Relations (6.5), (6.6) and (6.10) take the form

$$\Delta_{I_0} = \begin{pmatrix} 0 & 1 & 1 \\ 0 & 1 & -1 \\ 1 & 0 & 1 \end{pmatrix}^{-1} \circ \begin{pmatrix} 1 & 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 \end{pmatrix} A = \frac{1}{2} \begin{pmatrix} -a_1 & a_2 & a_3 & 0 & 0 \\ a_1 & a_2 & 0 & 0 & 0 \\ a_1 & -a_2 & 0 & 0 & 0 \end{pmatrix},$$

$$\begin{pmatrix} l_1 + \frac{b_3}{a_5} \\ -\frac{b_2}{a_4} + \frac{b_3}{a_5} \\ 0 \end{pmatrix} \in \text{cone} \left(\left\{ \begin{pmatrix} 1 \\ 0 \\ 1 \end{pmatrix}, \begin{pmatrix} 0 \\ 0 \\ -1 \end{pmatrix}, \begin{pmatrix} -1 \\ -1 \\ -1 \end{pmatrix} \right\} \right), \quad Ay_{I_0} = \begin{pmatrix} c_1^+ \\ c_2^+ \\ c_3^- \\ -c_3^- + c_1^+ \\ c_3^- + c_2^+ \end{pmatrix}.$$

Admissibility and irreducibility condition (6.7), feasibility condition (6.11), and convergence time (6.11) now compute as follows.

Admissibility and irreducibility condition:

$$l_1 + \frac{b_2}{a_4} > 0, \quad \frac{b_3}{a_5} + l_1 > 0, \quad \frac{b_2}{a_4} - \frac{b_3}{a_5} > 0. \quad (6.20)$$

Feasibility condition:

$$\begin{aligned} c_4^- &< -tb_2 - d_2 - c_3^- + c_1^+ < c_4^+, \\ c_5^- &< -tb_3 - d_3 + c_3^- + c_2^+ < c_5^+. \end{aligned} \quad (6.21)$$

Convergence time:

$$t \geq \max \left\{ \sqrt{\frac{a_4 a_5 (a_4 a_5 + a_3 (a_4 + a_5))}{a_3 (a_5 b_2 - a_4 b_3)^2}}, \sqrt{\frac{a_1 (a_1 + a_4)}{a_1 (b_2 + a_4 l_1)^2}}, \sqrt{\frac{a_5 (a_2 + a_5)}{a_2 (b_3 + a_5 l_1)^2}} \right\} \|A^{-1} c^+ - A^{-1} c^-\|^A. \quad (6.22)$$

A statement about convergence of the stress vector and about the terminal distribution of plastic deformations given by I_0 comes in direct analogy with the statement of Proposition 6.3.

6.4. Comparison of the cases considered

While the admissibility and irreducibility condition (6.6) for the case $I_0 = \{(+, 2), (+, 3), (+, 4)\}$ does hold when the stress-controlled loading is absent, this no longer the case for the cases $I_0 = \{(+, 2), (-, 3), (+, 4)\}$ and $I_0 = \{(+, 1), (+, 2), (-, 3)\}$ meaning that the latter two cases are possible only when stress-controlled loading is forcing the model of Figure 1. However, the requirements for the stress-controlled loading in cases $I_0 = \{(+, 2), (-, 3), (+, 4)\}$ and $I_0 = \{(+, 1), (+, 2), (-, 3)\}$ are qualitatively opposite. Indeed, we can see that arbitrary small amount of stress-controlled loading is sufficient to realize the case of $I_0 = \{(+, 1), (+, 2), (-, 3)\}$. In this case, the role of stress-controlled loading is to make admissible and reducible $I_0 = \{(+, 1), (+, 2), (-, 3)\}$ irreducible (because in the absence of the stress-controlled loading $I_0 = \{(+, 1), (+, 2), (-, 3)\}$ reduces to $I_0 = \{(+, 1), (+, 2)\}$). The case of $I_0 = \{(+, 2), (-, 3), (+, 4)\}$, in contrast, requires a significant amount of stress-controlled loading to make I_0 admissible (see e.g., the second inequality of (6.17) saying that stress-controlled loading should surpass the displacement-controlled loading). The elastic limits of the remaining springs 1 and 5 should be large enough to accommodate such a large stress-controlled loading, see feasibility condition (6.18).

6.5. Remaining cases of $|I_0| = 3$

We recall the reader that in the case where the model of Figure 1 is forced by just displacement-controlled loading, all admissible irreducible I_0 are [11] $I_0 = \{(+, 1), (+, 2)\}$, $I_0 = \{(+, 4), (+, 5)\}$, $I_0 = \{(+, 1), (-, 3), (+, 5)\}$, $I_0 = \{(+, 2), (+, 3), (+, 4)\}$, see formulas (5.4) and (5.5) for the definitions of admissibility and irreducibility. Addition of stress-controlled loading enlarges this list substantially even in this case of $|I_0| = 3$ that this work sticks to. Indeed, the only $|I_0| = 3$ that are singular (Definition 4.1) are

$$\begin{aligned} I_0 &= \{(\pm, 1), (\pm, 3), (\pm, 4)\}, \\ I_0 &= \{(\pm, 2), (\pm, 3), (\pm, 5)\}, \end{aligned}$$

with all possible combinations of pluses and minuses. And, additionally,

$$I_0 = \{(+, 1), (-, 2), (+, 3)\},$$

$$I_0 = \{(-, 1), (+, 2), (-, 3)\},$$

$$I_0 = \{(+, 3), (-, 4), (+, 5)\},$$

$$I_0 = \{(-, 3), (+, 4), (-, 5)\}$$

are not admissible, i.e., do not satisfy (5.4).

7. Conclusions

We developed an algorithm to determine finite-time convergence of sweeping processes with moving and shape changing polytope to a vertex of the polytope. The earlier results either considered special shapes of polytopes (parallelepipedal) or didn't allow change of shape. As an application we were able to understand the influence of stress-controlled loading on asymptotic (finite-time) elastoplastic behavior of lattice spring model. In particular, we discovered that addition of a small stress-controlled loading can reduce the dimension of attractor that is present in the system otherwise (i.e., add more plastically deforming springs to the springs that deform plastically already), while addition of a larger stress-controlled loading can create such combinations of plastically deforming springs I_0 that none of $I \subset I_0$ constitute an eligible (i.e., admissible and irreducible) set of plastically deforming springs in the absence of stress-controlled loading.

Use of Generative-AI tools declaration

The authors declare they have not used Artificial Intelligence (AI) tools in the creation of this article.

Conflict of interest

The authors declare no conflicts of interest.

References

1. S. Adly, H. Attouch, A. Cabot, Finite time stabilization of nonlinear oscillators subject to dry friction, In: P. Alart, O. Maisonneuve, R. T. Rockafellar, *Nonsmooth mechanics and analysis*, Advances in Mechanics and Mathematics, Springer, **12** (2006), 289–304. https://doi.org/10.1007/0-387-29195-4_24
2. R. B. Bapat, *Graphs and matrices*, Universitext, 1 Ed., Springer, 2010. <https://doi.org/10.1007/978-1-84882-981-7>
3. H. H. Bauschke, P. L. Combettes, *Convex analysis and monotone operator theory in Hilbert spaces*, 2 Eds., Springer, 2017. <https://doi.org/10.1007/978-3-319-48311-5>
4. E. Bernuau, A. Polyakov, D. Efimov, W. Perruquetti, Verification of ISS, iISS and IOSS properties applying weighted homogeneity, *Syst. Control Lett.*, **62** (2013), 1159–1167. <https://doi.org/10.1016/j.sysconle.2013.09.004>
5. S. P. Bhat, D. S. Bernstein, Continuous finite-time stabilization of the translational and rotational double integrators, *IEEE Trans. Automat. Contr.*, **43** (1998), 678–682. <https://doi.org/10.1109/9.668834>

6. C. Bouby, G. de Saxcé, J. B. Tritsch, A comparison between analytical calculations of the shakedown load by the bipotential approach and step-by-step computations for elastoplastic materials with nonlinear kinematic hardening, *Int. J. Solids Struct.*, **43** (2006), 2670–2692. <https://doi.org/10.1016/j.ijsolstr.2005.06.042>
7. G. Colombo, P. Gidoni, E. Vilches, Stabilization of periodic sweeping processes and asymptotic average velocity for soft locomotors with dry friction, *Discrete Contin. Dyn. Syst.*, **42** (2022), 737–757. <https://doi.org/10.3934/dcds.2021135>
8. I. Gudoshnikov, M. Kamenskii, O. Makarenkov, N. Voskovskaia, One-period stability analysis of polygonal sweeping processes with application to an elastoplastic model, *Math. Model. Nat. Phenom.*, **15** (2020), 25. <https://doi.org/10.1051/mmnp/2019030>
9. I. Gudoshnikov, O. Makarenkov, Stabilization of the response of cyclically loaded lattice spring models with plasticity, *ESAIM: COCV*, **27** (2021), S8. <https://doi.org/10.1051/cocv/2020043>
10. I. Gudoshnikov, O. Makarenkov, Structurally stable families of periodic solutions in sweeping processes of networks of elastoplastic springs, *Phys. D*, **406** (2020), 132443. <https://doi.org/10.1016/j.physd.2020.132443>
11. I. Gudoshnikov, O. Makarenkov, D. Rachinskii, Finite-time stability of polyhedral sweeping processes with application to elastoplastic systems, *SIAM J. Control Optim.*, **60** (2022), 1320–1346. <https://doi.org/10.1137/20M1388796>
12. I. Gudoshnikov, O. Makarenkov, D. Rachinskii, Formation of a nontrivial finite-time stable attractor in a class of polyhedral sweeping processes with periodic input, *ESAIM: COCV*, **29** (2023), 84. <https://doi.org/10.1051/cocv/2023074>
13. P. Krejčí, *Hysteresis, convexity and dissipation in hyperbolic equations*, Tokyo: Gakkotosho, 1996.
14. M. Kunze, M. D. P. M. Marques, An introduction to Moreau's sweeping process, In: B. Brogliato, *Impacts in mechanical systems*, Lecture Notes in Physics, Springer, **551** (2000), 1–60. https://doi.org/10.1007/3-540-45501-9_1
15. J. J. Moreau, On unilateral constraints, friction and plasticity, In: G. Capriz, G. Stampacchia, *New variational techniques in mathematical physics*, C.I.M.E. Summer Schools, Springer, **63** (1974), 171–322. https://doi.org/10.1007/978-3-642-10960-7_7
16. H. B. Oza, Y. V. Orlov, S. K. Spurgeon, Continuous uniform finite time stabilization of planar controllable systems, *SIAM J. Control Optim.*, **53** (2015), 1154–1181. <https://doi.org/10.1137/120877155>
17. D. Rachinskii, On geometric conditions for reduction of the Moreau sweeping process to the Prandtl-Ishlinskii operator, *Discrete Contin. Dyn. Syst.*, **23** (2018), 3361–3386. <https://doi.org/10.3934/dcdsb.2018246>
18. R. T. Rockafellar, R. J. B. Wets, *Variational analysis*, Grundlehren der mathematischen Wissenschaften, Vol. 317, Springer, 1998. <https://doi.org/10.1007/978-3-642-02431-3>
19. T. Sanchez, J. A. Moreno, L. M. Fridman, Output feedback continuous twisting algorithm, *Automatica*, **96** (2018), 298–305. <https://doi.org/10.1016/j.automatica.2018.06.049>
20. H. S. Yu, *Plasticity and geotechnics*, Advances in Mechanics and Mathematics, Vol. 13, Springer, 2006. <https://doi.org/10.1007/978-0-387-33599-5>

21. N. Zouain, R. SantAnna, Computational formulation for the asymptotic response of elastoplastic solids under cyclic loads, *Eur. J. Mech.*, **61** (2017), 267–278. <https://doi.org/10.1016/j.euromechsol.2016.09.013>

Appendix

Proof of Theorem 2.1. This proof follows the proof of [11, Theorem 3.1] with C replaced by $C(t) - c(t)$.

Let $y(t)$ be an arbitrary solution of (1.4). For the function $x(t)$ given by (1.10) consider

$$v(t) = V(x(t)).$$

Note, that $x(t)$ is differentiable almost everywhere on $[0, \infty)$ because $c(t)$ is Lipschitz continuous. Let us fix some $t \geq 0$ such that $x(t)$ is differentiable at t . Without loss of generality we can assume that $t \geq 0$ is chosen also so that $V(x(t))$ is differentiable at t . Then

$$v'(t) = 2 \langle x'(t), A(x(t) - y_*) \rangle. \quad (\text{A.1})$$

By the definition of normal cone, (2.6) implies

$$\langle -\dot{x}(t) - \dot{c}(t), A(\xi - x(t)) \rangle \leq 0, \quad \text{for any } \xi \in C(t) - c(t).$$

Therefore, taking $\xi = y_*$ we conclude from (A.1) that

$$v'(t) \leq 2 \langle -c'(t), A(x(t) - y_*) \rangle. \quad (\text{A.2})$$

Now we use assumption (1.9), which is equivalent to

$$-c'(t) + \varepsilon \frac{\zeta}{\|\zeta\|^A} \in N_{C(t)-c(t)}^A(y_*), \quad \text{for any } \zeta \in V,$$

or, using the definition of the normal cone,

$$\left\langle -c'(t) + \varepsilon \frac{\zeta}{\|\zeta\|^A}, A(\xi - y_*) \right\rangle \leq 0, \quad \text{for any } \zeta \in \mathbb{R}^m, \xi \in C(t) - c(t).$$

Therefore, letting $\xi = x(t)$ and $\zeta = x(t) - y_*$, we get

$$\left\langle -c'(t) + \varepsilon \frac{x(t) - y_*}{\|x(t) - y_*\|^A}, A(x(t) - y_*) \right\rangle \leq 0,$$

which allows to further rewrite inequality (A.2) as

$$v'(t) \leq -2\varepsilon \left\langle \frac{x(t) - y_*}{\|x(t) - y_*\|^A}, A(x(t) - y_*) \right\rangle = -2\varepsilon \sqrt{v(t)}.$$

Therefore, the Lyapunov function (2.5) satisfies estimate (1.1). The proof is complete. \square



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